

FHI-aims User and Developer Workshop 2023, Hamburg, Germany

Lessons learned in
Heat and Charge Transport
over the last couple of years

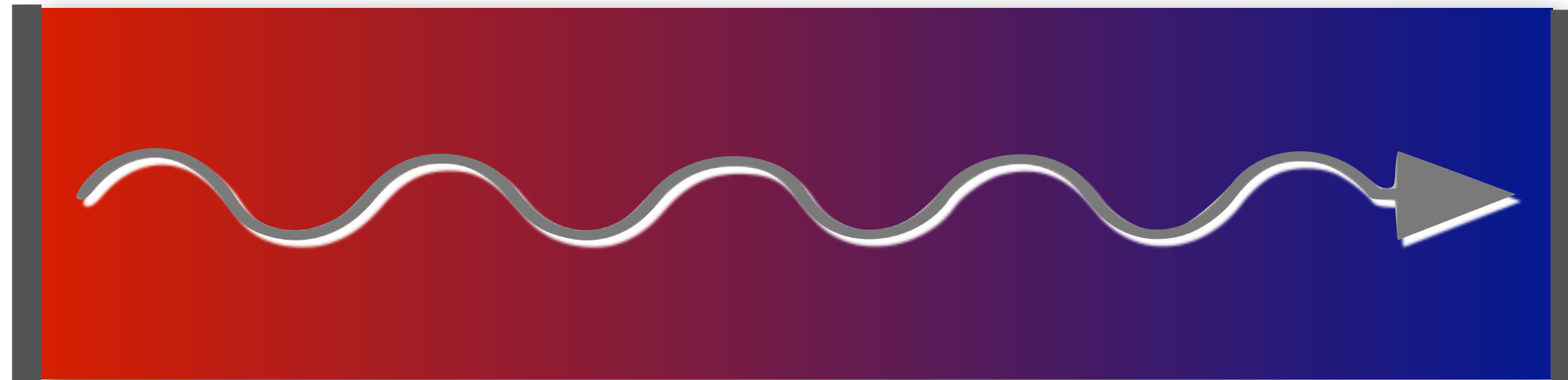
Christian Carbogno

FRITZ-HABER-INSTITUT
MAX-PLANCK-GESELLSCHAFT



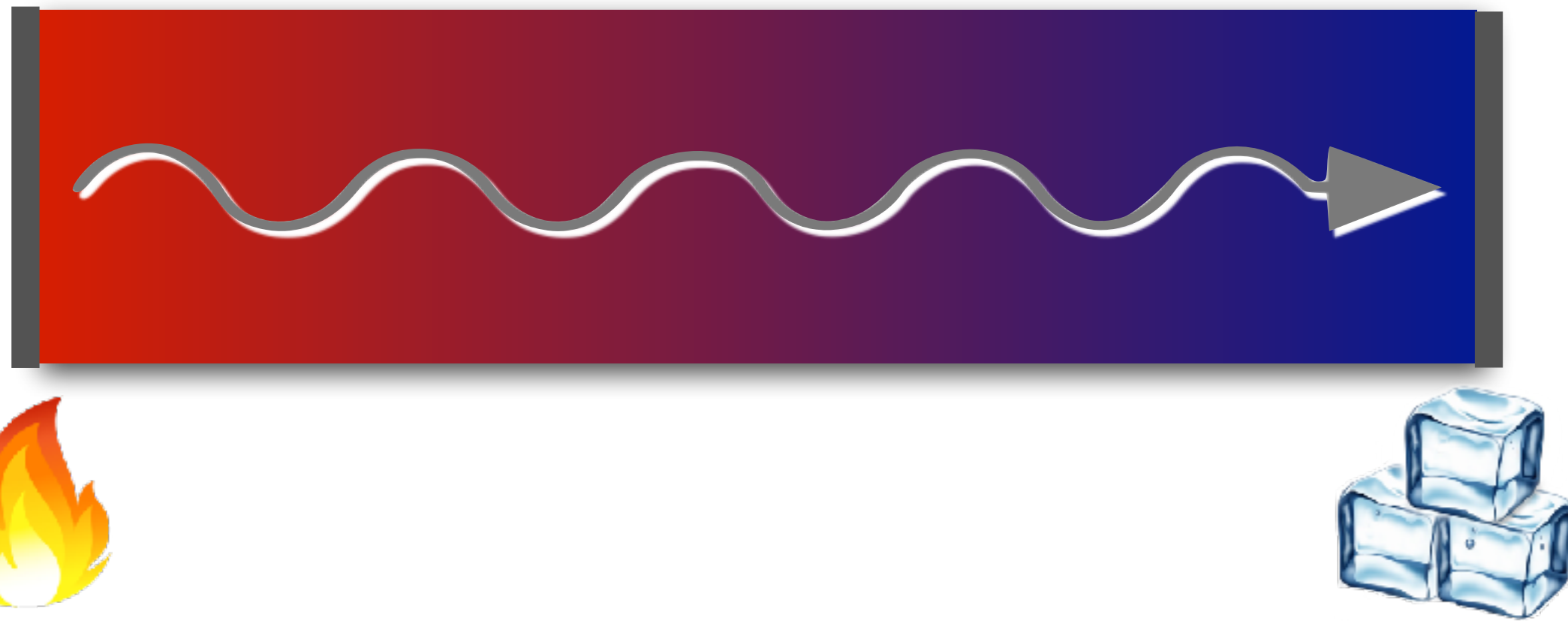
Heat & Charge Transport in Solids

Heat Flux: Photons, Electrons, and Nuclei

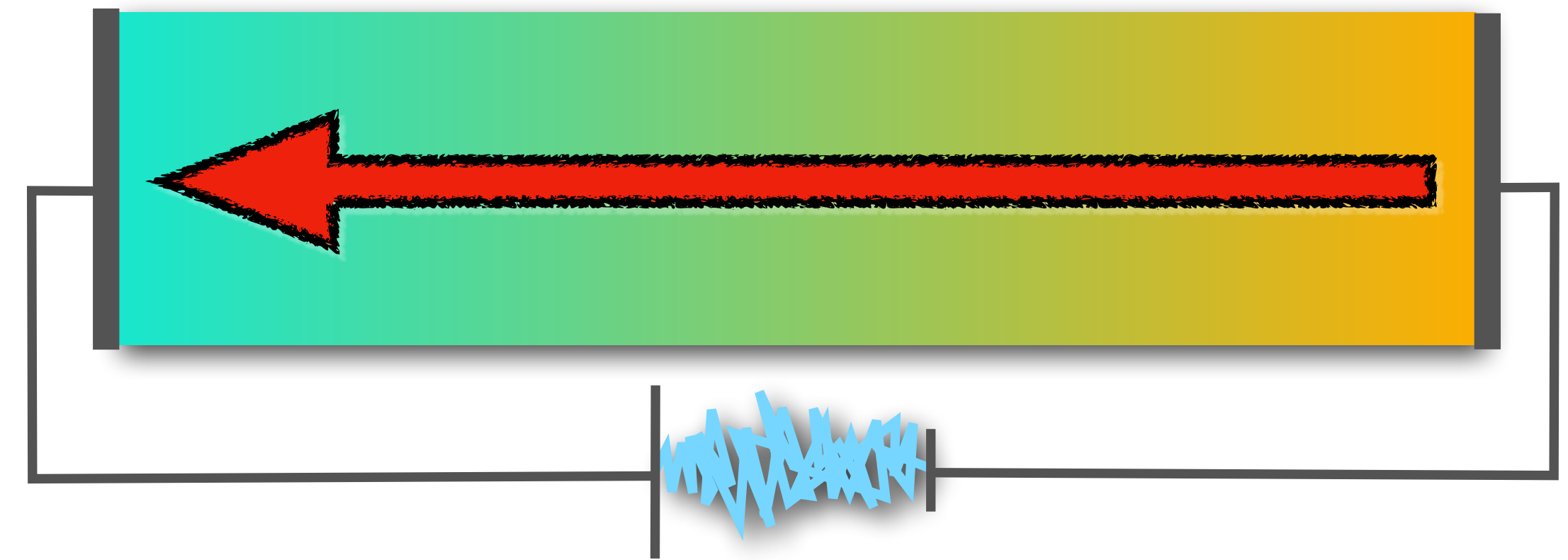


Heat & Charge Transport in Solids

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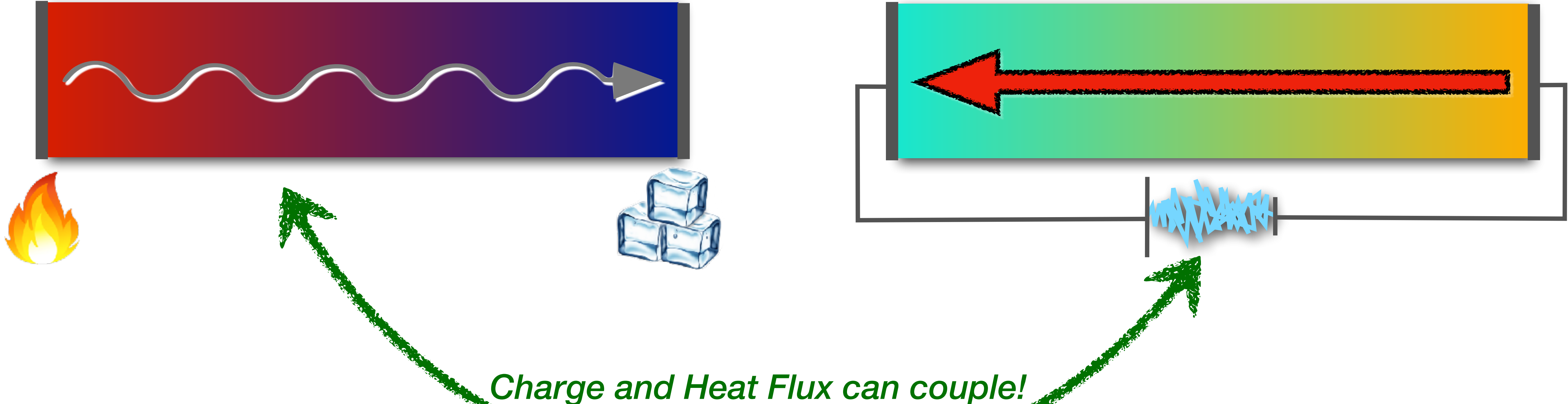
Charge Flux: Electrons & Ions



Heat & Charge Transport in Solids

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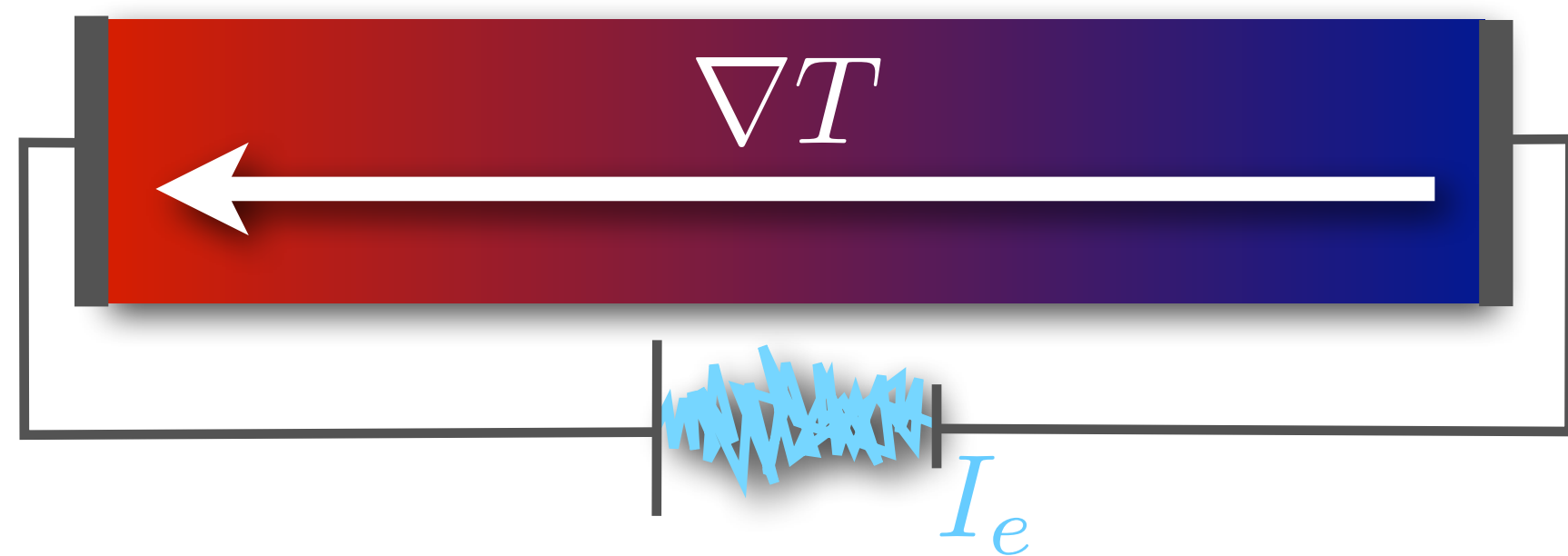
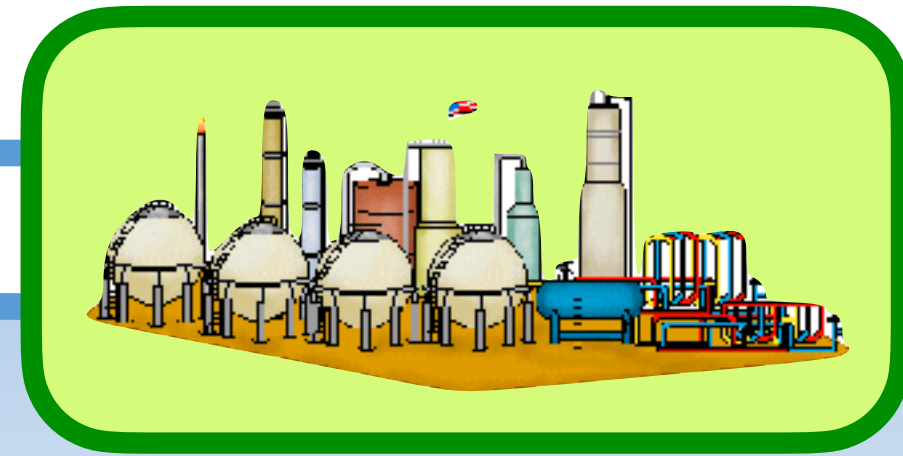
Charge Flux: Electrons & Ions



Charge and Heat Flux can couple!

Thermoelectric Materials:
Seebeck Effect: Heat flux induces a Potential Gradient
Peltier Effect: Charge Flux induces a Temperature Gradient

Thermoelectric Materials

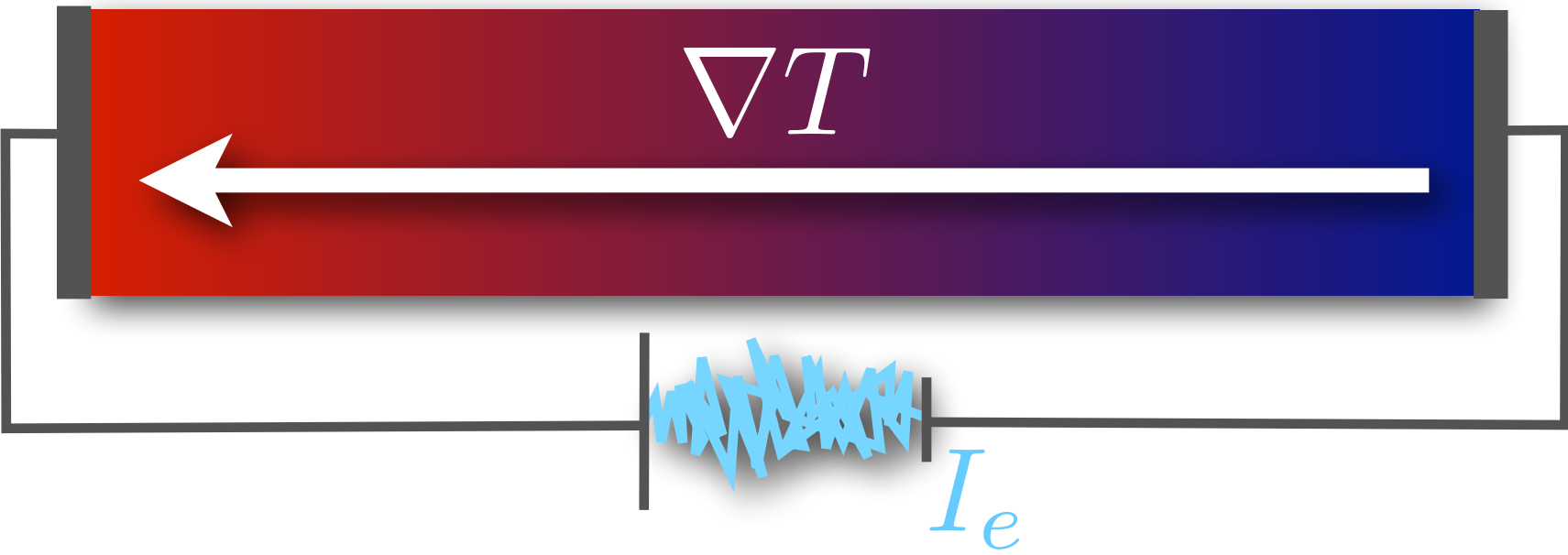


Thermoelectric materials are able to convert **waste-heat** into useful **voltage**.
=> **Substantial energy efficiency gains.**

G. J. Snyder and E. S. Toberer, *Nat. Mater.* **7**, 105 (2008).

More **efficient thermoelectric materials** can **enable** these applications!

Thermoelectric Materials



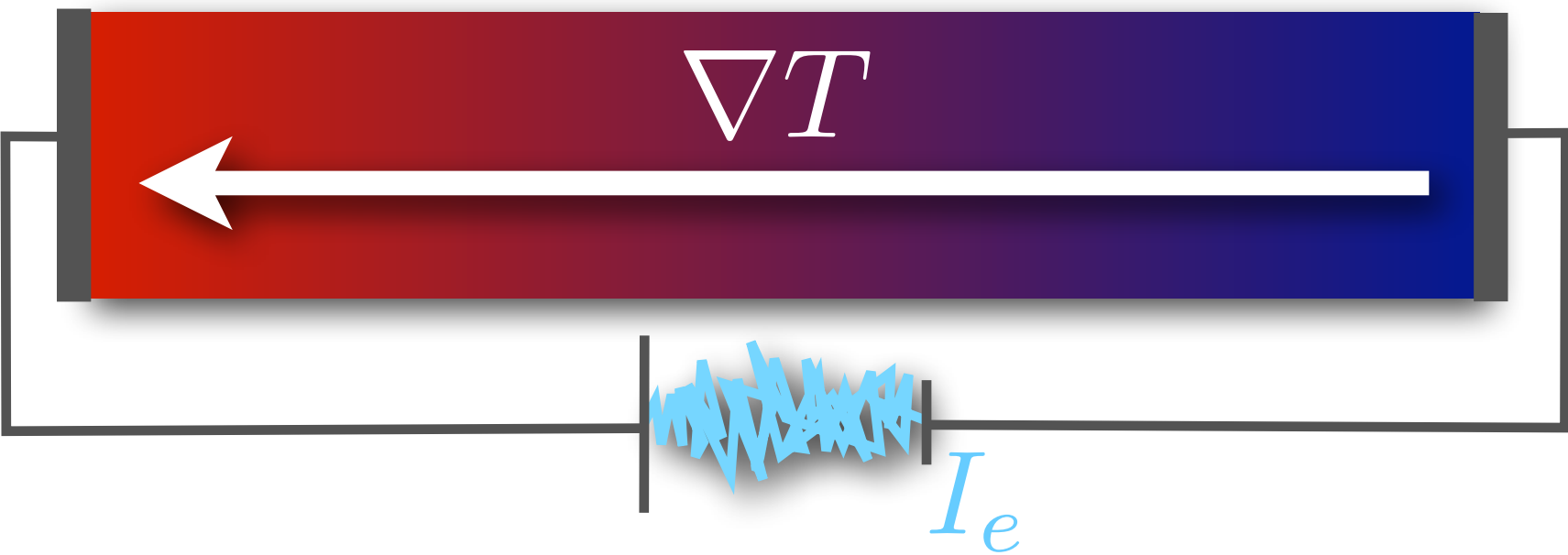
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Minimize *vibrational* heat transport!

Thermoelectric Materials

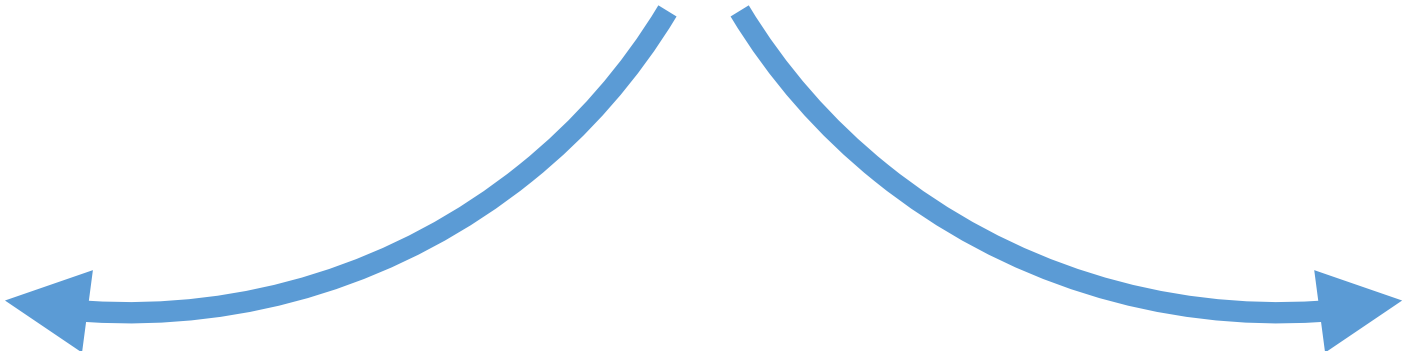


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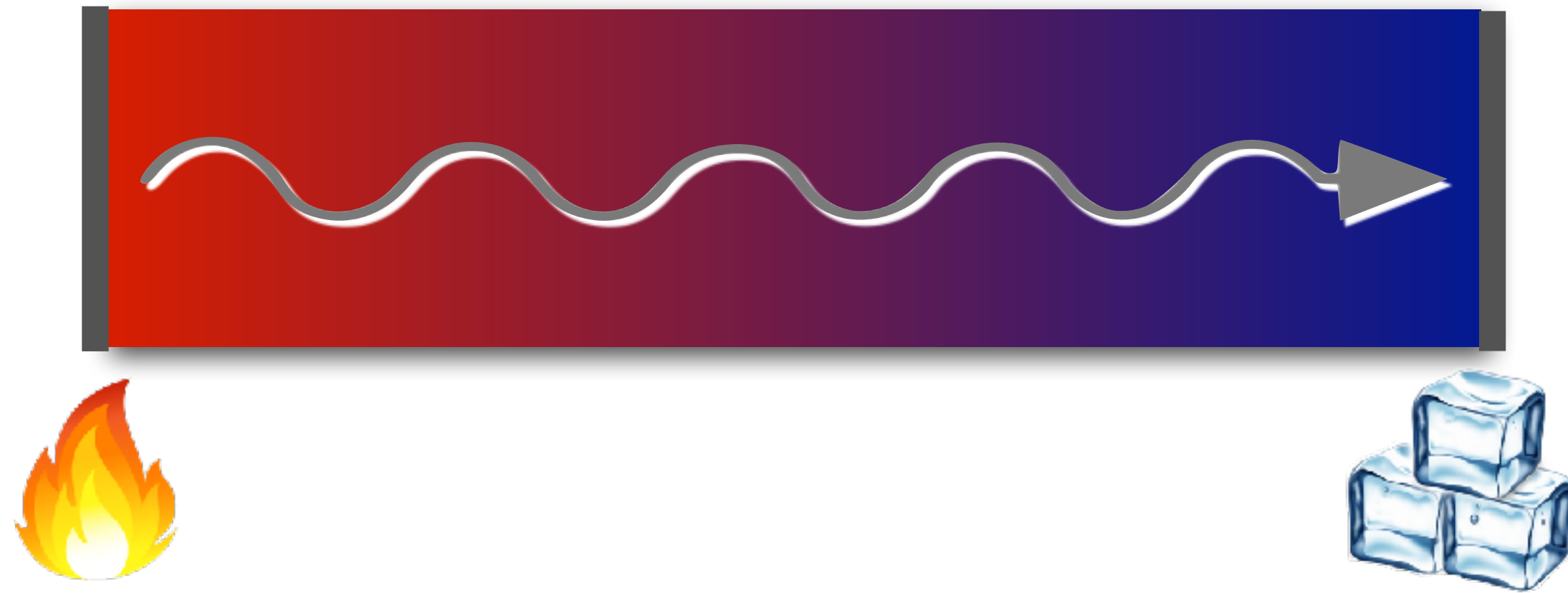
Minimize *vibrational* heat transport!



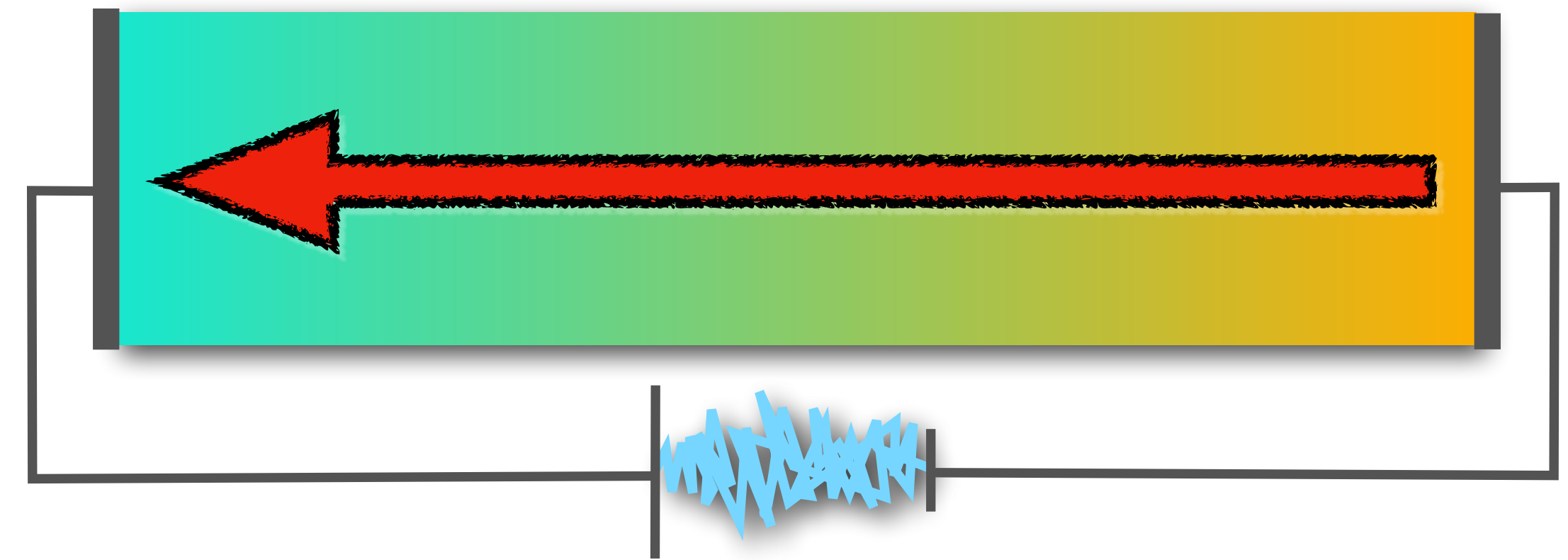
Maximize *charge* transport!

Heat & Charge Transport in Solids

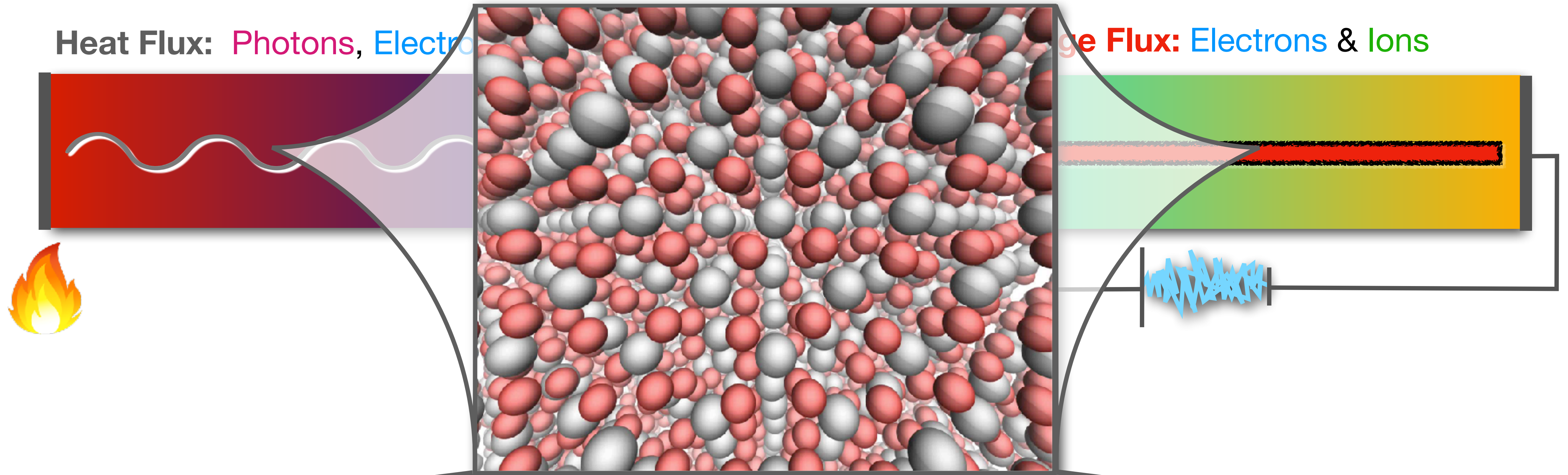
Heat Flux: Photons, Electrons, and Nuclei



Charge Flux: Electrons & Ions



Heat & Charge Transport in Solids



Nuclear Motion causes Dissipation!

Fourier's Law:

$$\mathbf{J}^e = -\kappa \nabla T$$

This Talk:

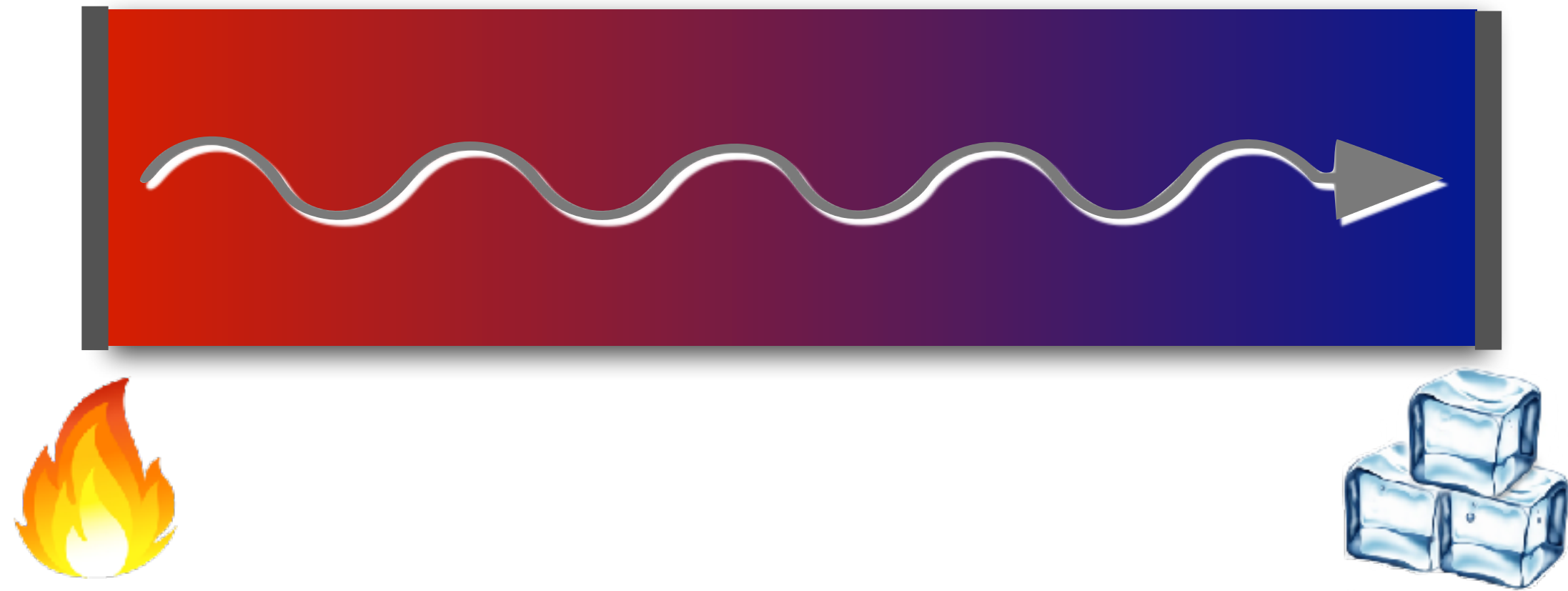
How can we *accurately and predictively* compute, understand, and tailor κ and σ ?

Ohm's Law:

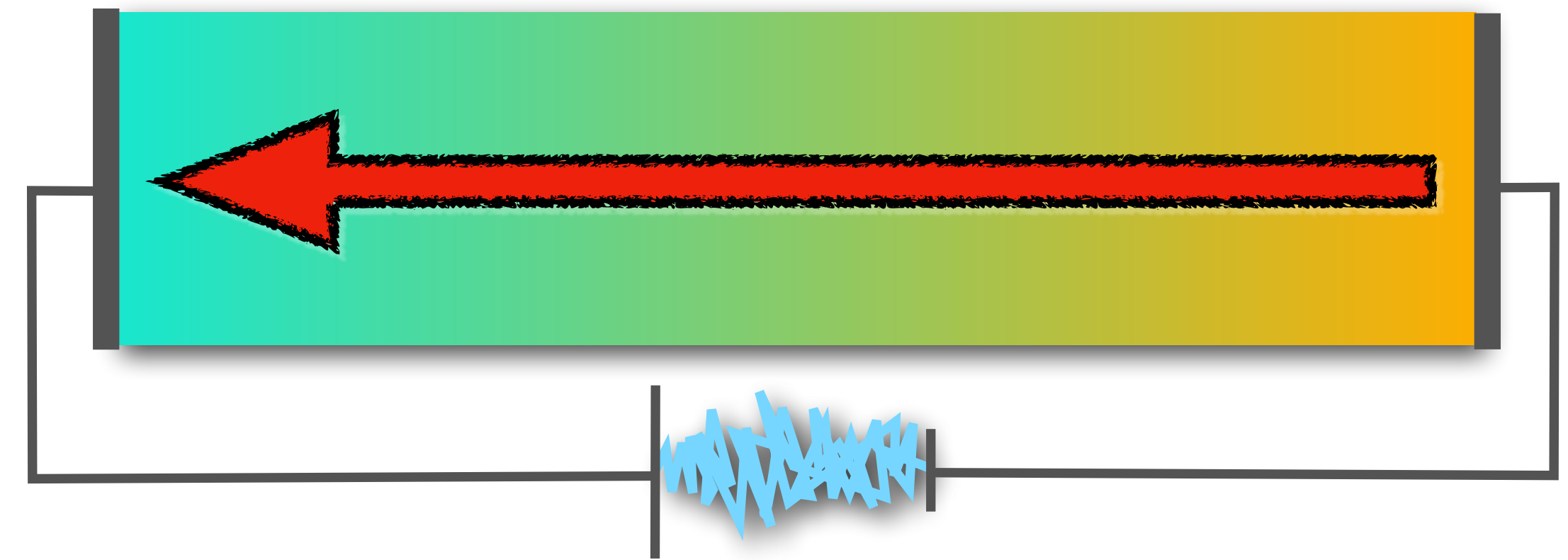
$$\mathbf{J}^c = -\sigma \nabla U$$

Heat & Charge Transport in Solids

Heat Flux: Photons, Electrons, and Nuclei

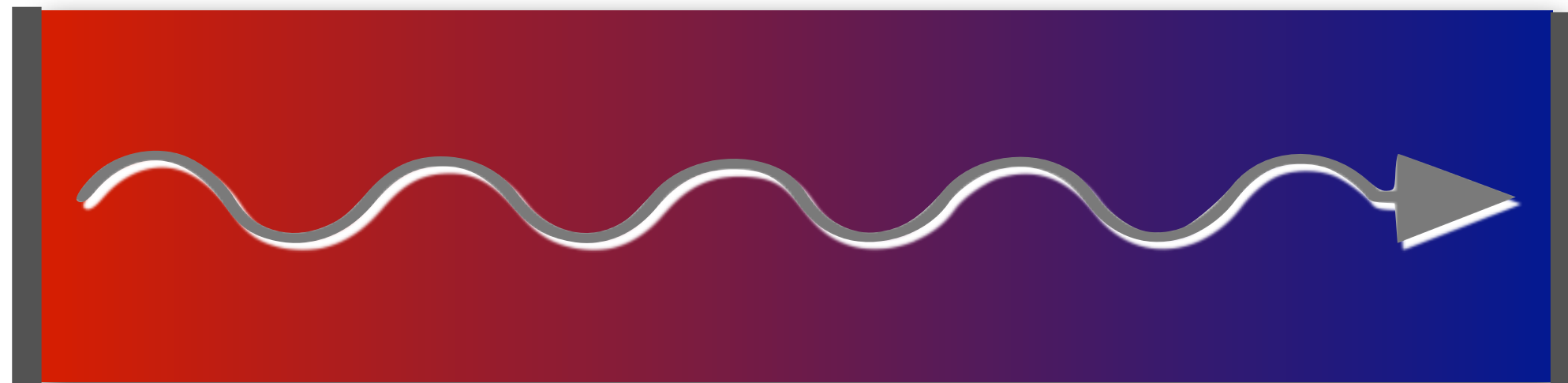


Charge Flux: Electrons & Ions



Heat & Charge Transport in Solids

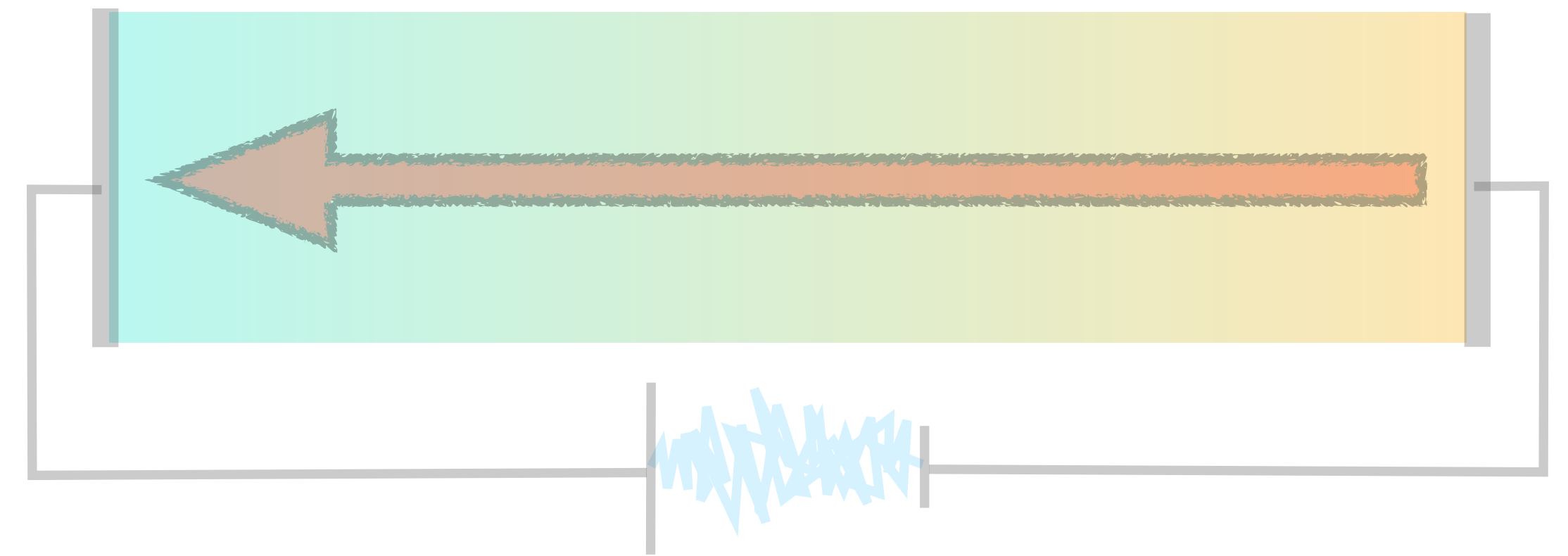
Heat Flux: Photons, Electrons, and Nuclei



$$\kappa = \kappa_{\text{photon}} + \kappa_{\text{elec.}} + \kappa_{\text{nucl.}}$$

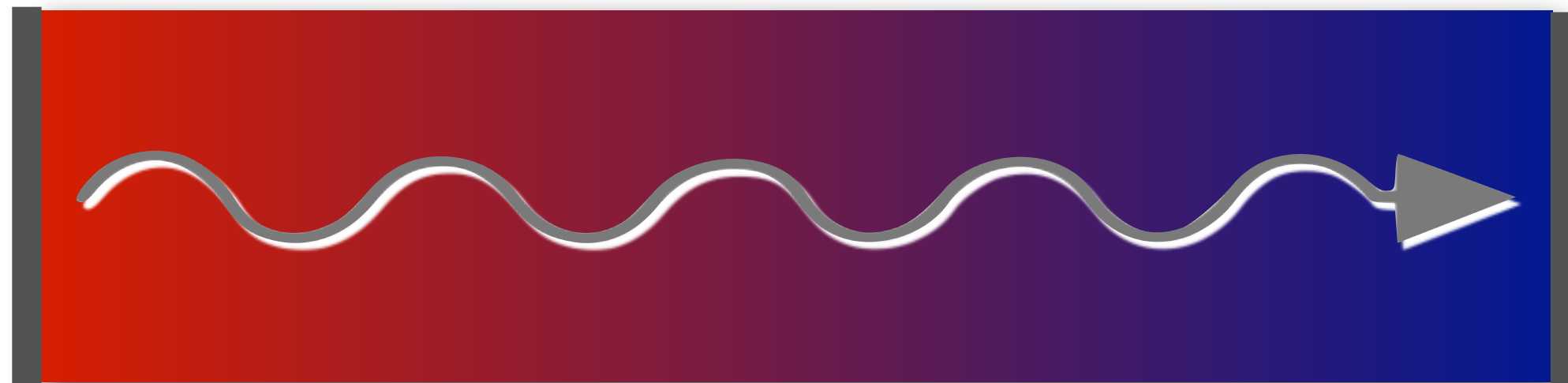


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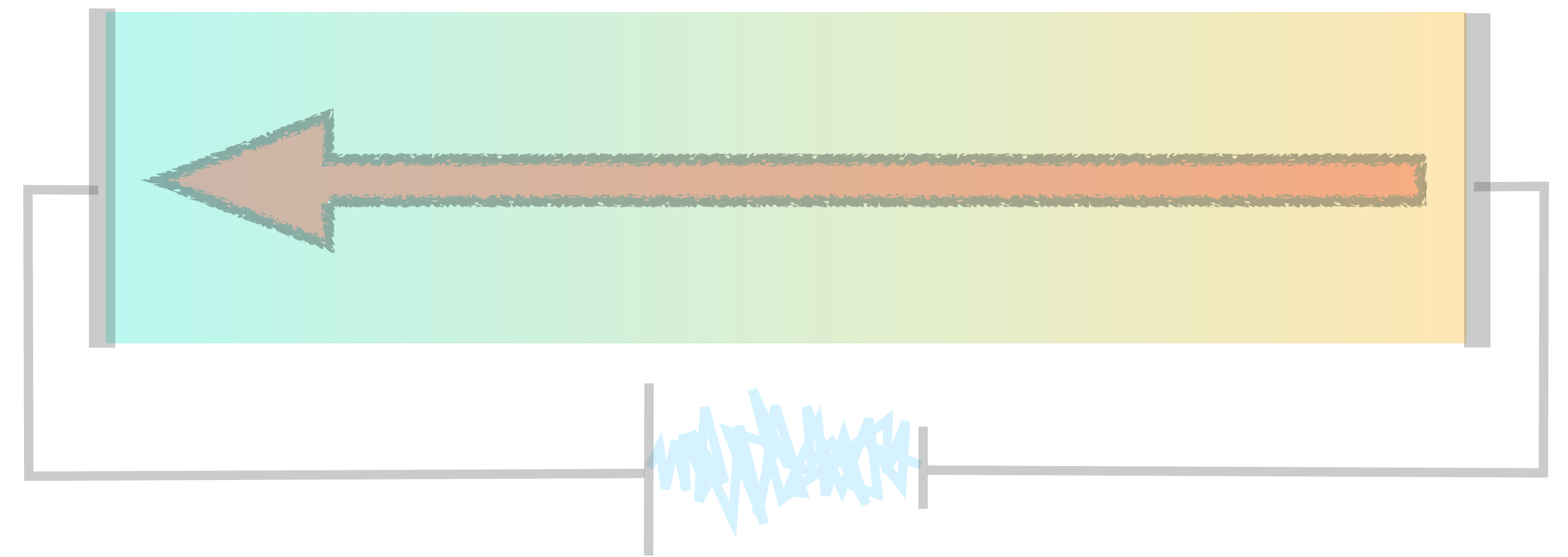


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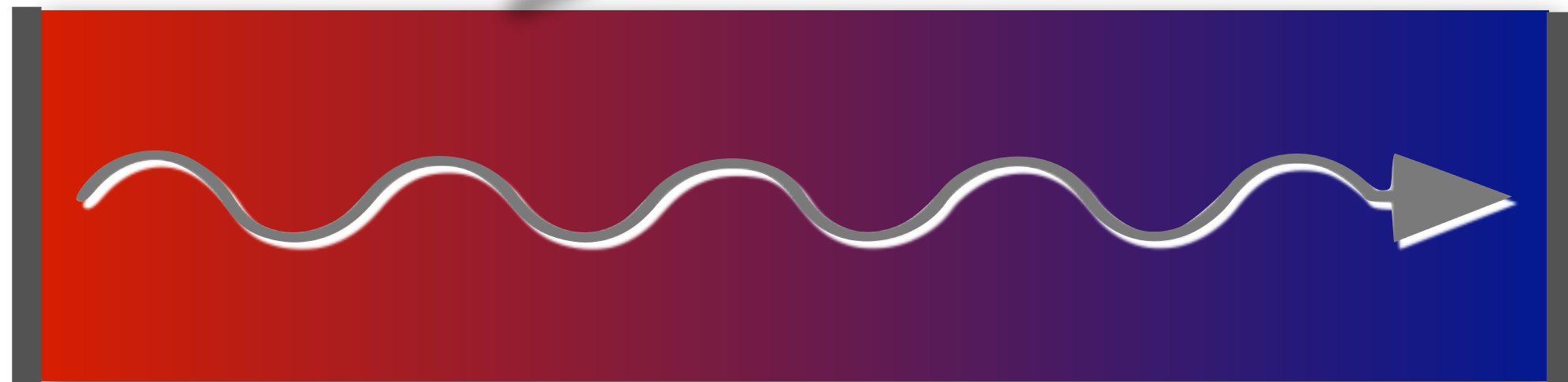
Relevant in **transparent, incandescent** materials!

Charge Flux: Electrons & Ions



Heat & Charge Transport in Solids

Heat Flux: ~~Photons~~, Electrons, and Nuclei



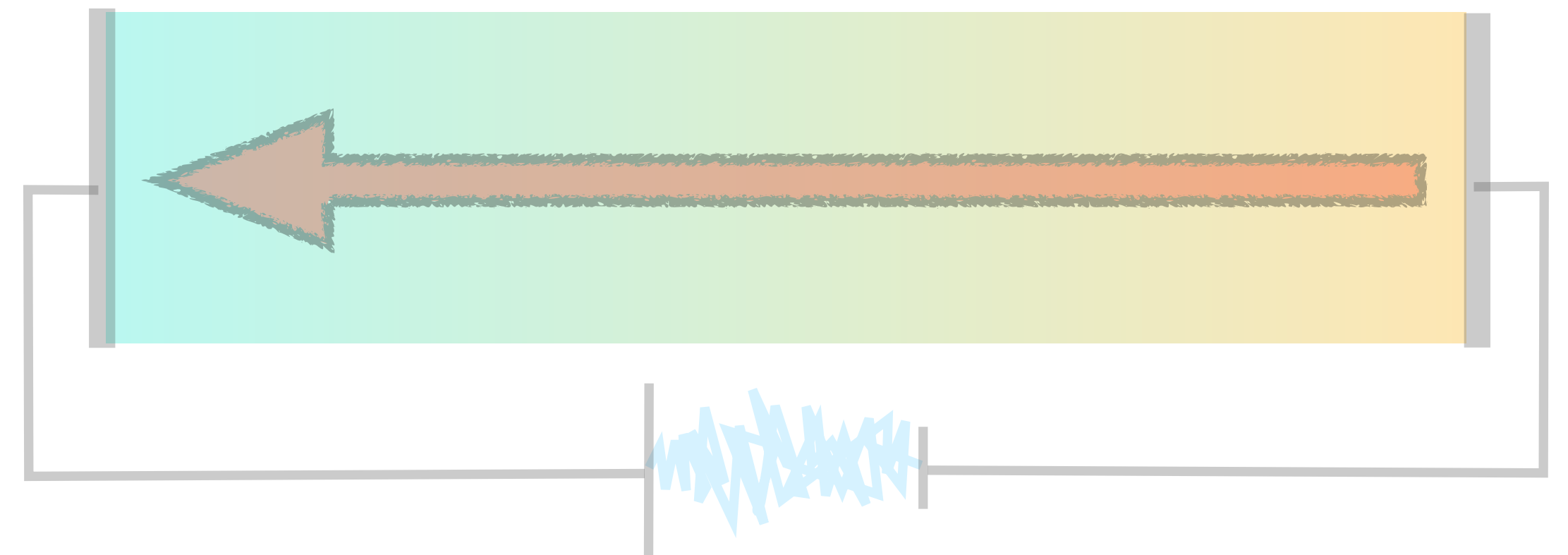
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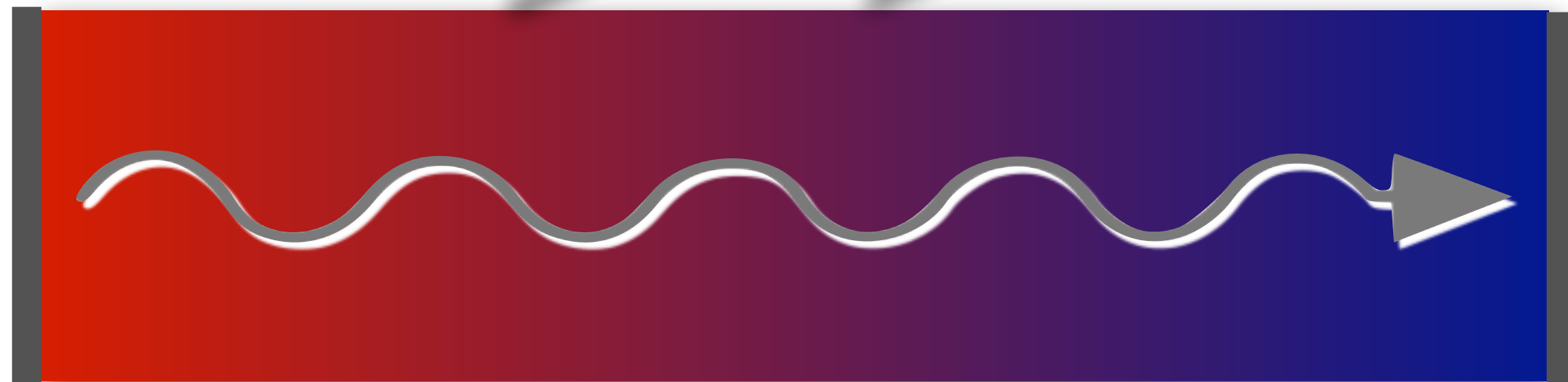
Relevant in **metals** and **degenerate semiconductors**.

Charge Flux: Electrons & Ions



Heat & Charge Transport in Solids

Heat Flux: ~~Photons~~, ~~Electrons~~, and Nuclei



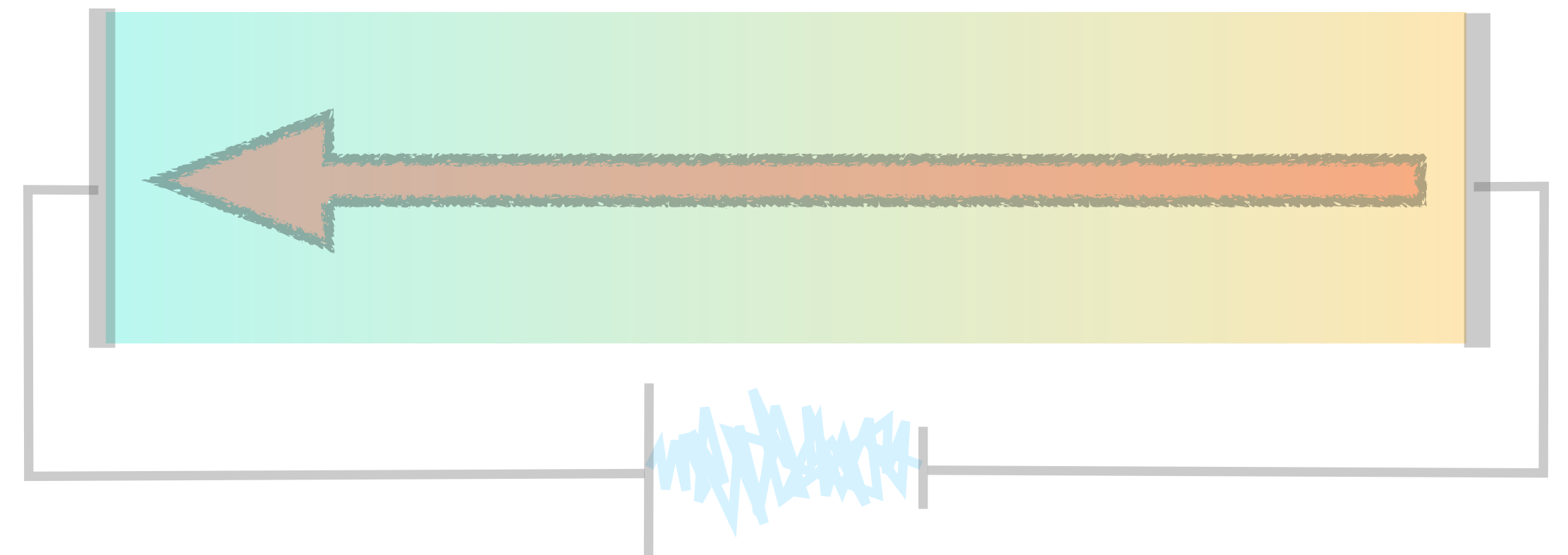
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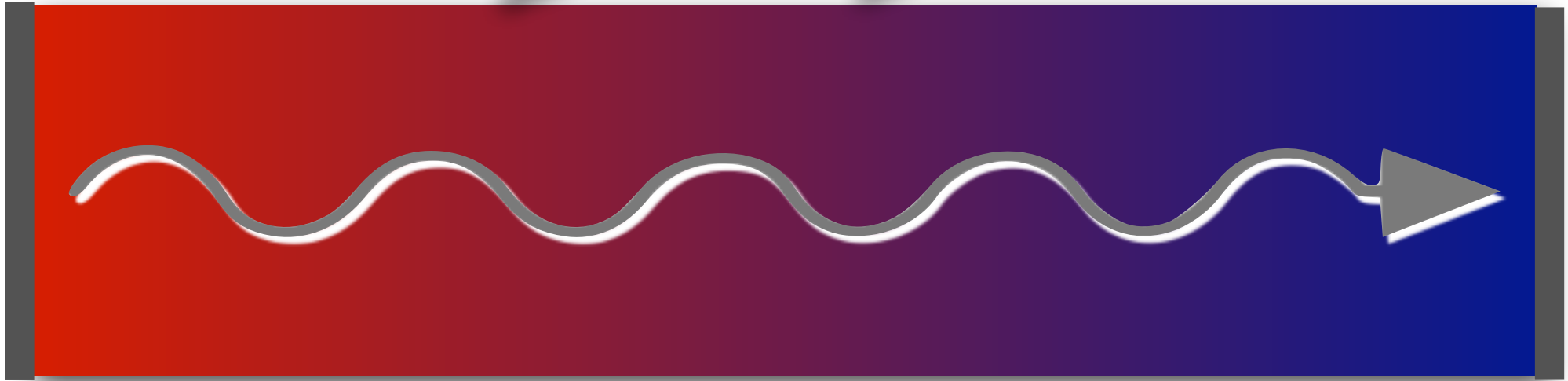
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Charge Flux: Electrons & Ions

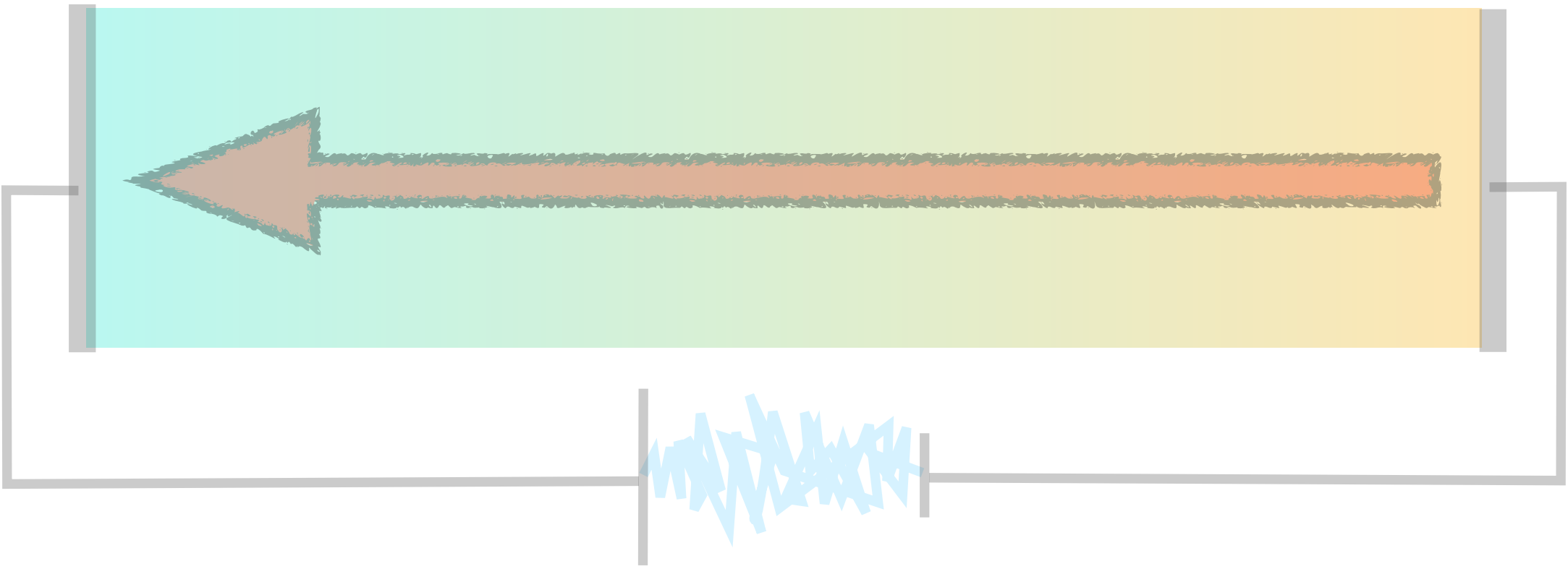


Heat & Charge Transport in Solids

Heat Flux: ~~Photons~~, ~~Electrons~~, and Nuclei



Charge Flux: Electrons & Ions



$$\kappa = \kappa_{\text{photon}} + \kappa_{\text{elec.}} + \kappa_{\text{nucl.}}$$

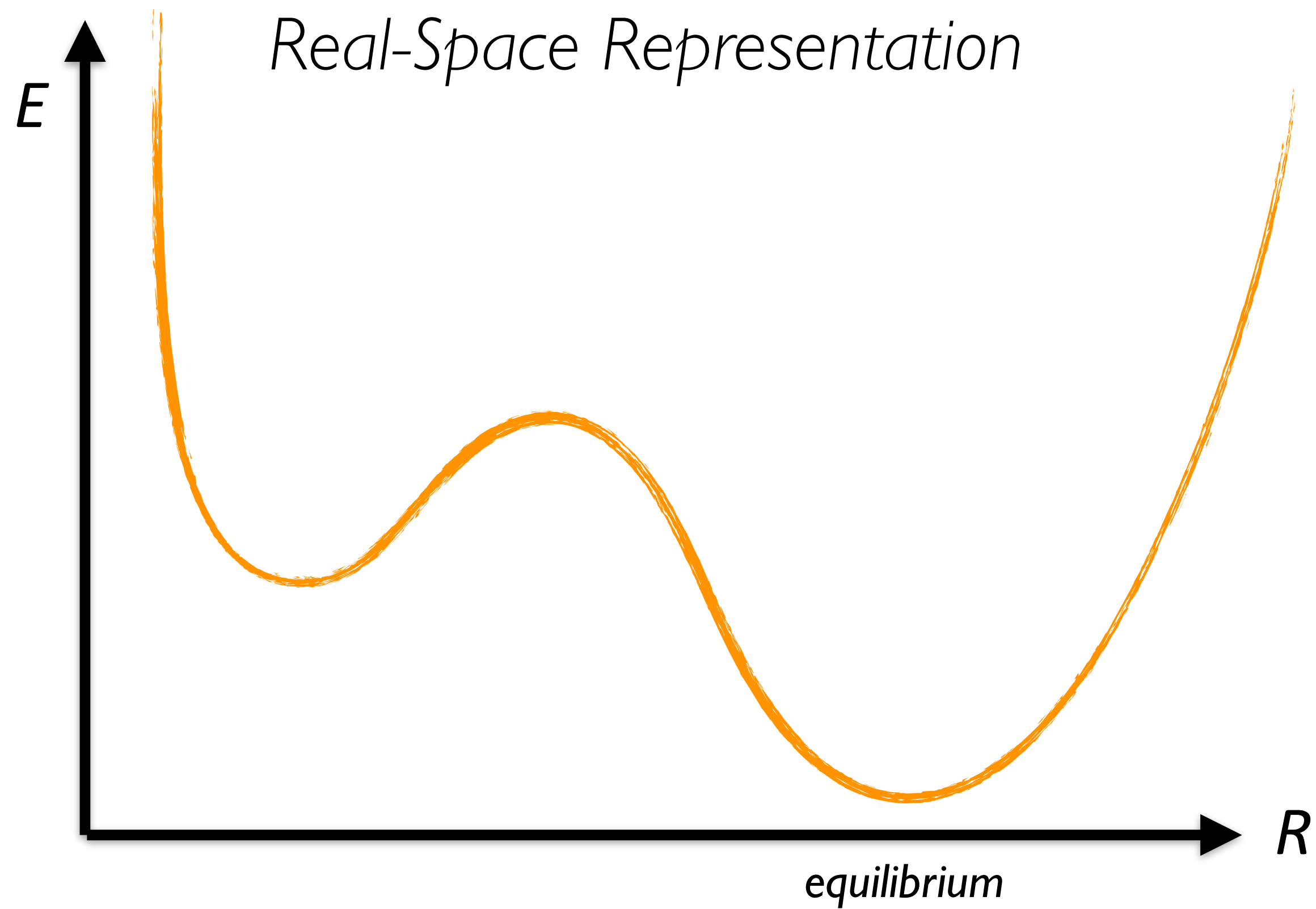


Relevant in transparent, incandescent materials!

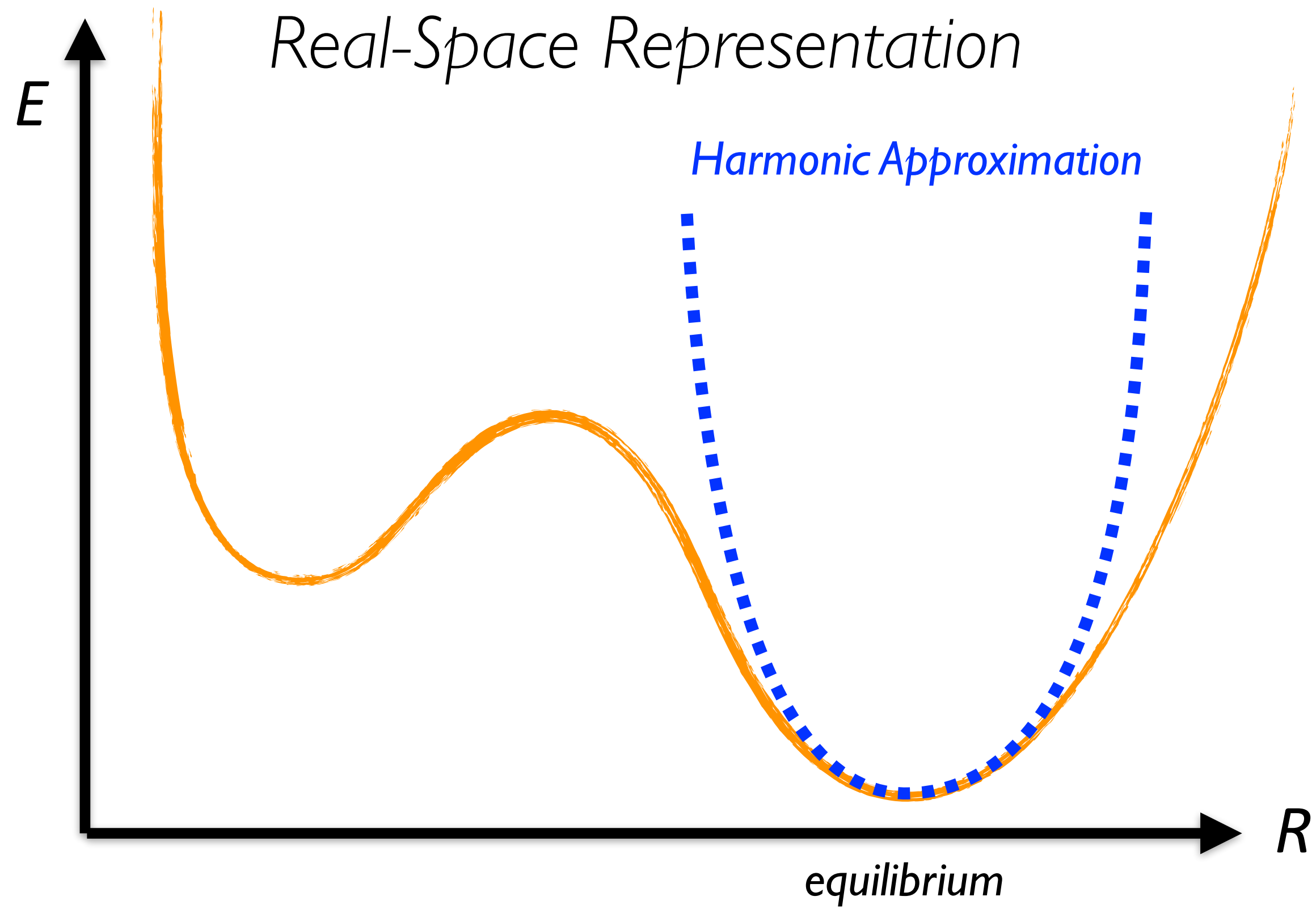
Relevant in metals and degenerate semiconductors.

The dominant contribution in insulators and non-degenerate semiconductors.

Heat Transport Theory 101

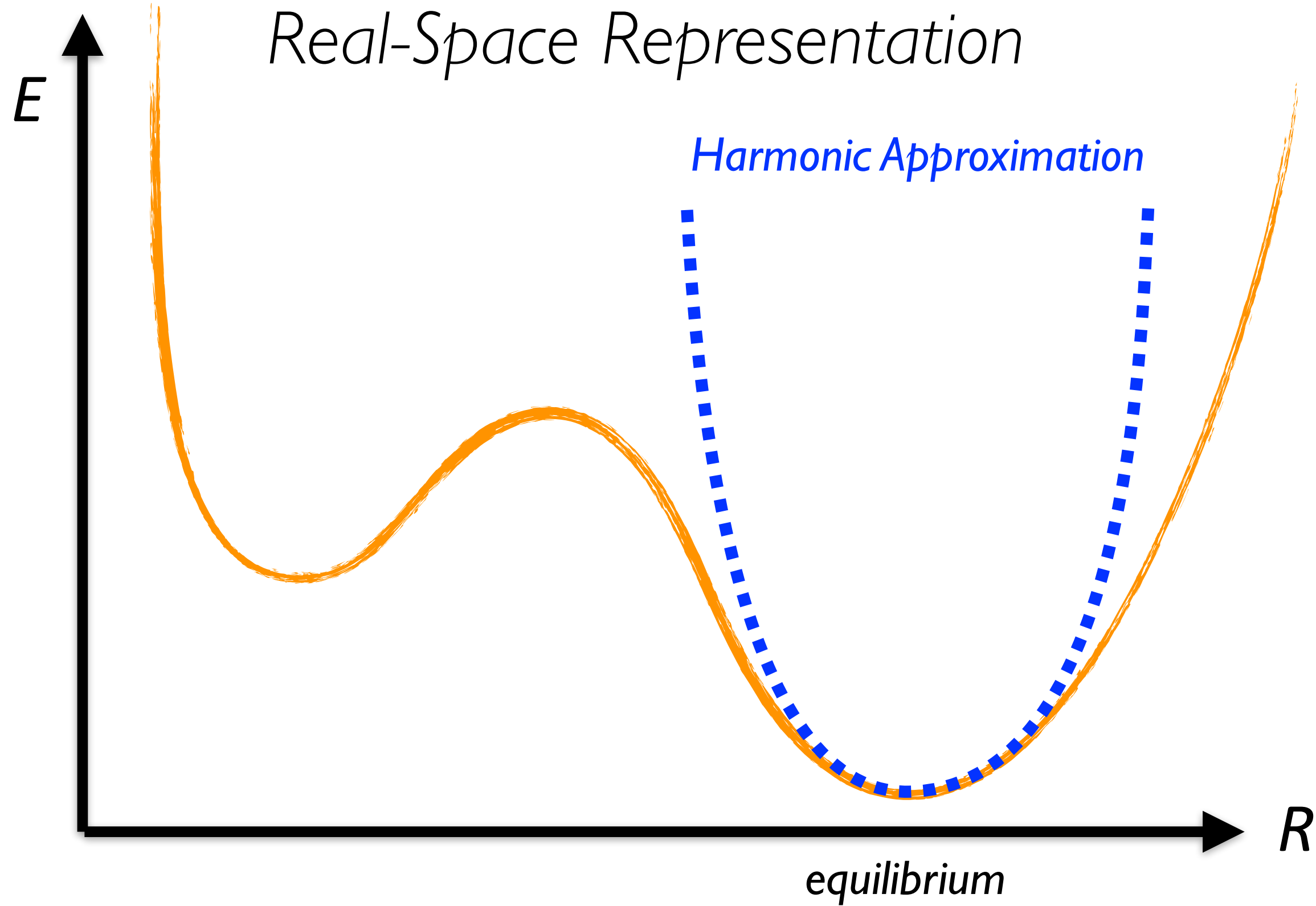


Heat Transport Theory 101

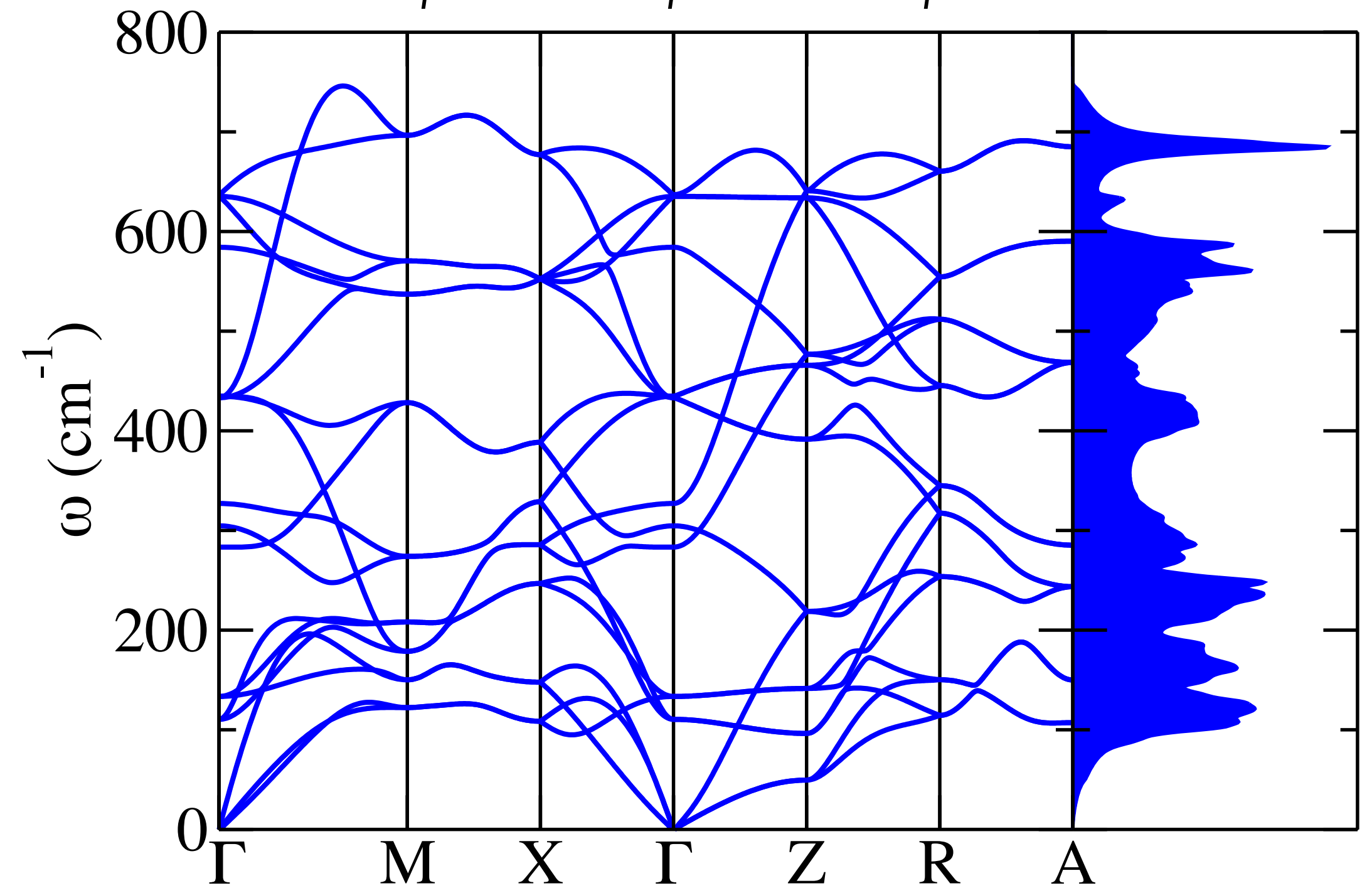


Heat Transport Theory I 01

Real-Space Representation

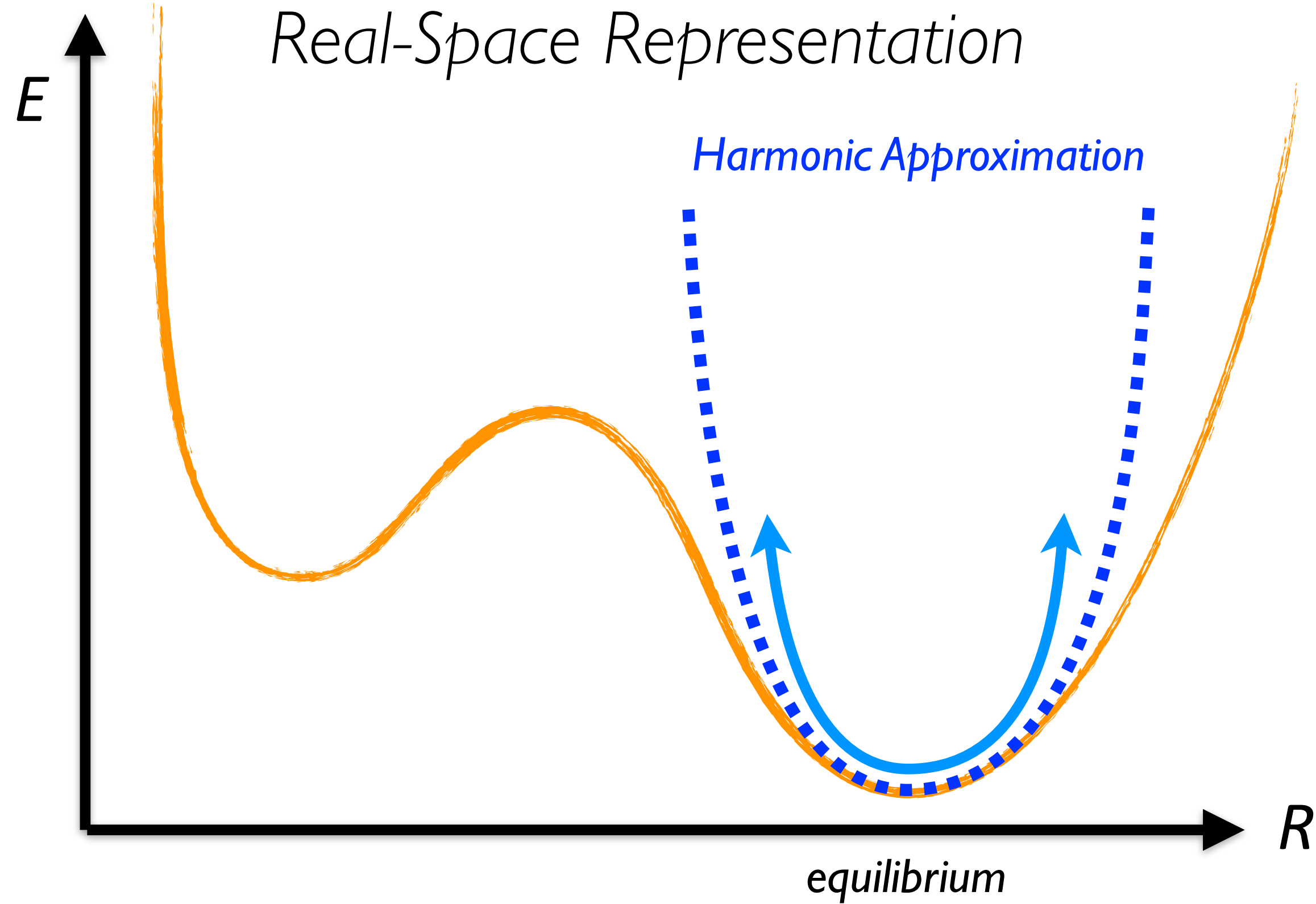


Reciprocal-Space Representation



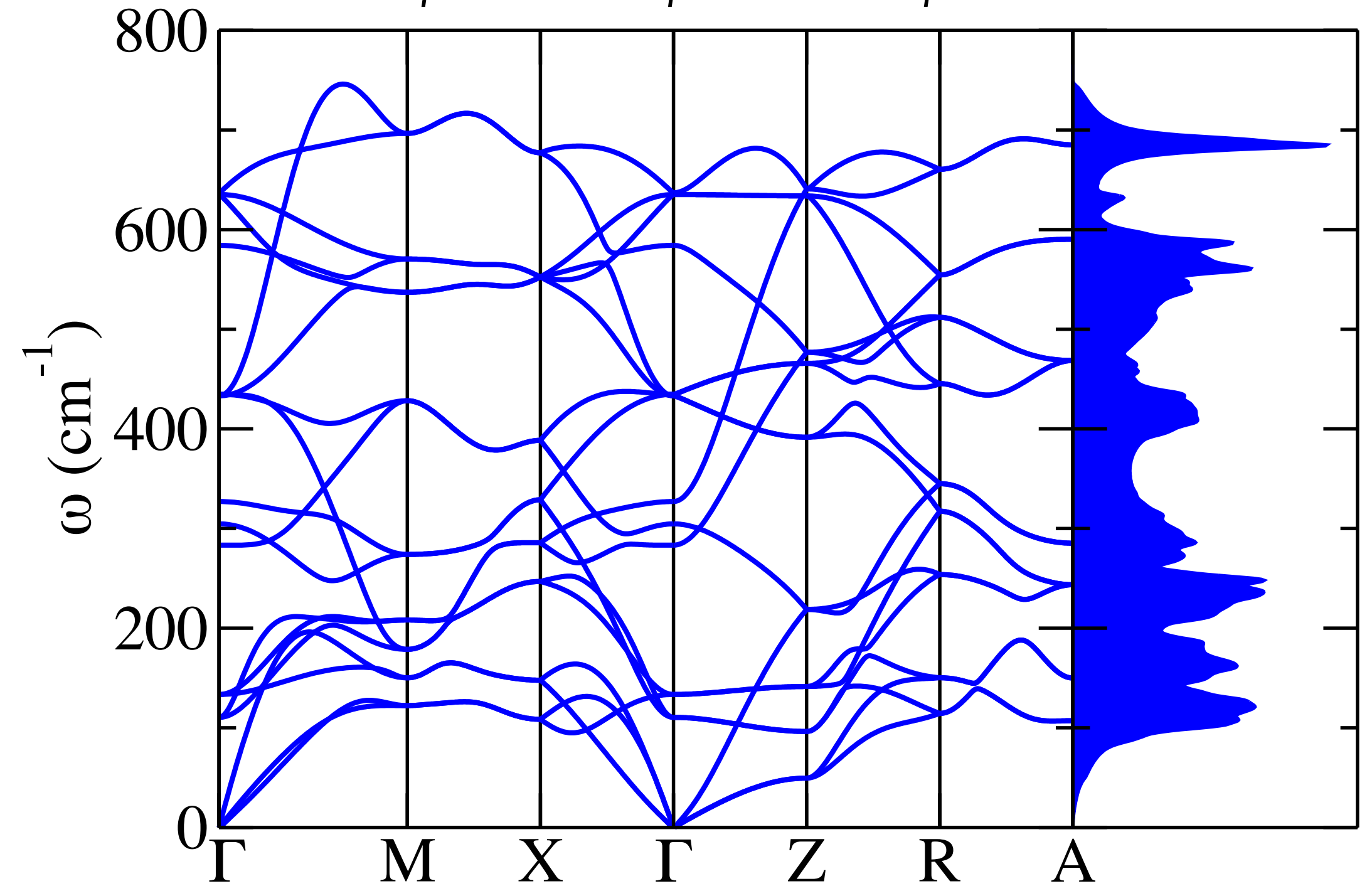
Heat Transport Theory 101

Real-Space Representation



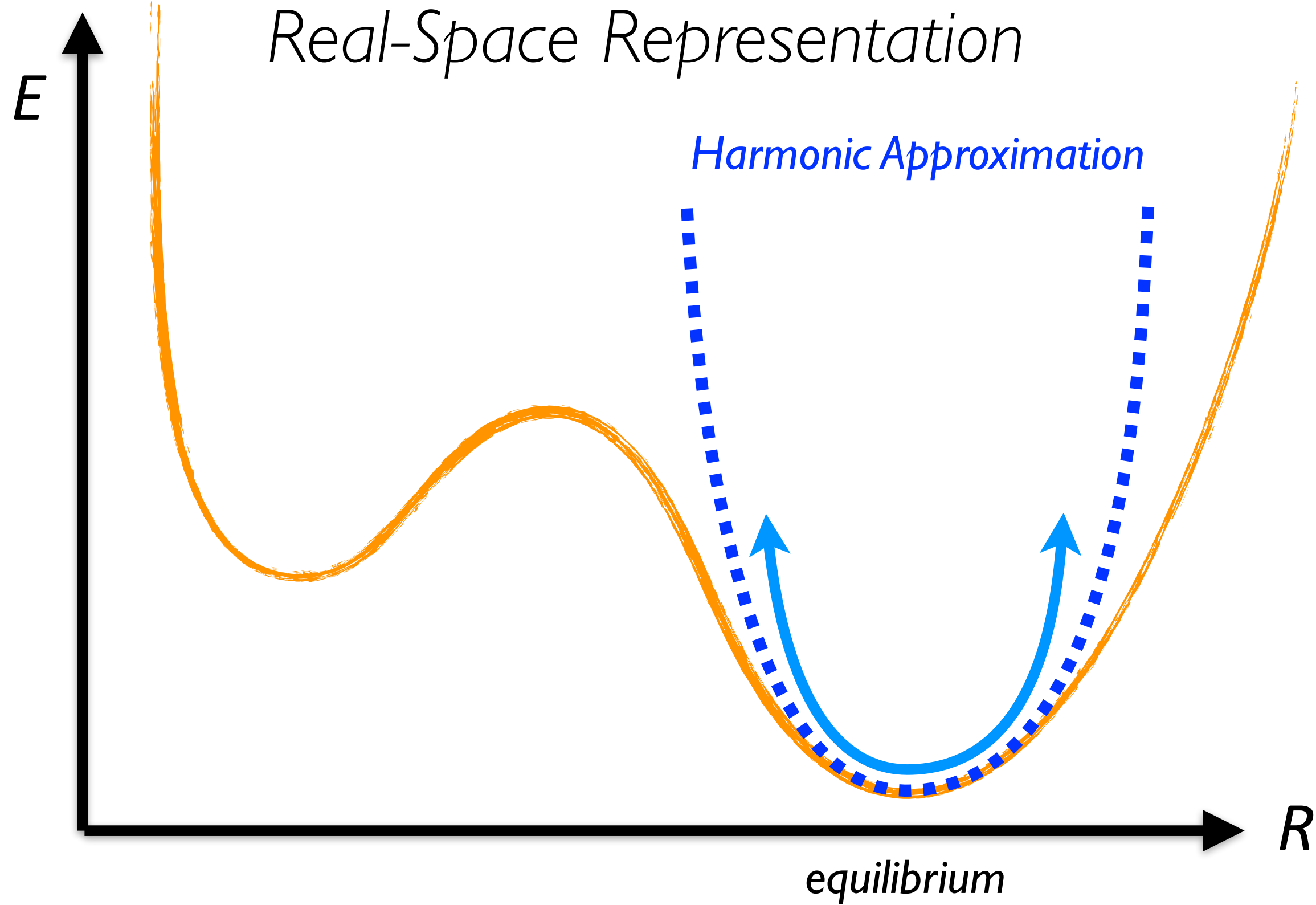
Decoupled Modes

Reciprocal-Space Representation



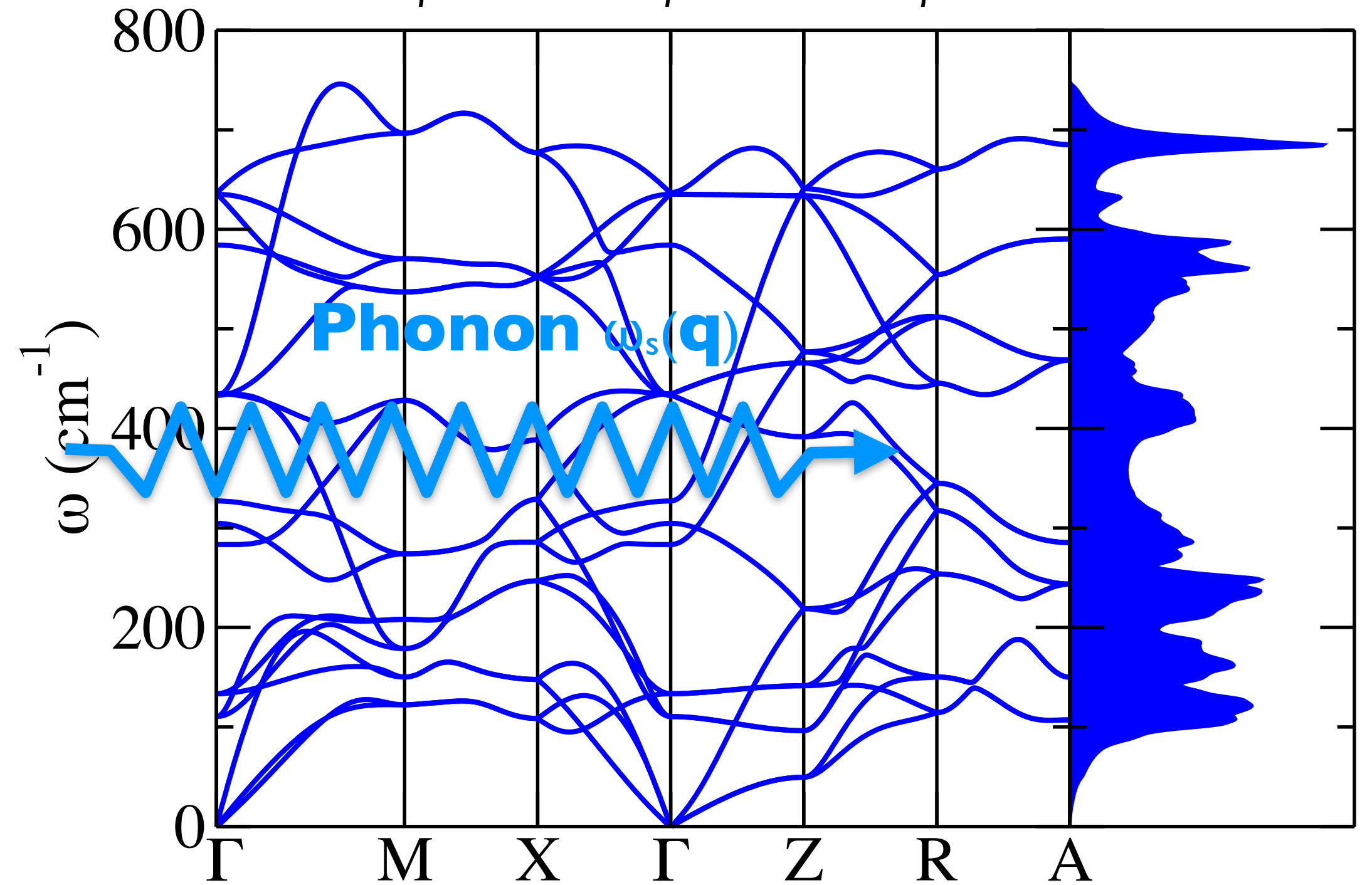
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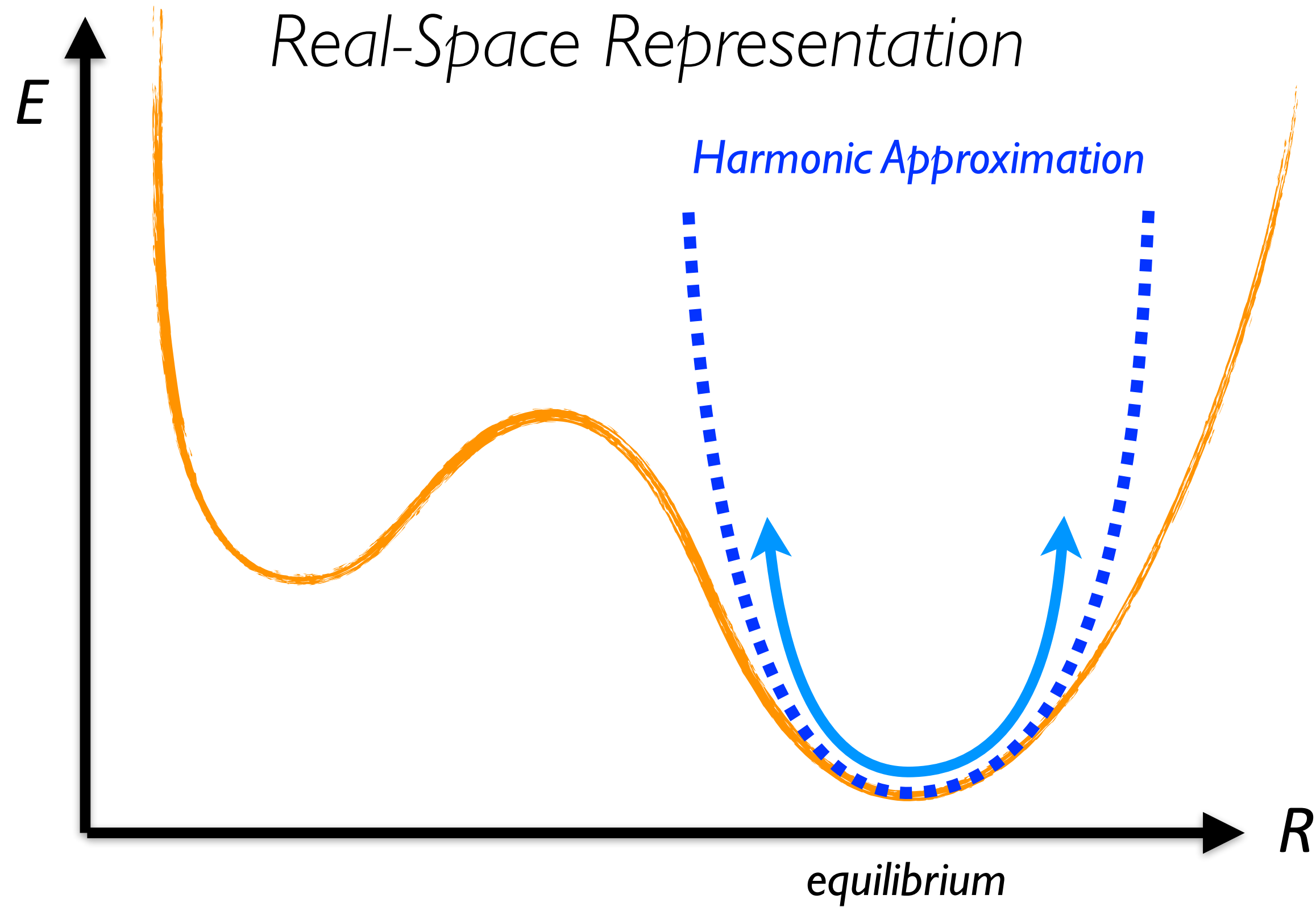
Reciprocal-Space Representation



Infinite Lifetimes

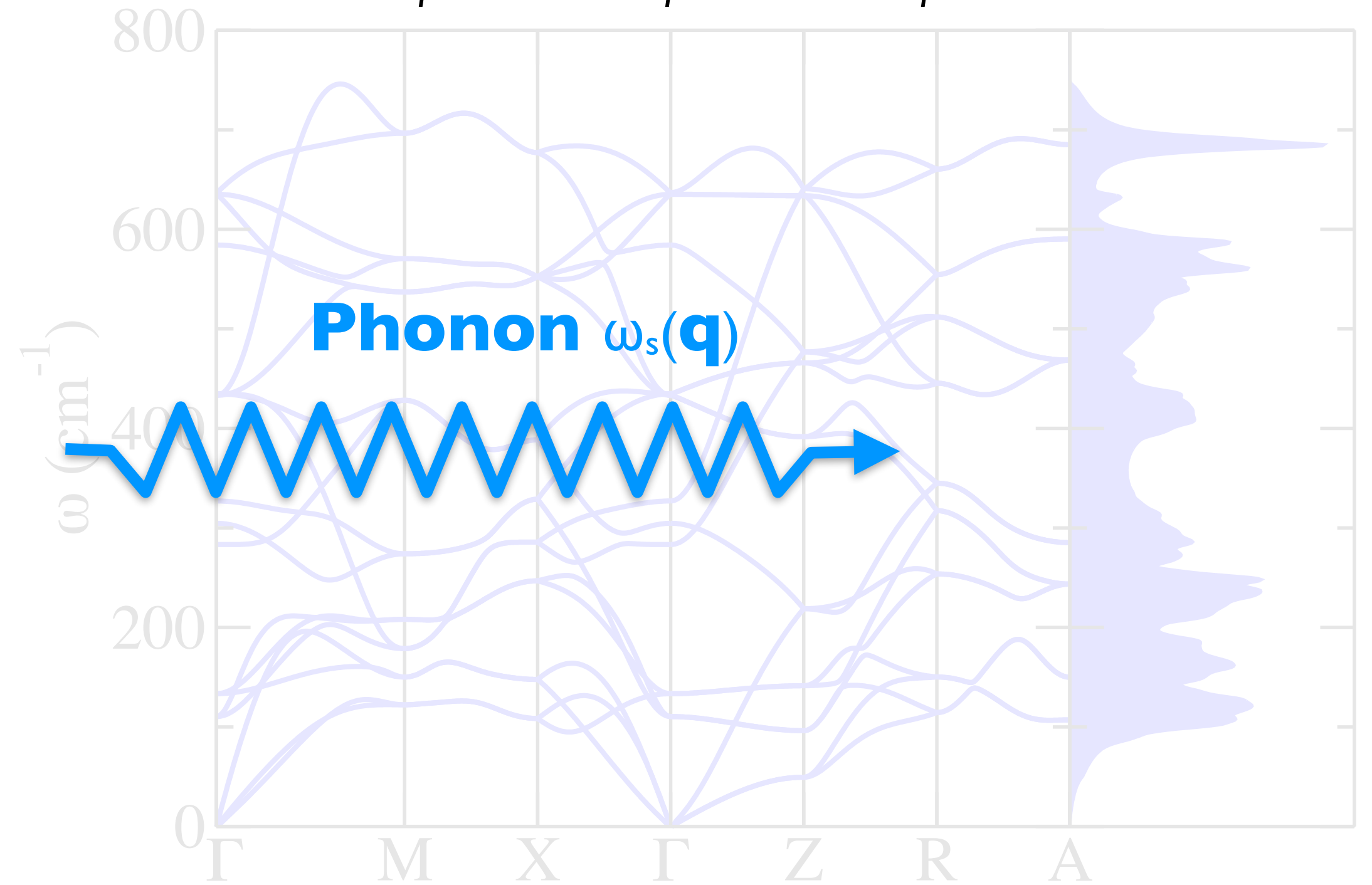
Heat Transport Theory 101

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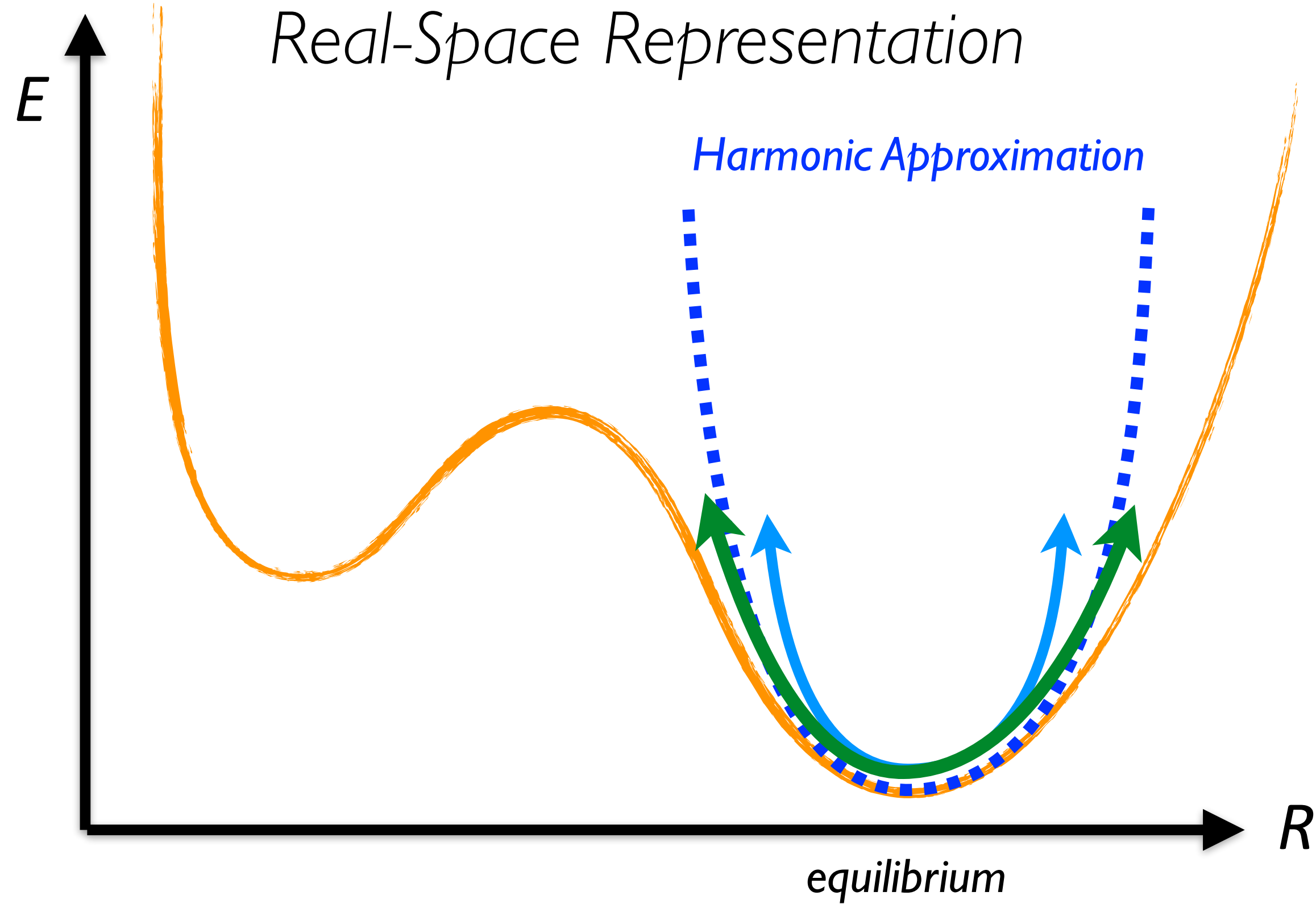


Infinite Lifetimes

Infinite Thermal Conductivity

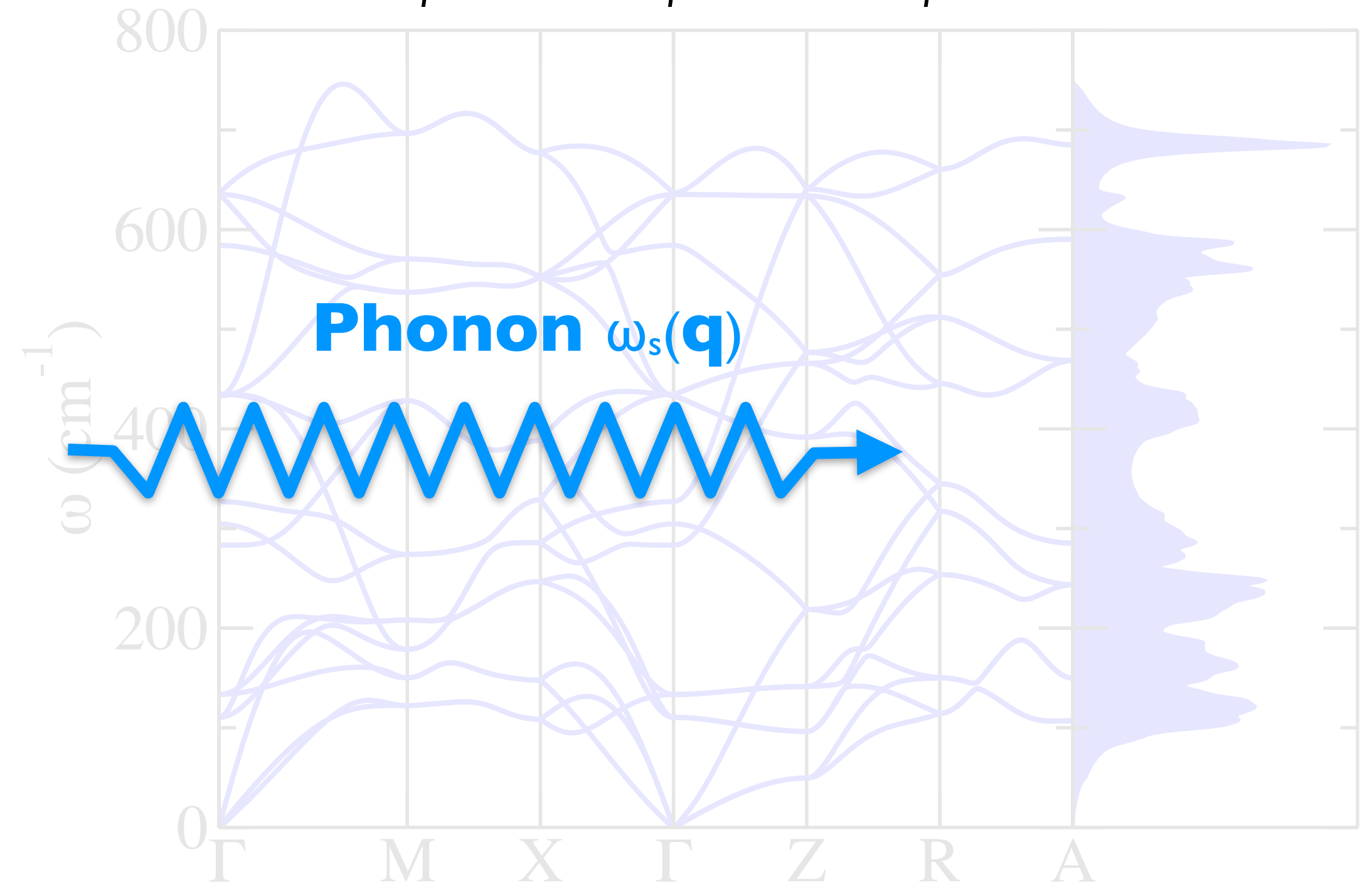
Heat Transport Theory 101

Real-Space Representation



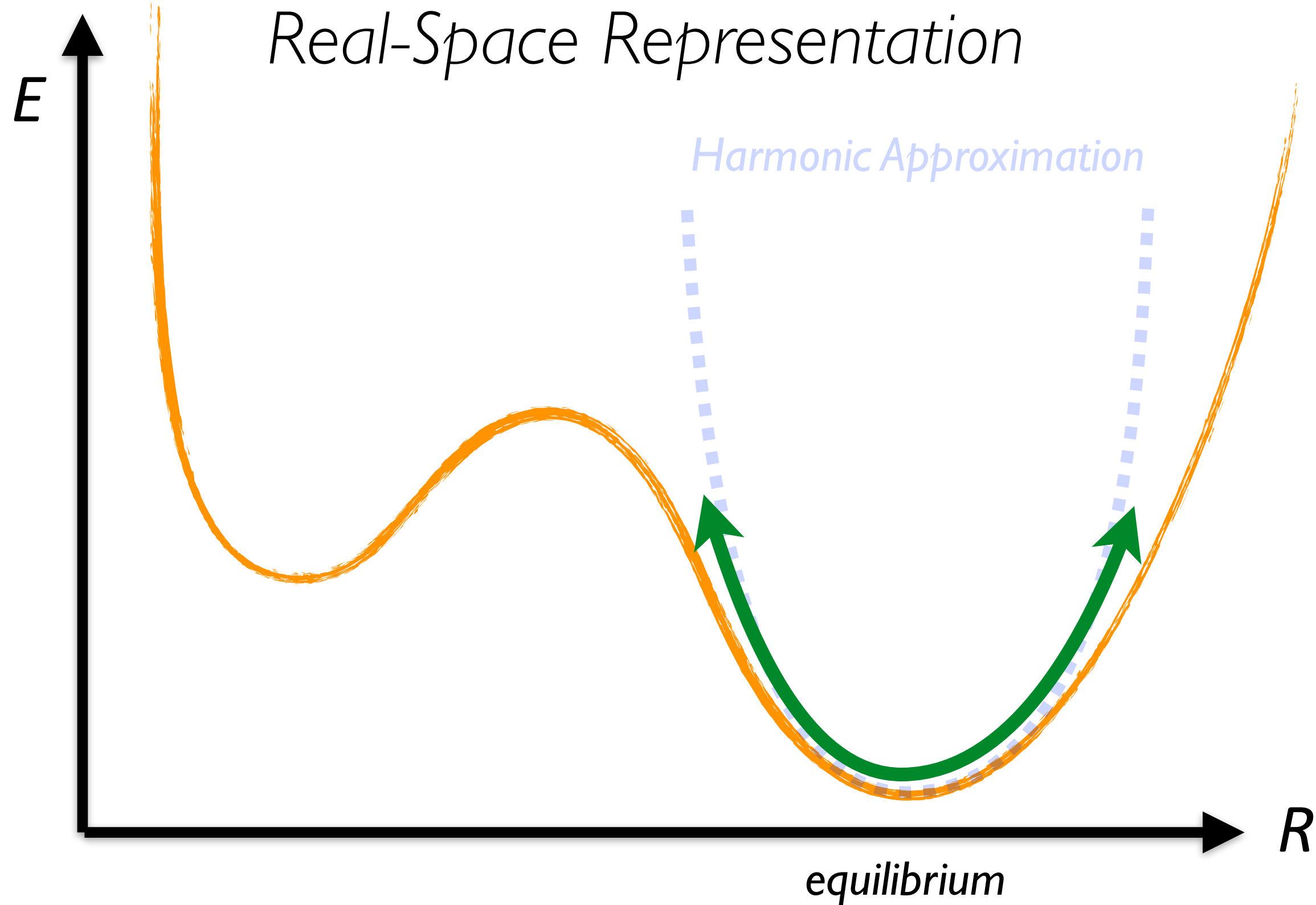
Anharmonic Dynamics

Reciprocal-Space Representation



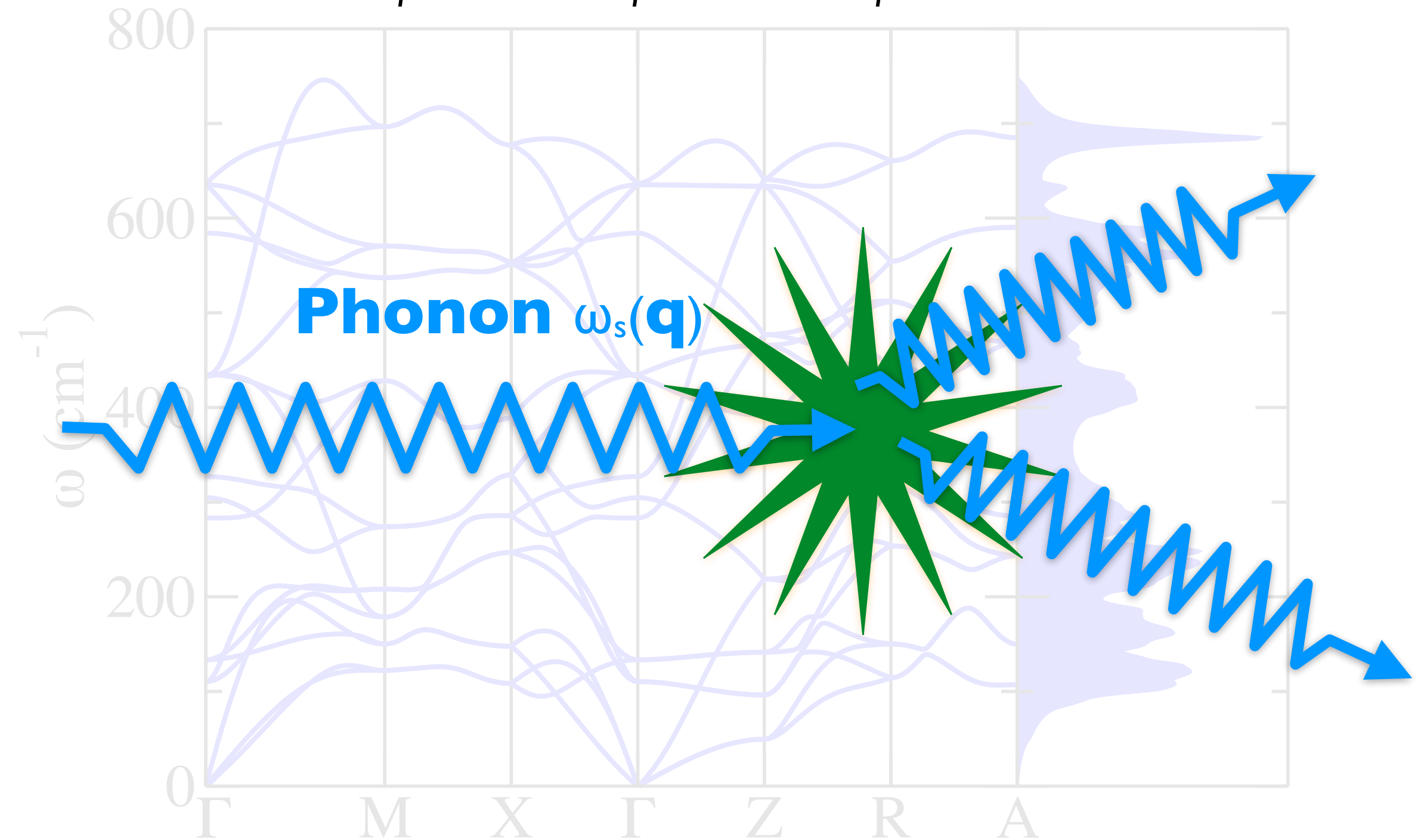
Heat Transport Theory 101

Real-Space Representation



Anharmonic Dynamics

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Phonon Scattering

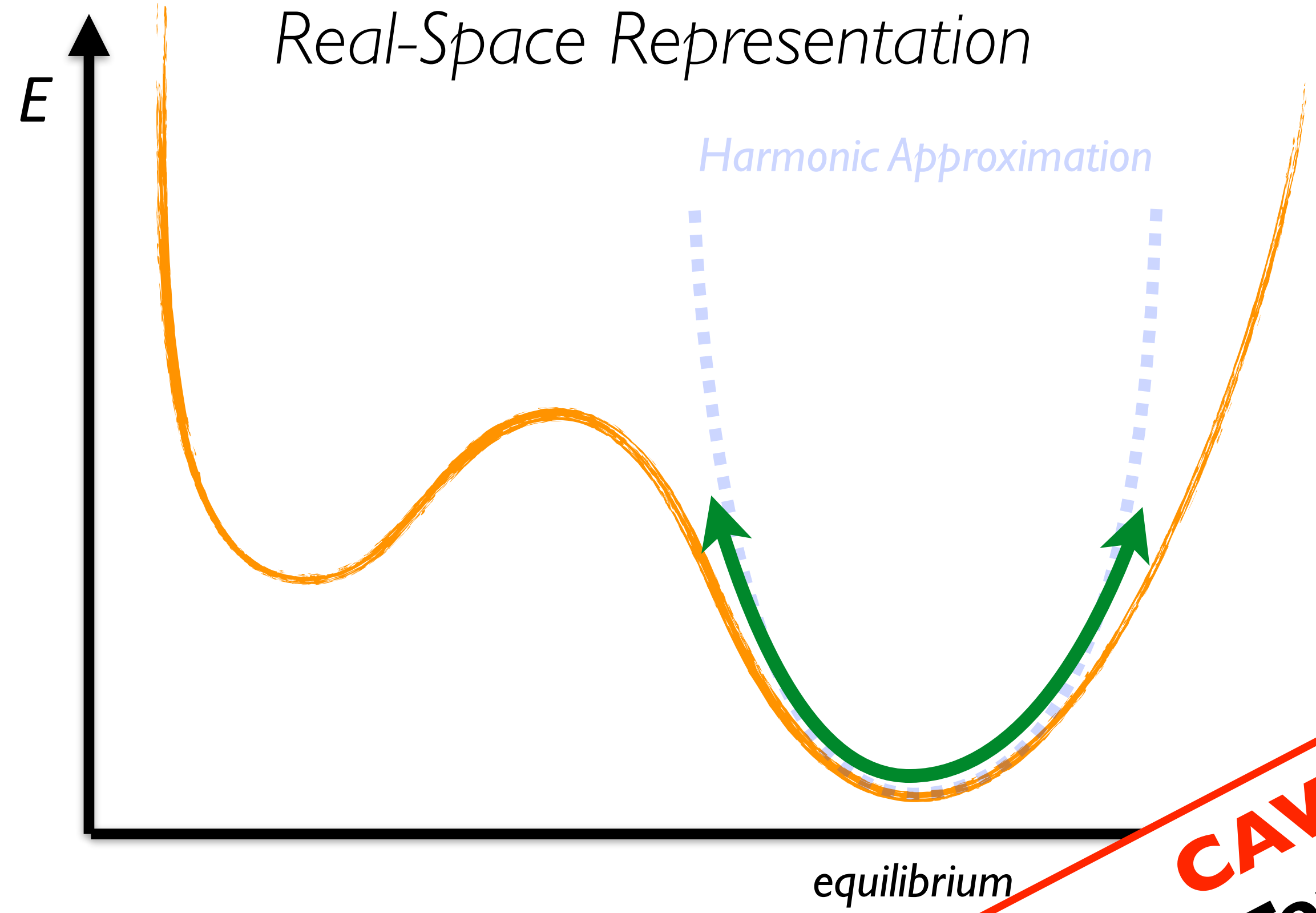
Perturbation Theory:

Scattering Cross Sections + Transport Equation

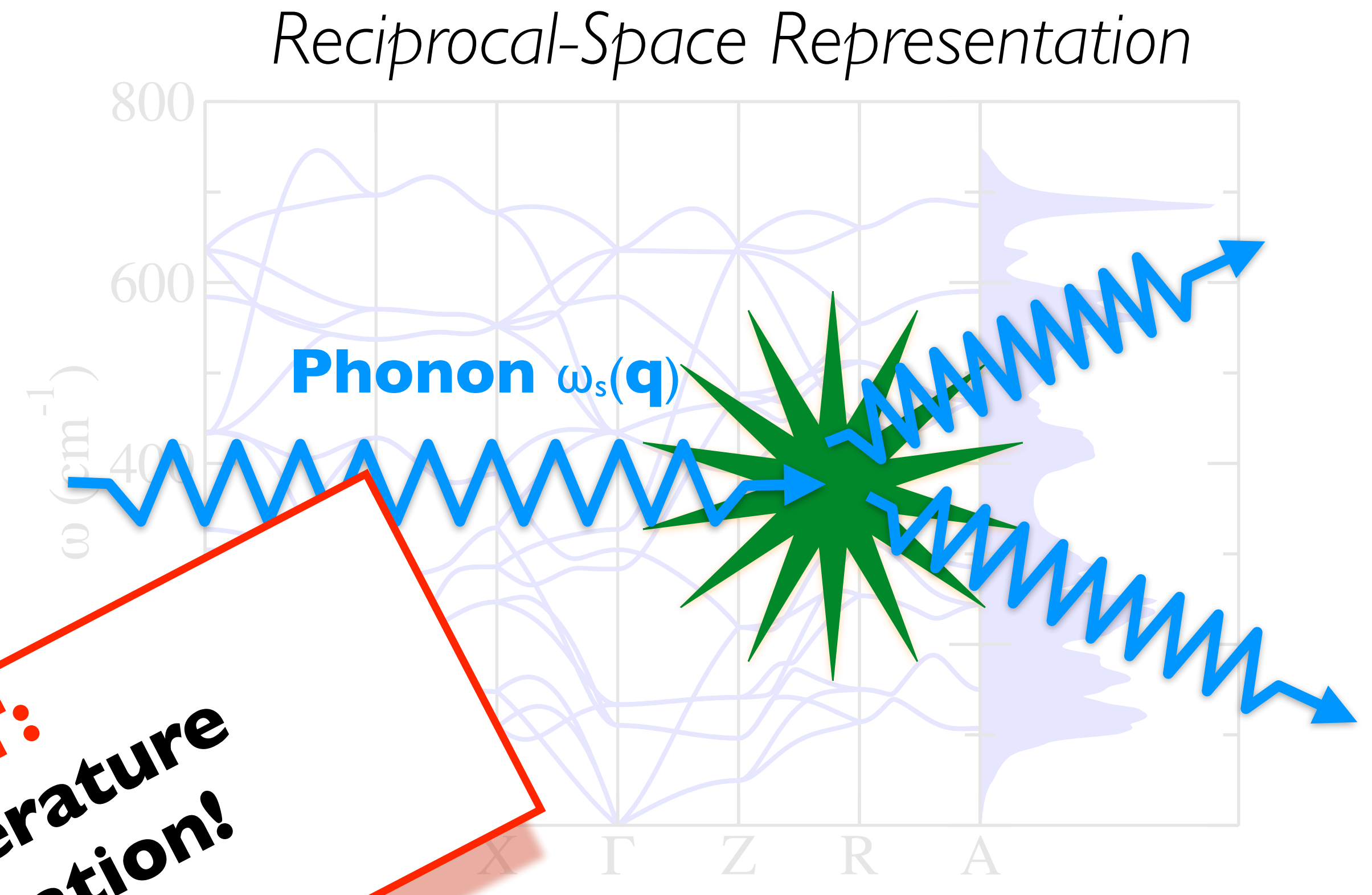
D. A. Broido et al., *Appl. Phys. Lett.* **91**, 231922 (2007).

Heat Transport Theory 101

Real-Space Representation



Reciprocal-Space Representation



Anharmonic D

CAVEAT:
Low Temperature
Approximation!
 $E_{\text{harm}} \gg E_{\text{anha}}$

Phonon Scattering

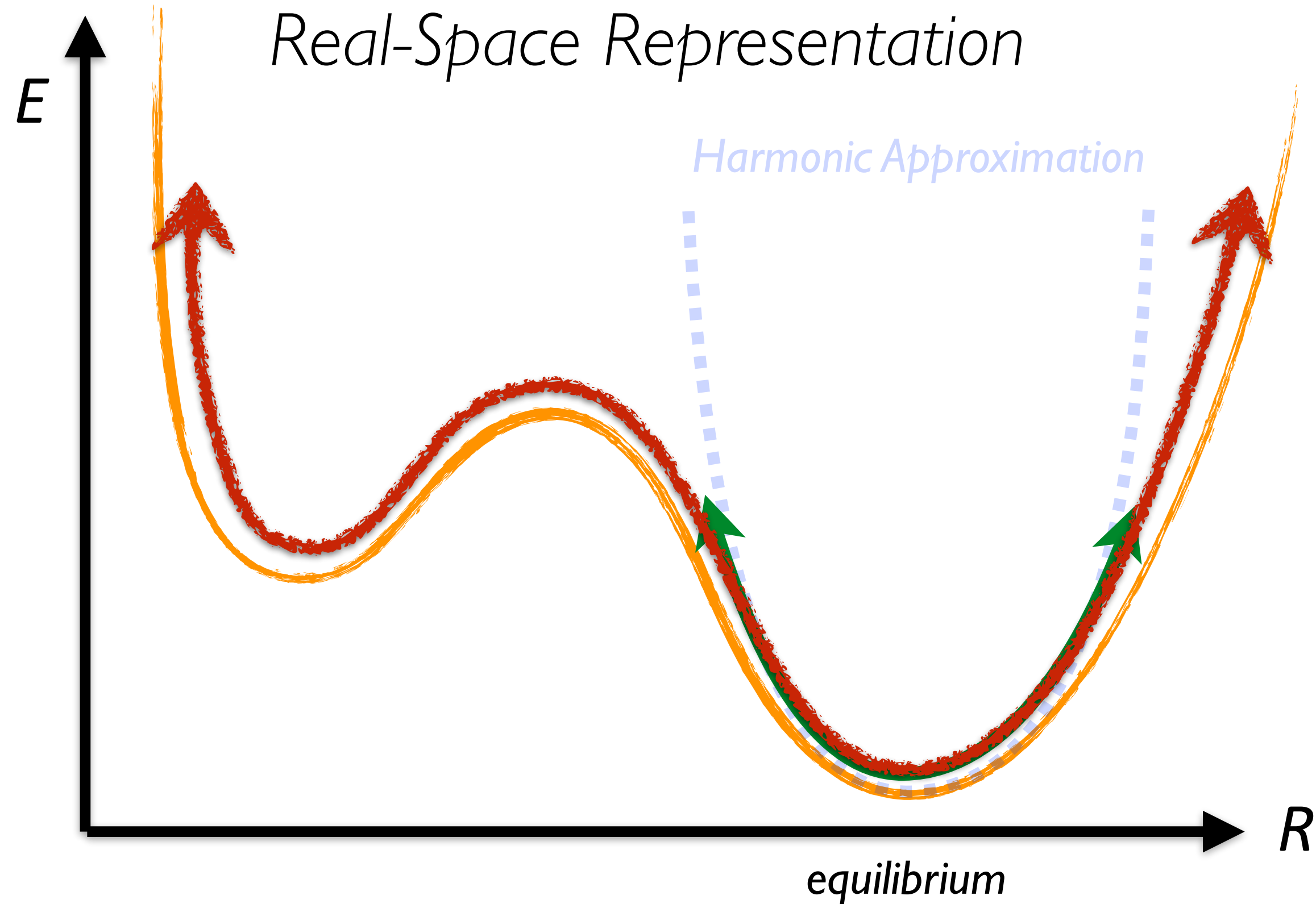
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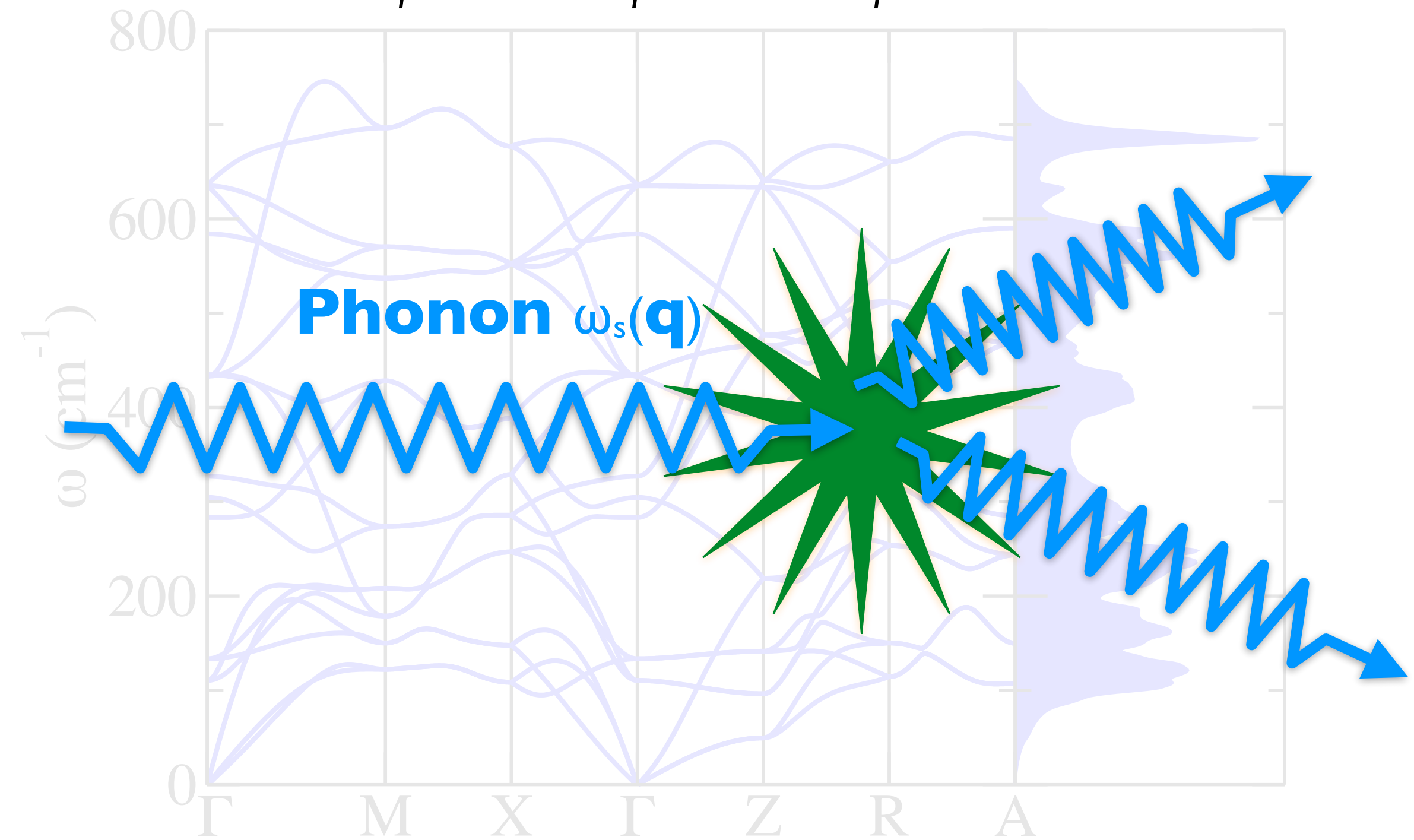
Heat Transport Theory 101

Real-Space Representation



Anharmonic Dynamics

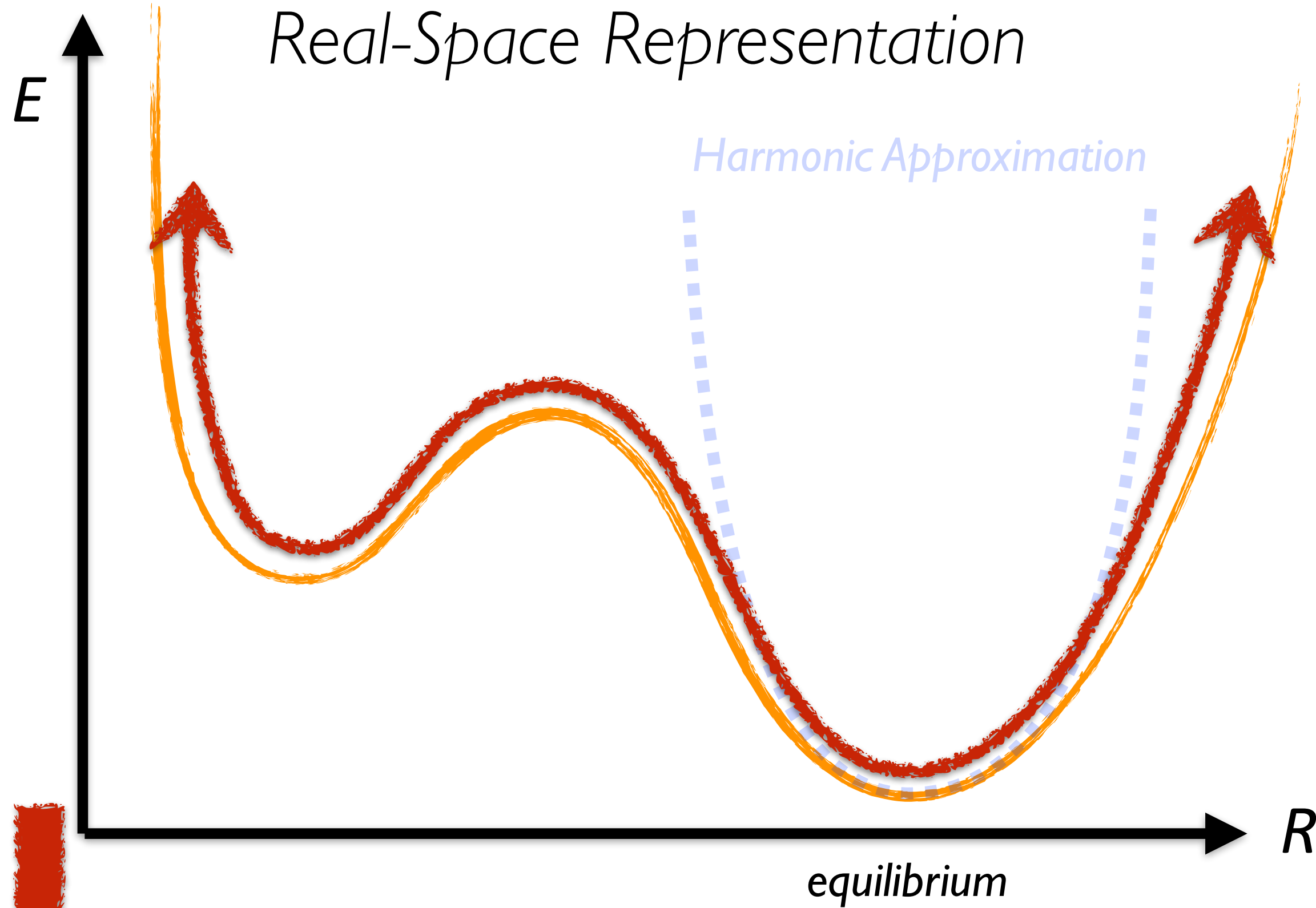
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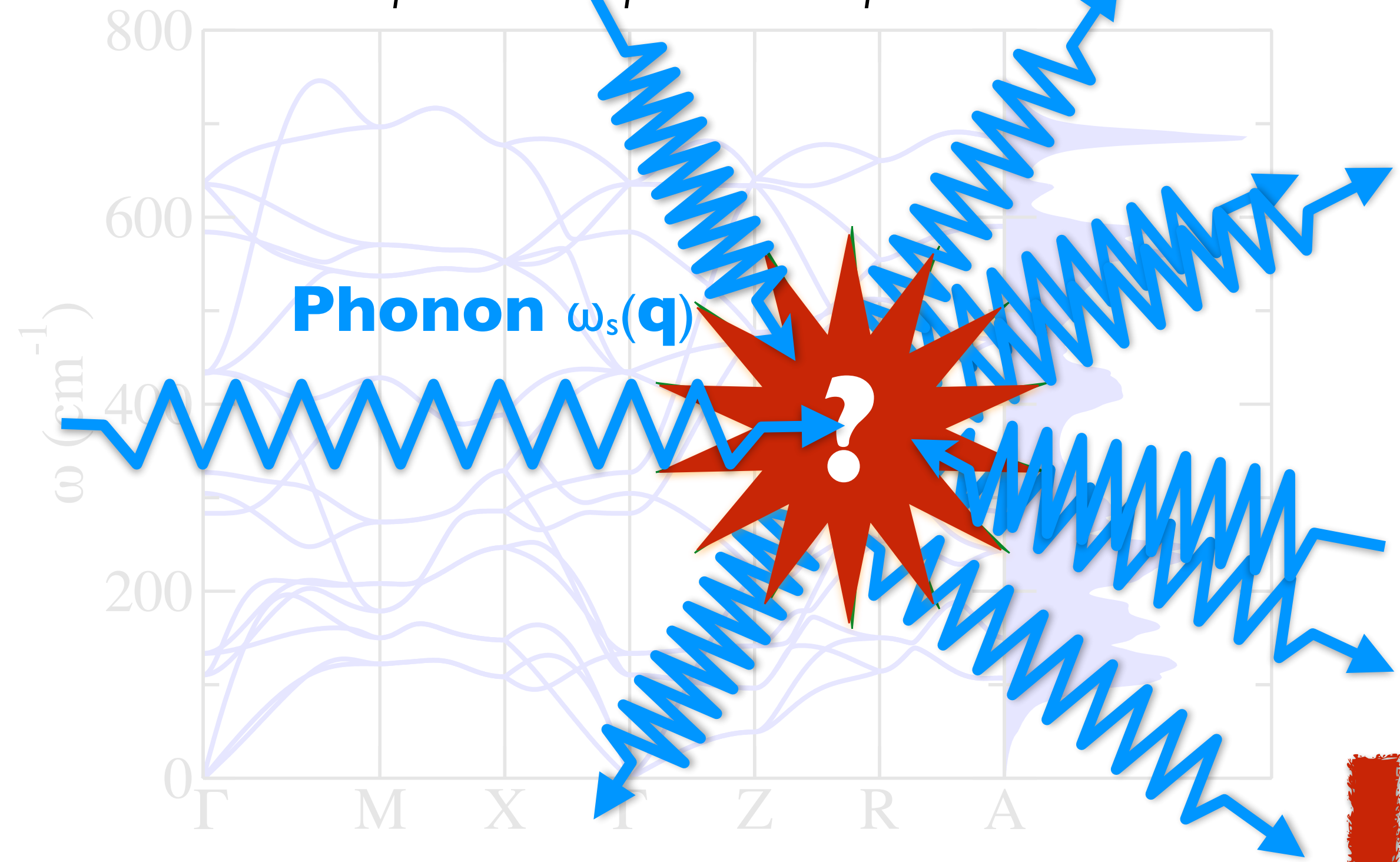
Heat Transport Theory 101

Real-Space Representation



Anharmonic Dynamics

Reciprocal-Space Representation



Phonon Scattering

Strong Anharmonic Effects beyond the Realm of Perturbation Theory:

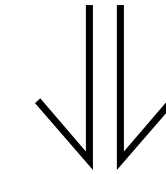
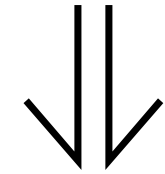
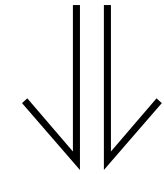
$$E_{\text{harm}} \ll E_{\text{anha}} !$$

GREEN-KUBO METHOD

R. Kubo, M. Yokota, and S. Nakajima, *J. Phys. Soc. Japan* **12**, 1203 (1957).

Fluctuation-Dissipation Theorem

Simulations of the **thermodynamic equilibrium**



Information about **non-equilibrium processes**

$$\kappa \sim \int_0^{\infty} d\tau \langle \mathbf{J}(0) \mathbf{J}(\tau) \rangle_{eq}$$

The **thermal conductivity** is related to the **autocorrelation function** of the **heat flux**.

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Green-Kubo method...

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...accounts for **anharmonic** effects **to all orders**

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Green-Kubo simulations in *FHI-aims*:



- Conductive **heat flux** evaluation using **virials**.

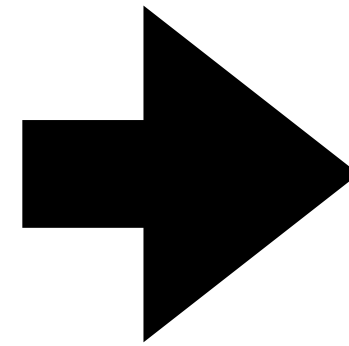
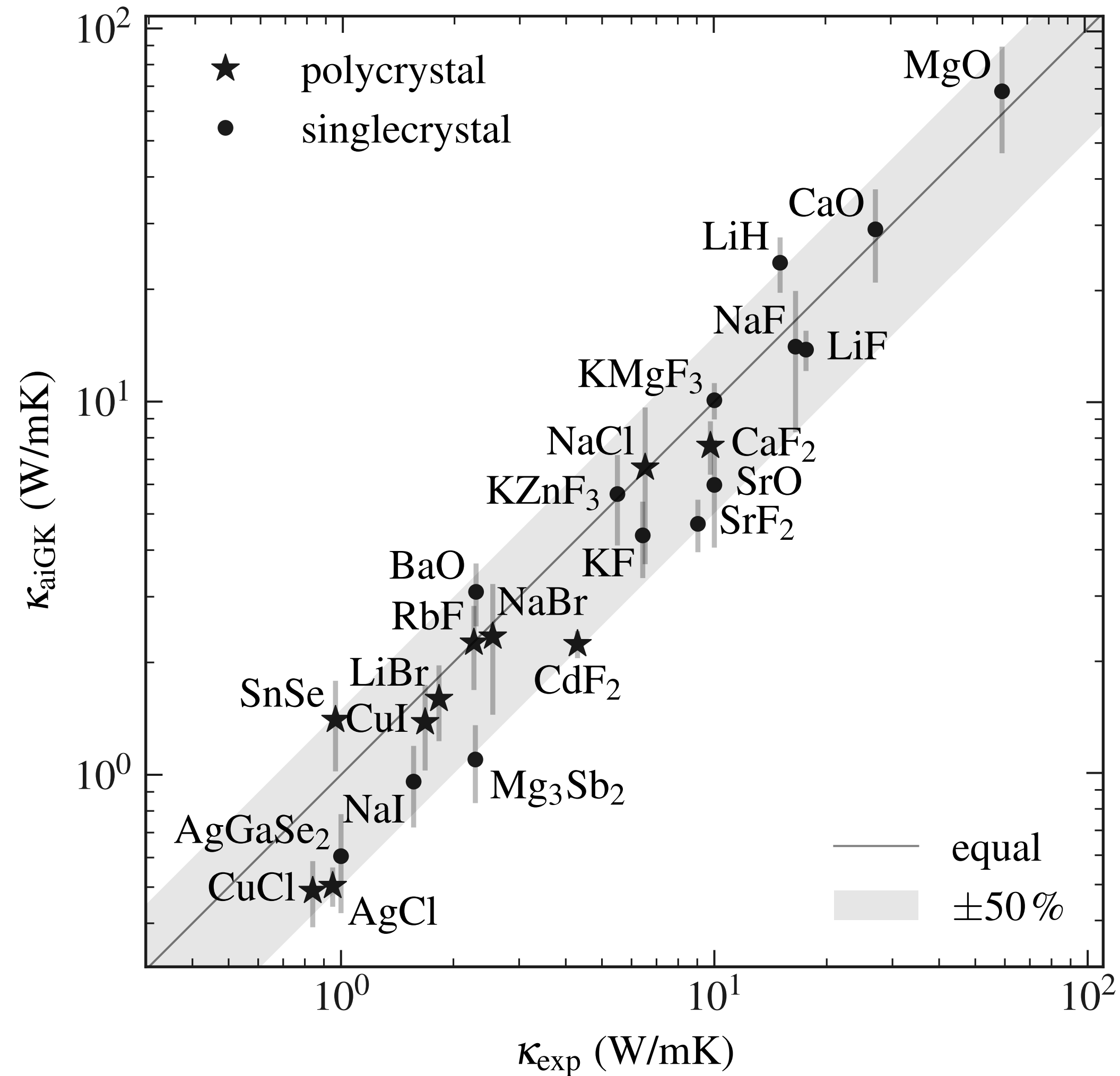
C. Carbogno, R. Ramprasad, and M. Scheffler, *Phys. Rev. Lett.* **118**, 175901 (2017).

- Numerical **post-processing** via *FHI-vibes*

F. Knoop, M. Scheffler, and C. Carbogno, *Phys. Rev. B* **107**, 224304 (2023).

Ab initio Green-Kubo Validation

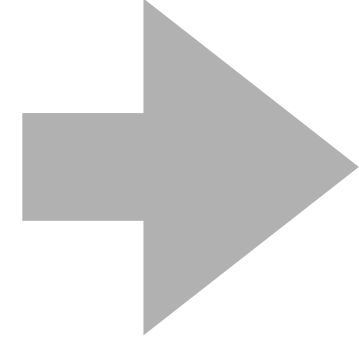
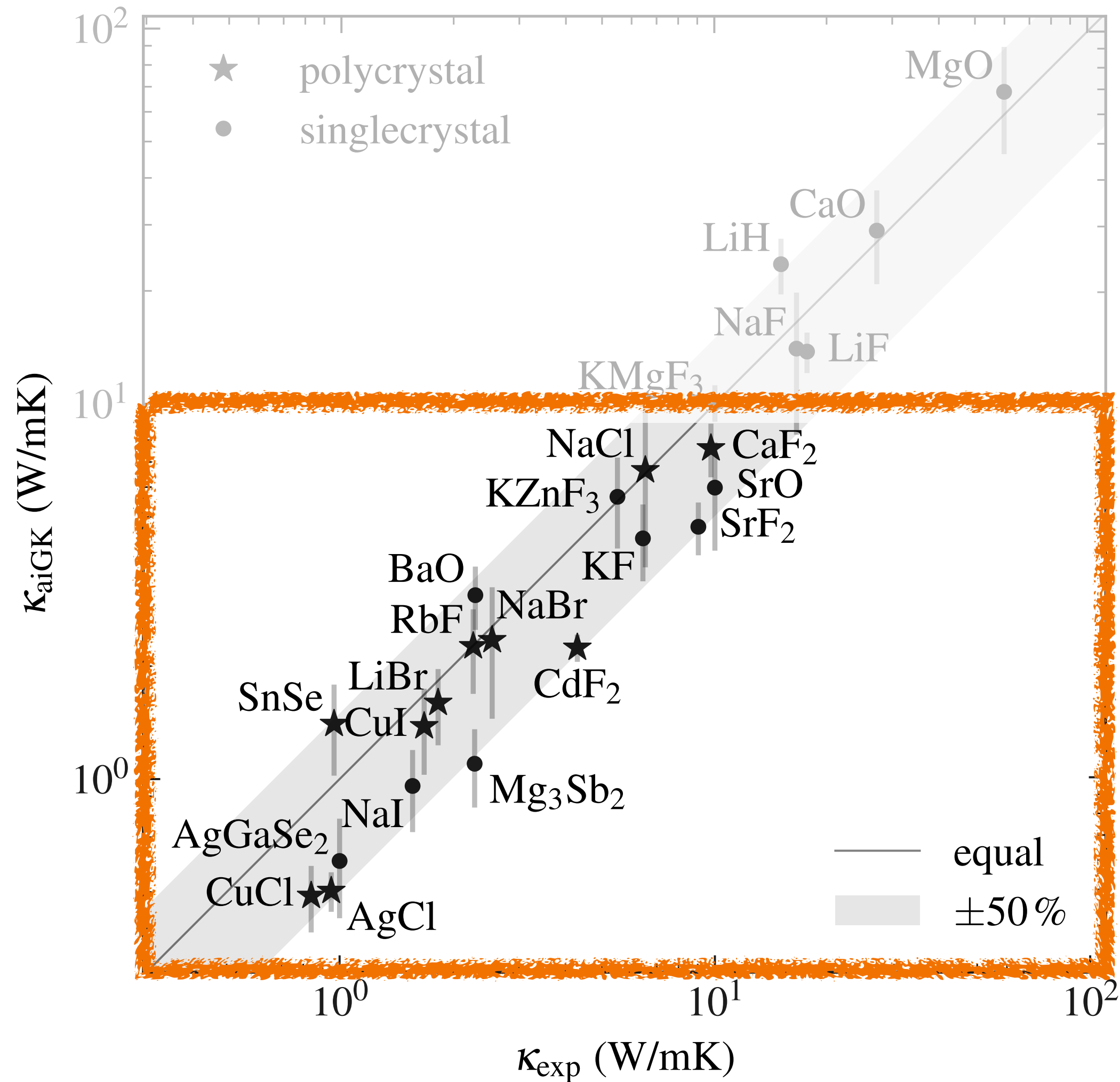
F. Knoop, T. A. R. Purcell, M. Scheffler, and C. Carbogno, *Phys. Rev. Lett.* **130**, 236301 (2023).



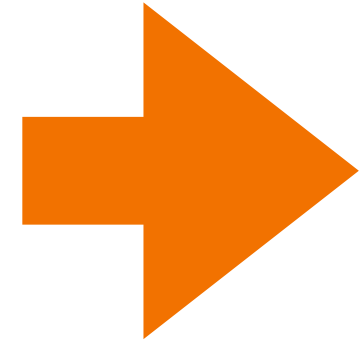
**Mean Absolute Percentage Error
MAPE ~ 25%**

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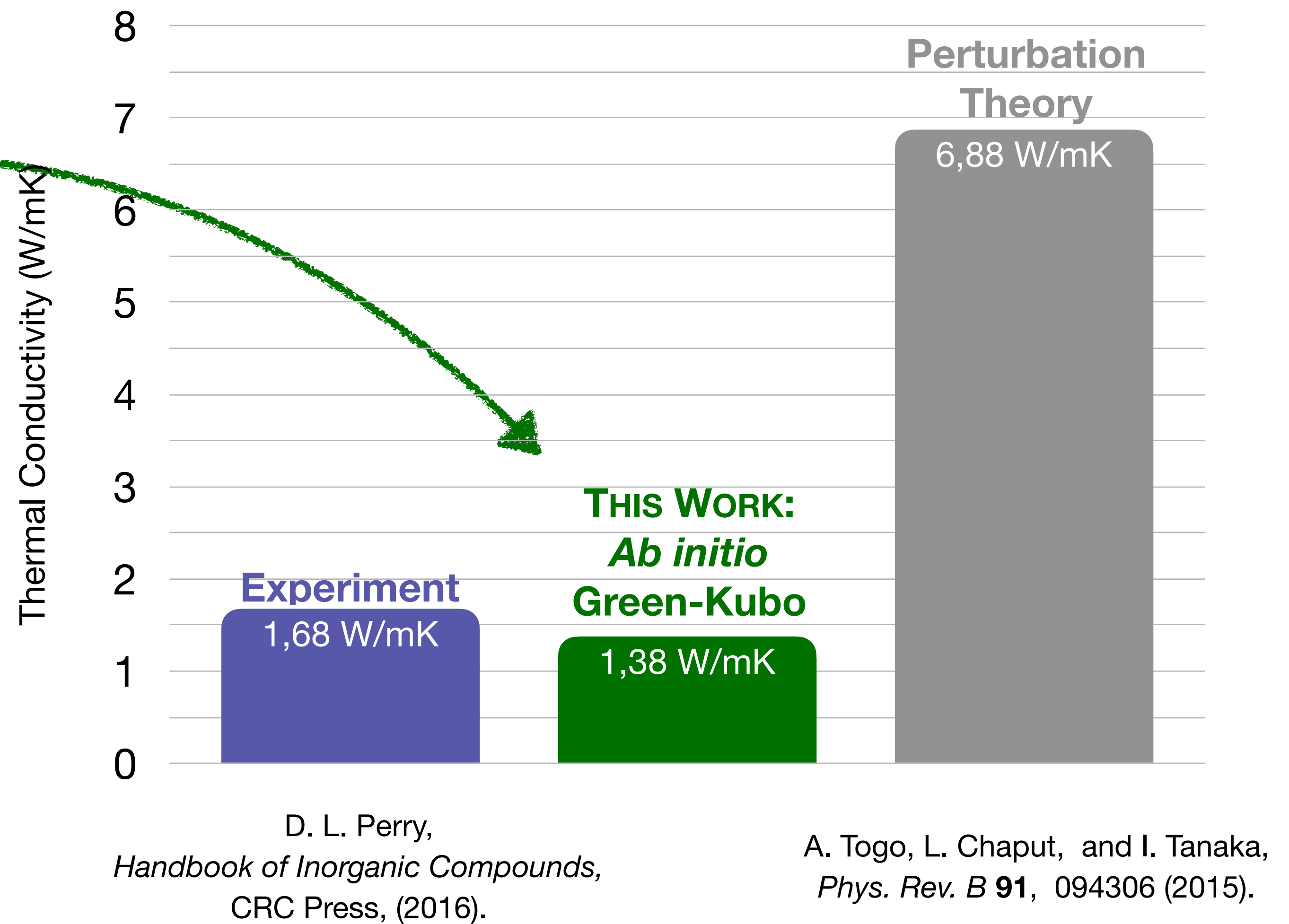
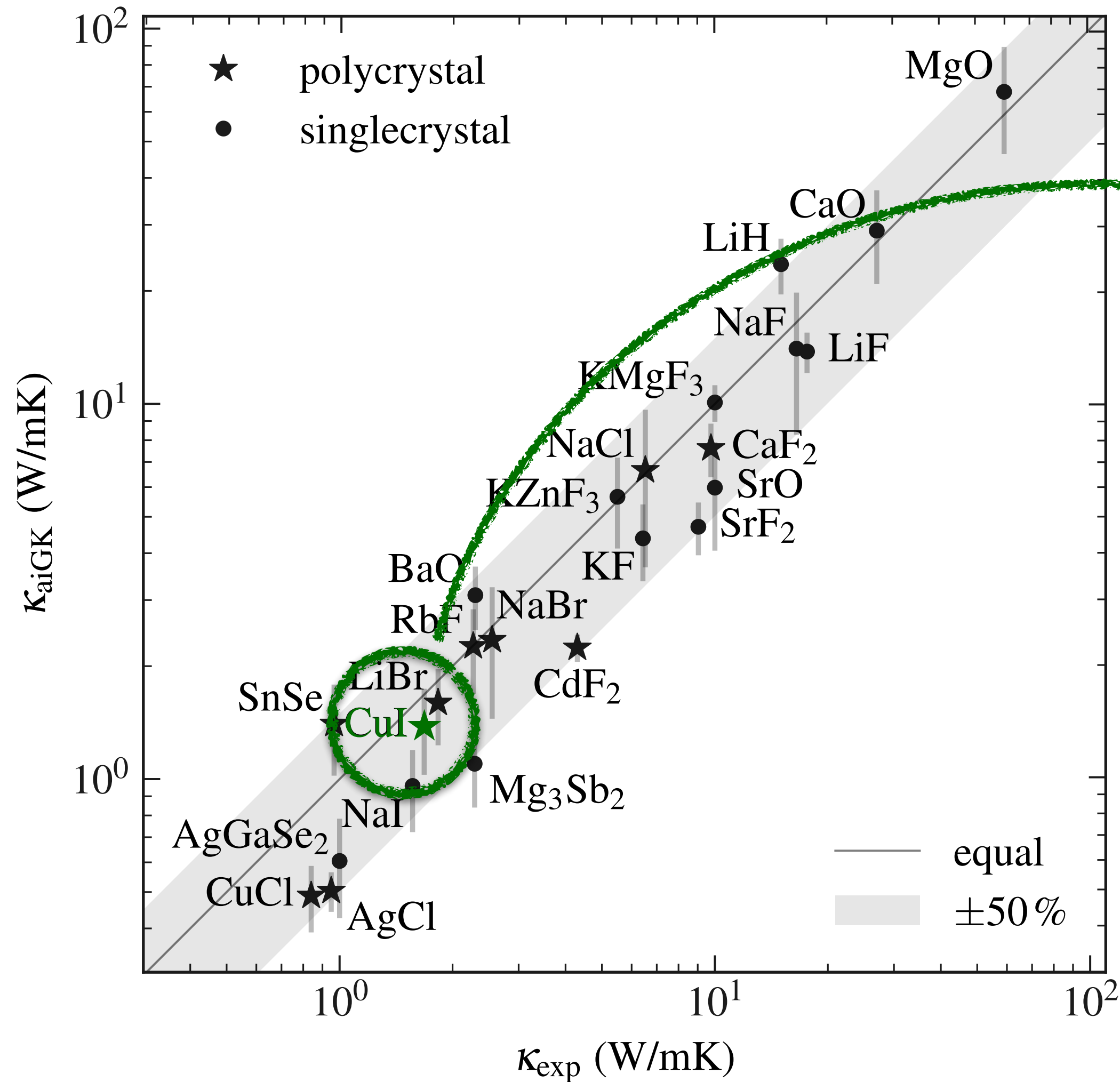
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Mean Absolute Error
MAE ~ 1 W/mK

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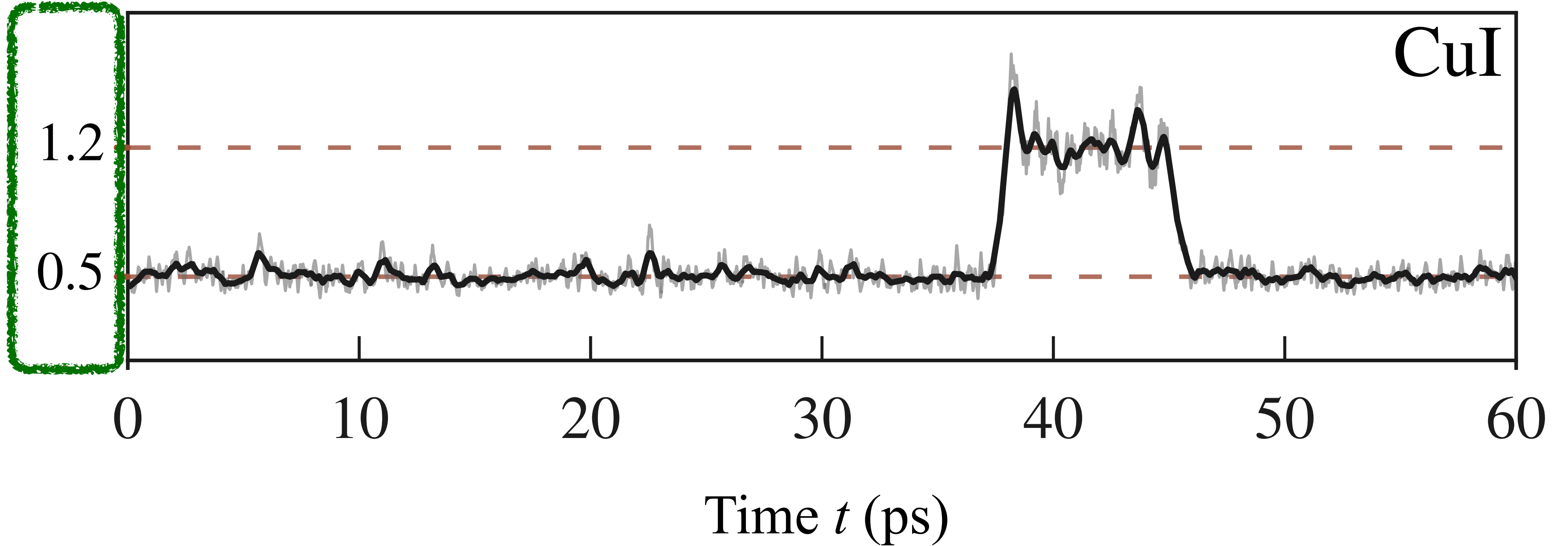
Anharmonicity in Copper Iodide

F. Knoop, T. A. R. Purcell, M. Scheffler, and C. Carbogno, *Phys. Rev. Lett.* **130**, 236301 (2023).

Anharmonicity

Metric σ^A

F. Knoop *et al.*,
Phys. Rev. Mater. **4**,
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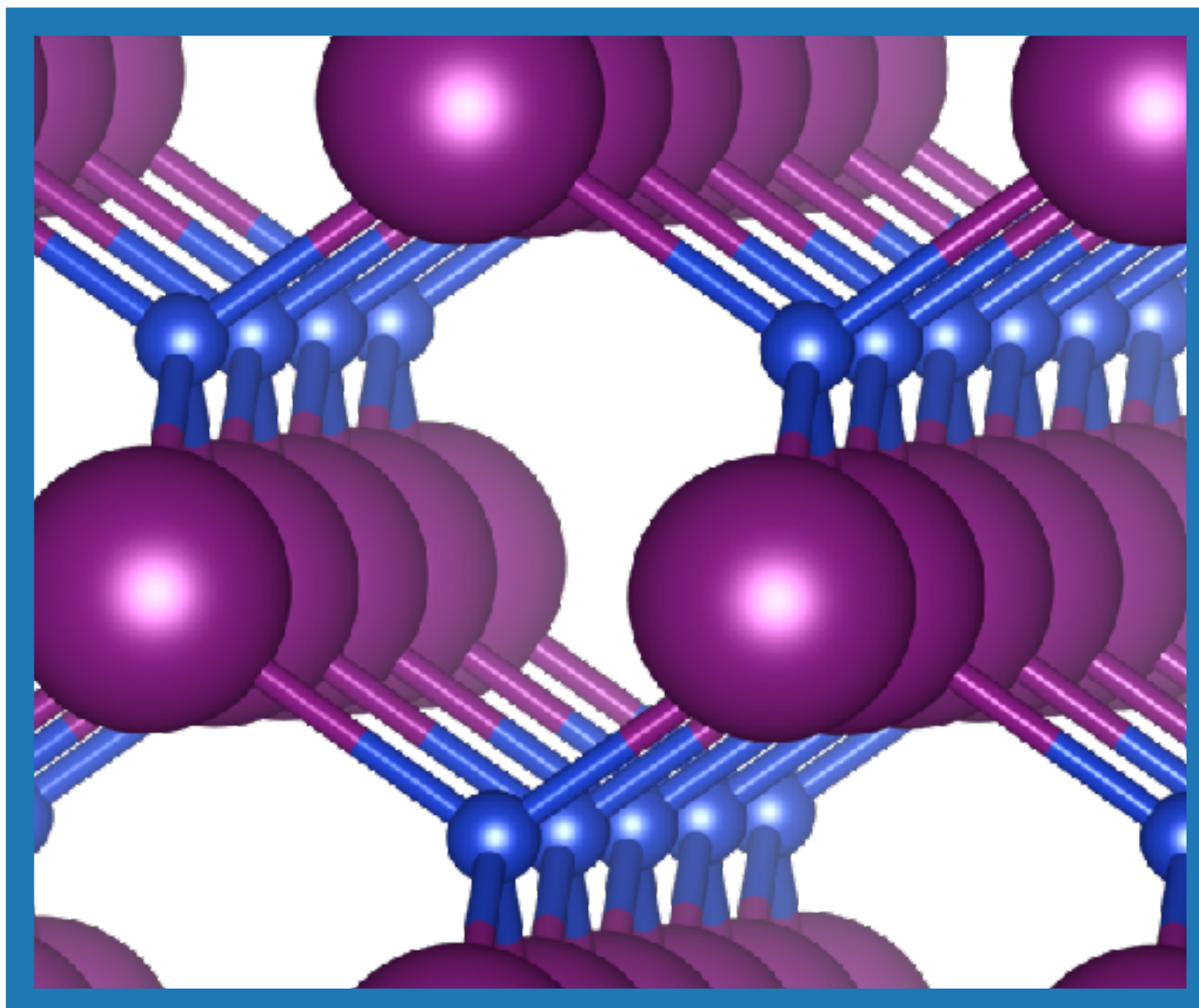
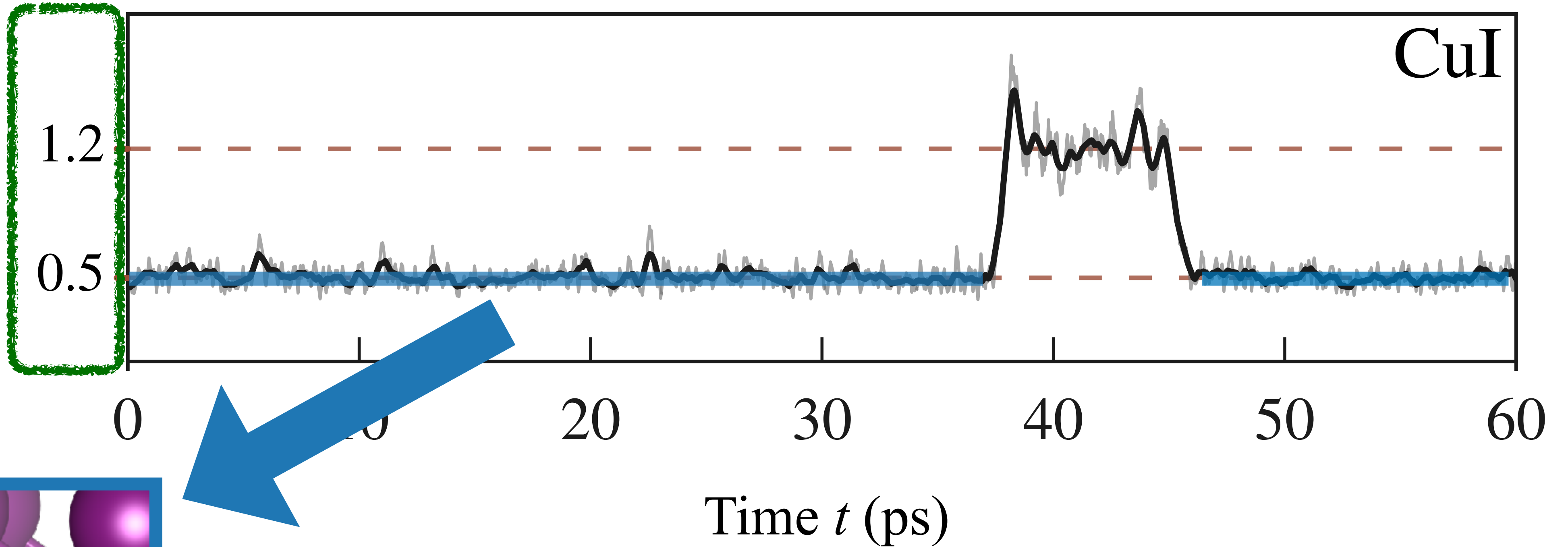
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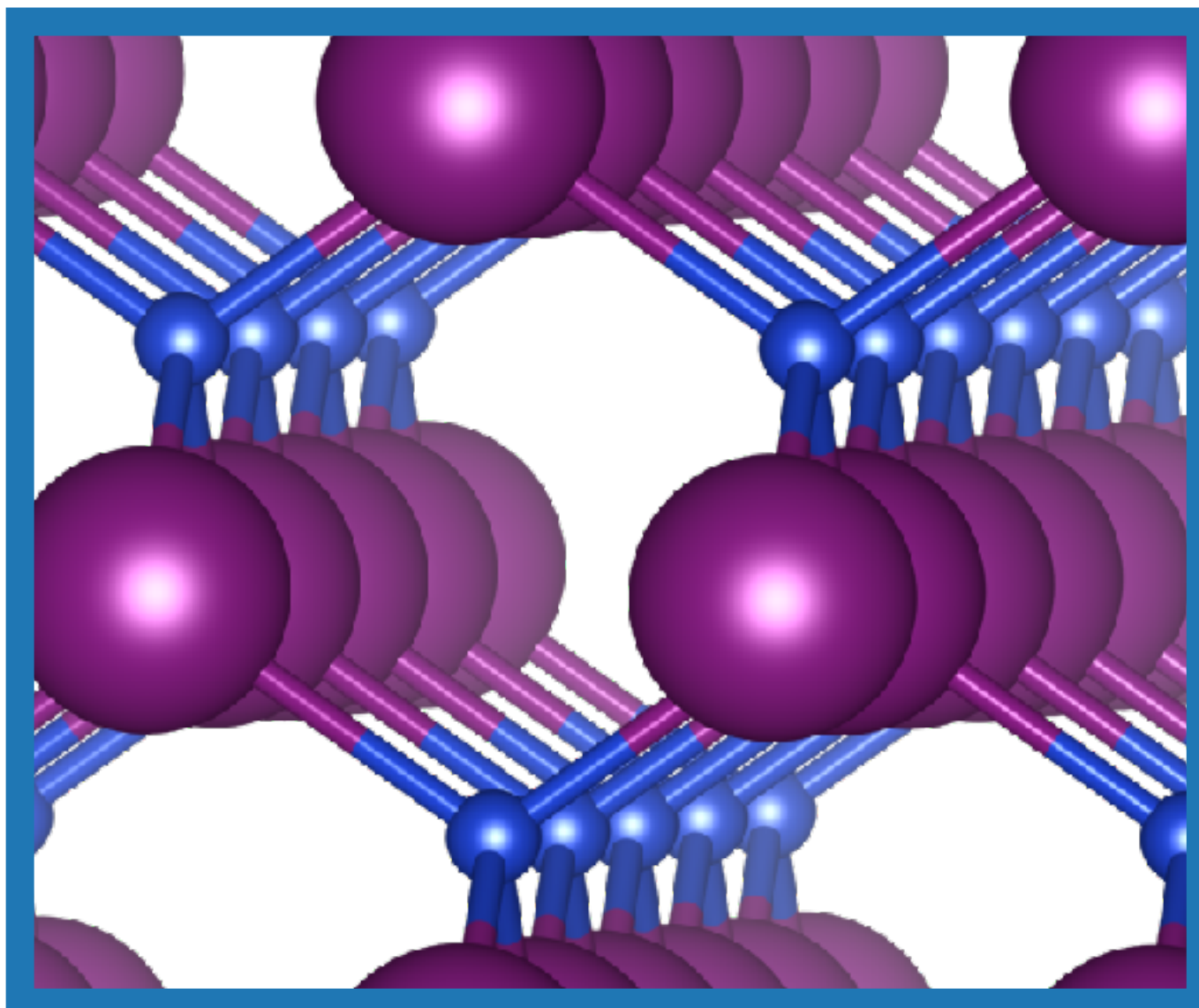
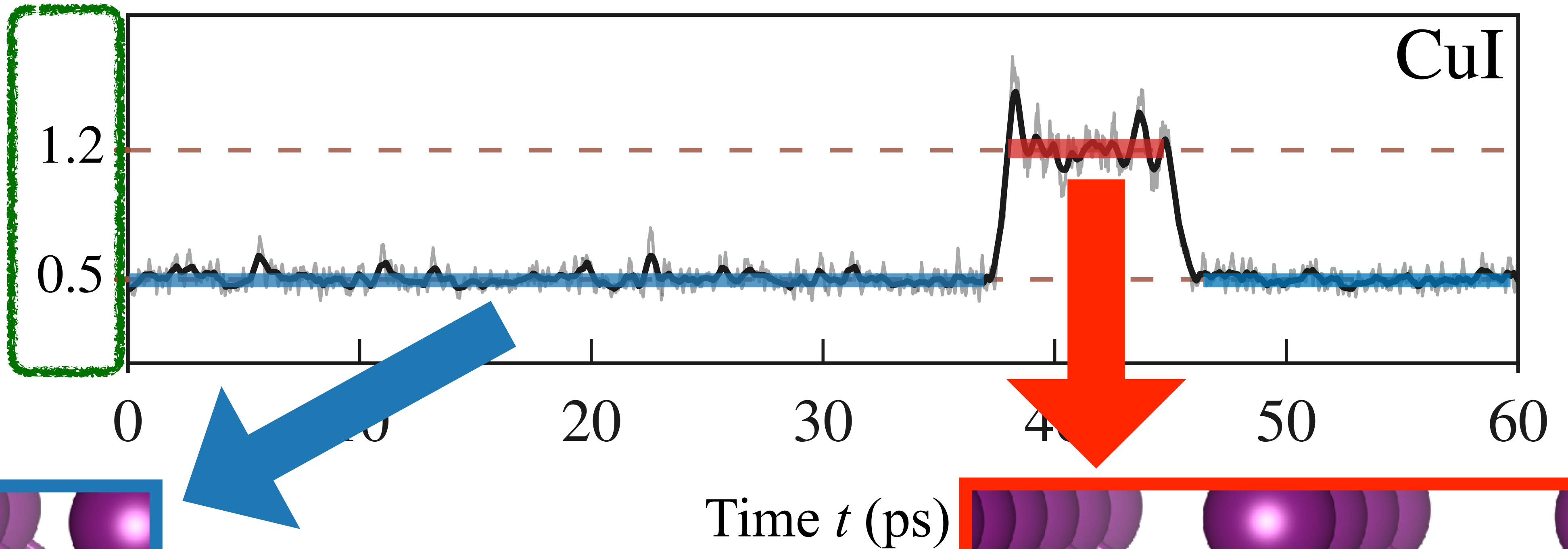
Time-averaged Structure:
Pristine *Zincblende*
Structure of Copper Iodide

Anharmonicity in Copper Iodide

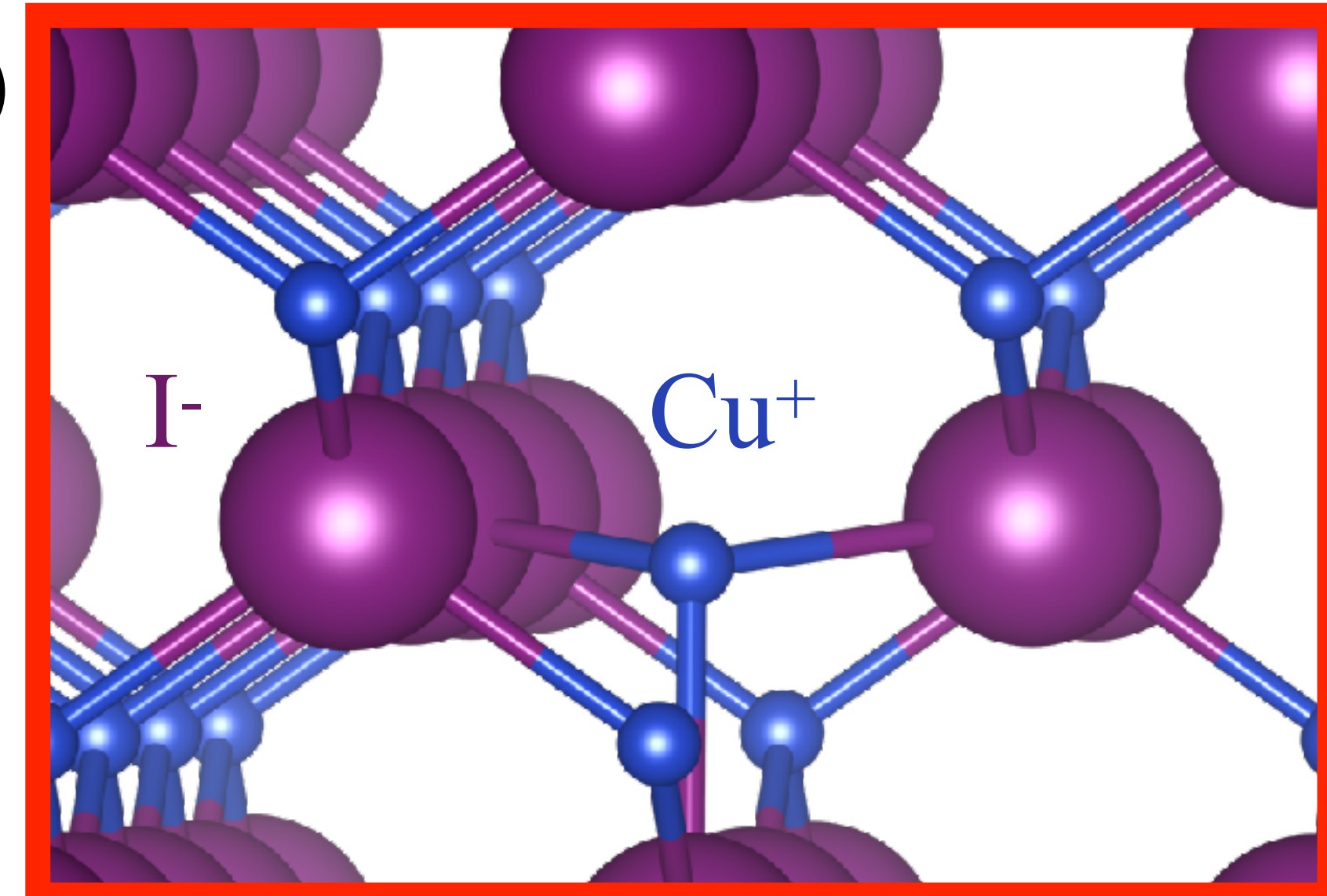
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Anharmonicity
Metric σ^A

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Phys. Rev. Mater. **4**,
083809 (2020).



Time-averaged Structure:
A metastable Copper
Self-Interstitial has formed!

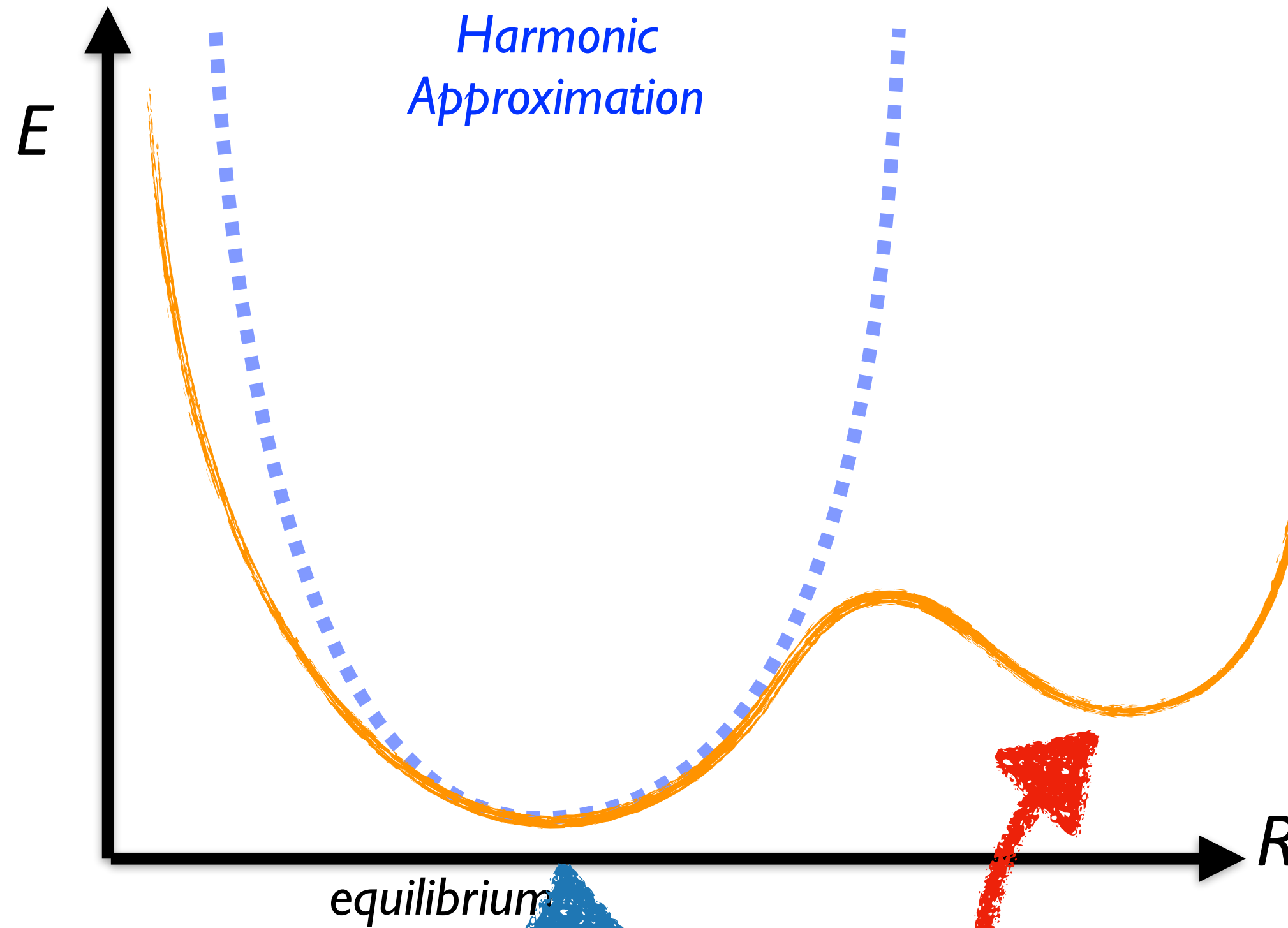


Anharmonicity in Copper Iodide

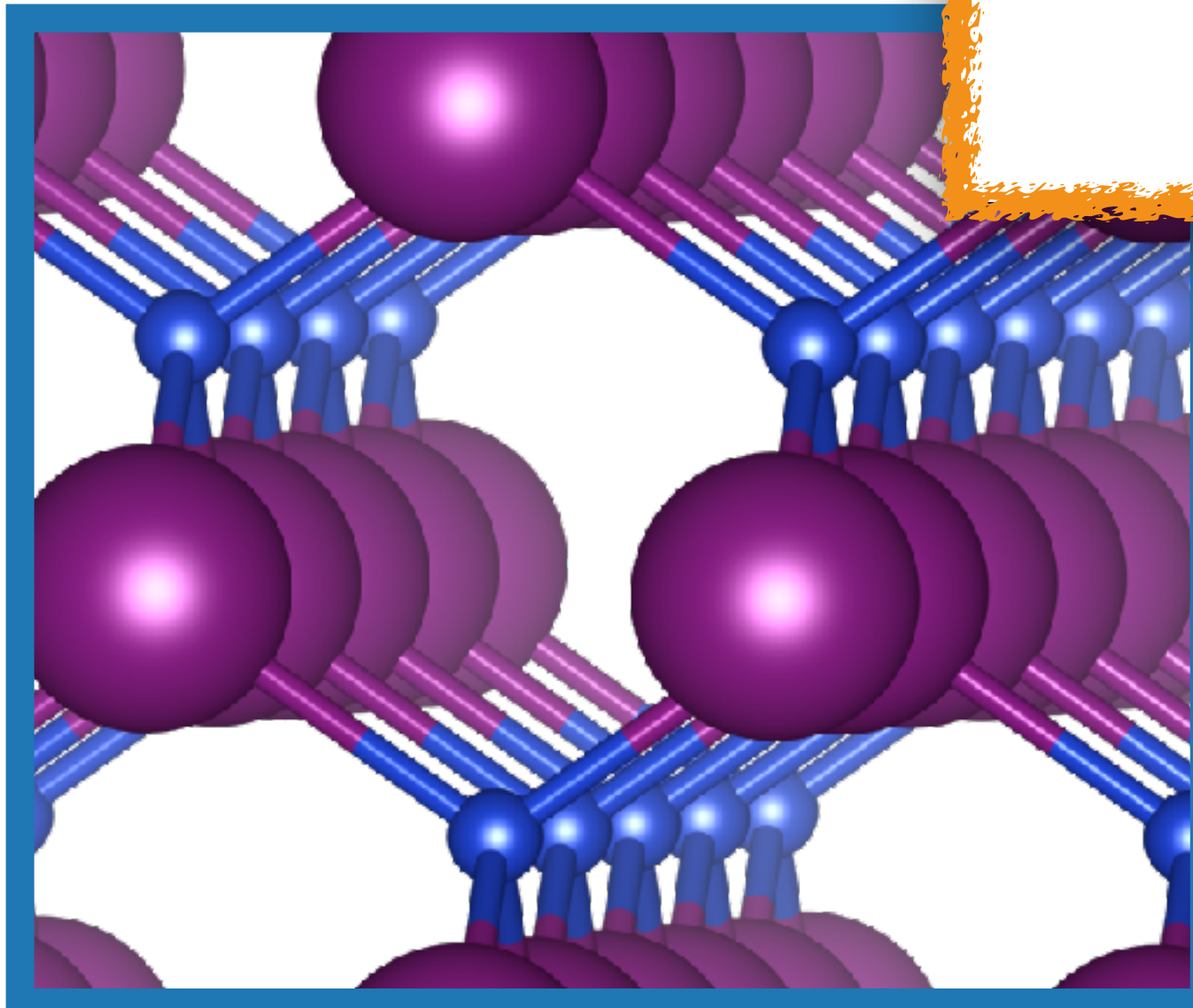
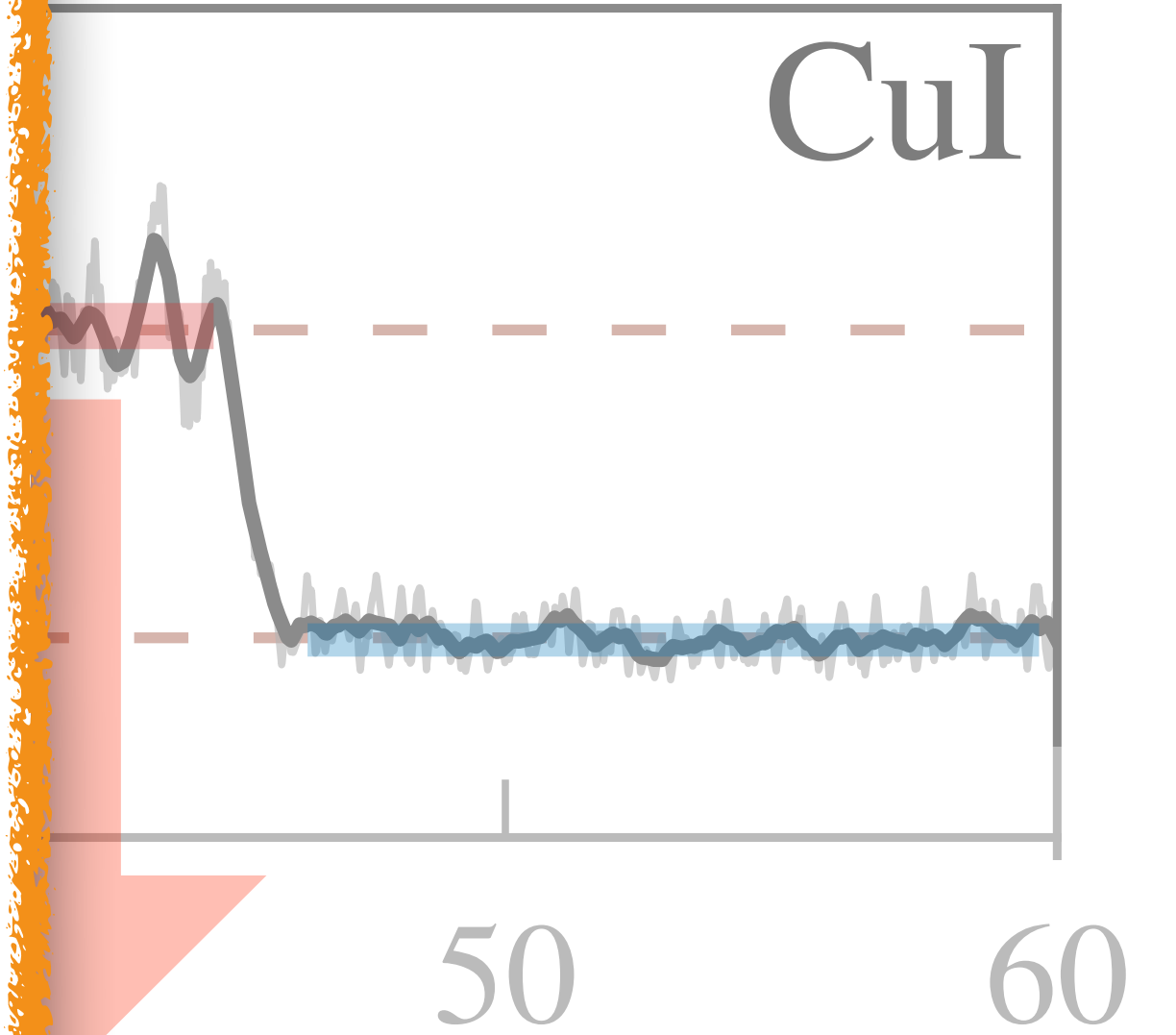
F. Knoop

Anharmonicity Metric σ^A

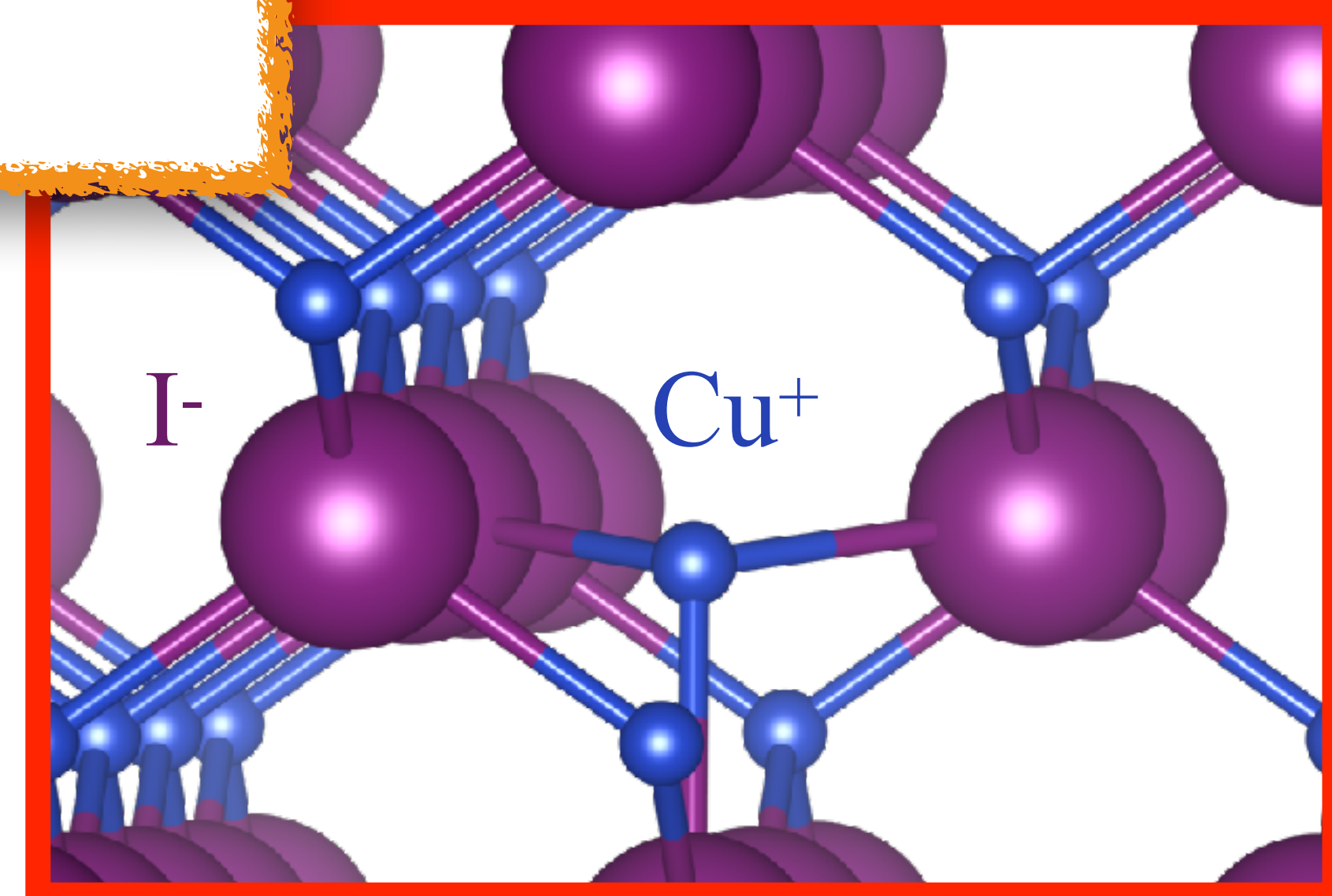
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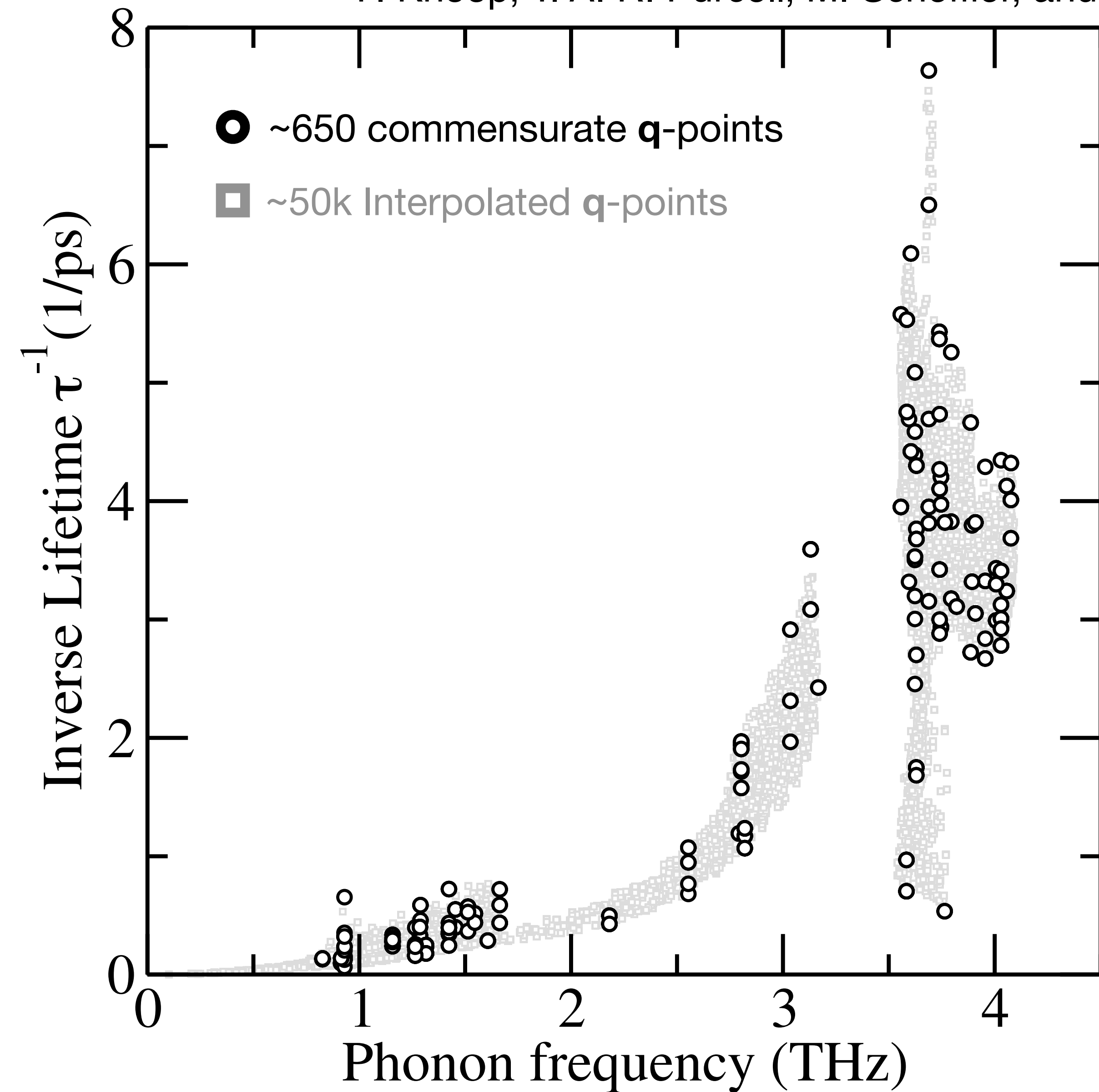
Time-averaged Structure:
Pristine Znblende
Structure of Copper Iodide



Time-averaged Structure:
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Self-Interstitial has formed!

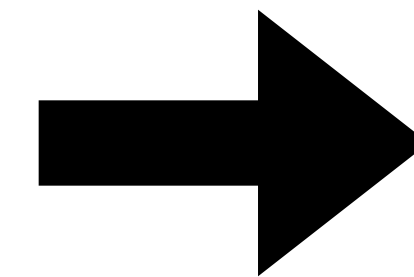
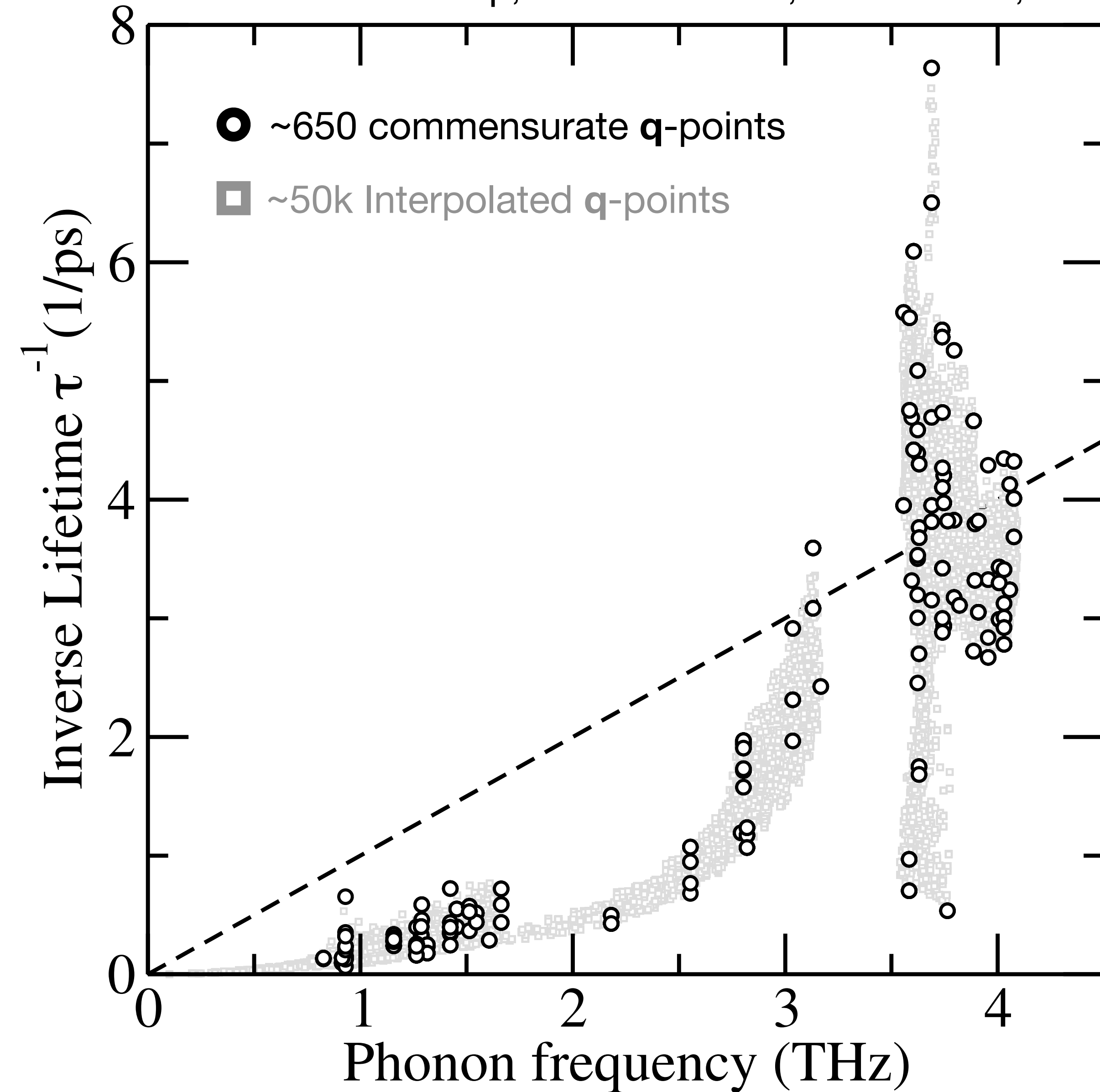
Anharmonicity in Copper Iodide

F. Knoop, T. A. R. Purcell, M. Scheffler, and C. Carbogno, *Phys. Rev. Lett.* **130**, 236301 (2023).



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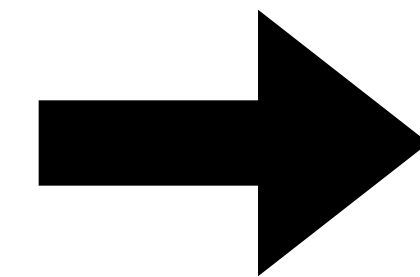
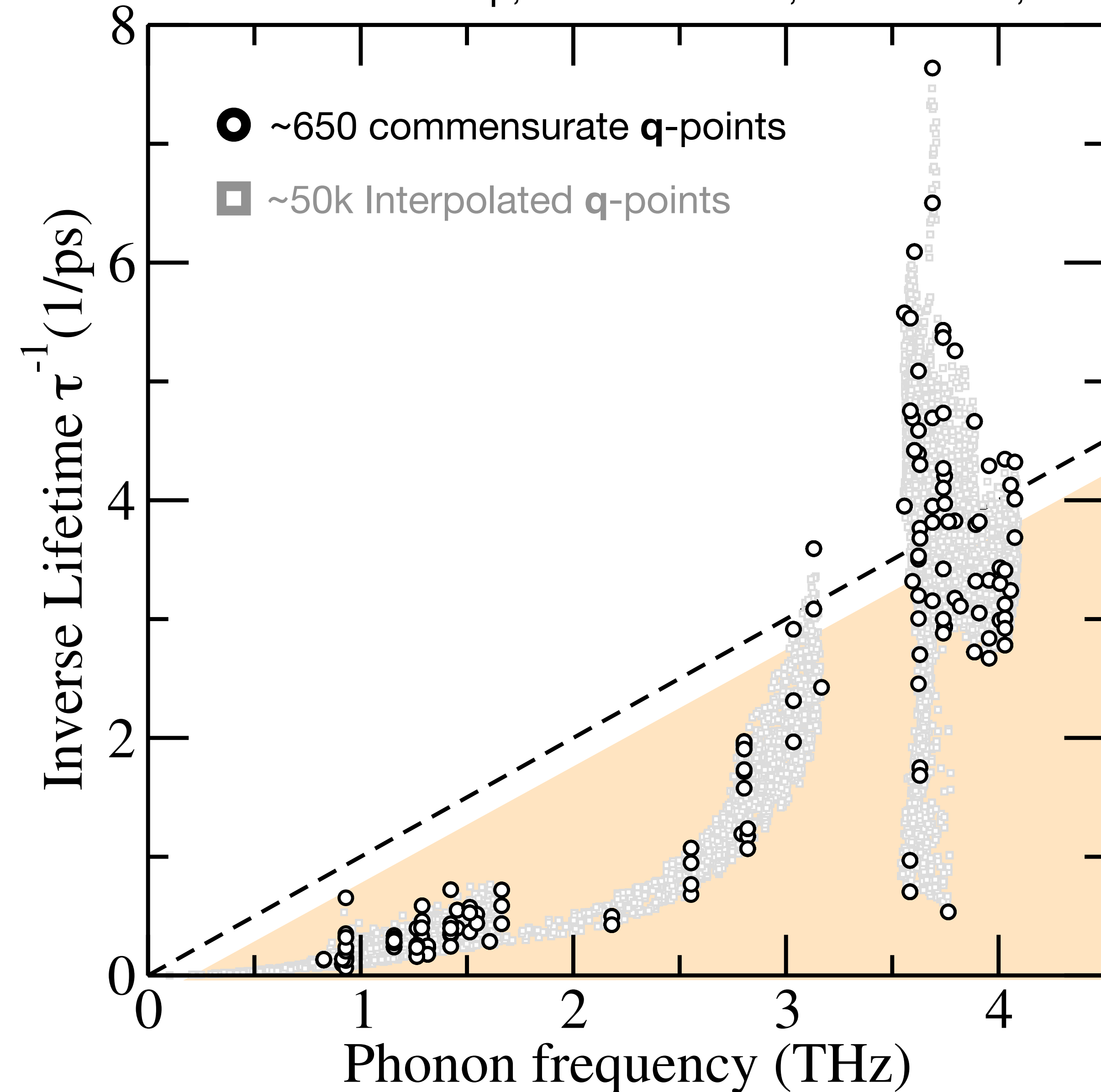


Ioffe-Regel Limit: $\tau > 1/\omega$

A. F. Ioffe and A. R. Regel, *Prog. Semicond.* **4**, 237 (1960).

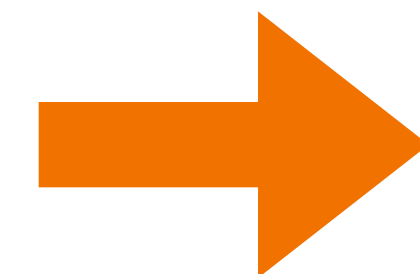
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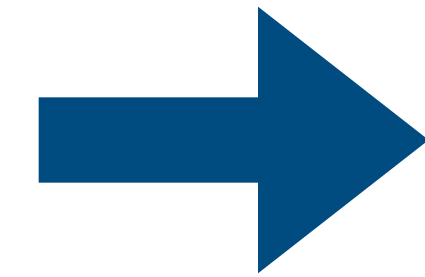
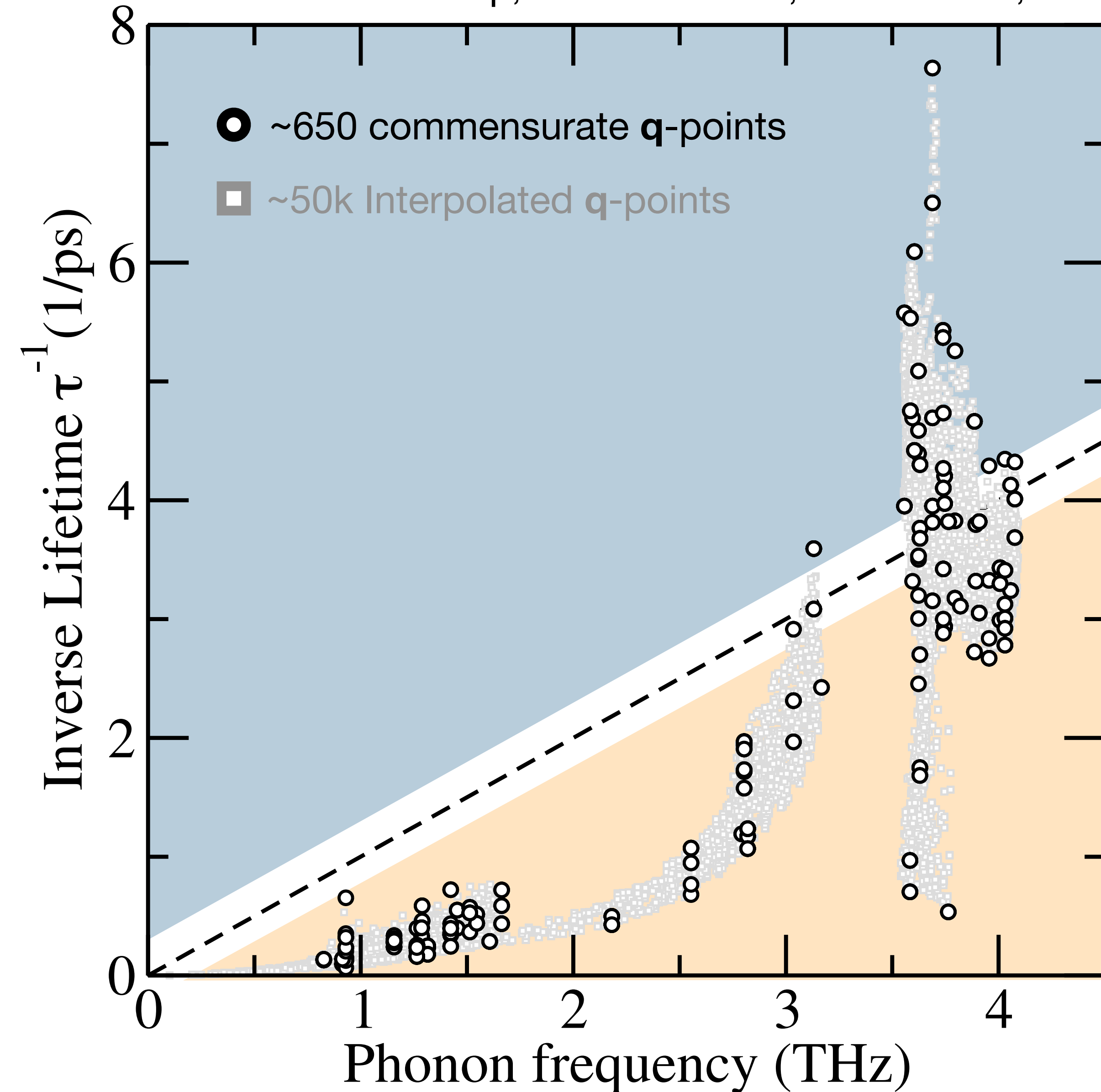


Phonon Picture

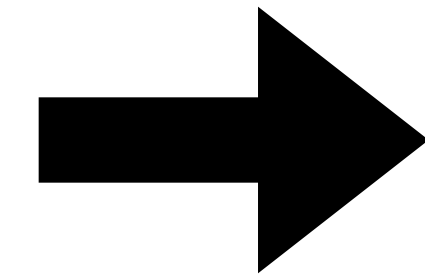
M. Simoncelli, N. Marzari, and F. Mauri,
Phys Rev X **12**, 041011 (2022).

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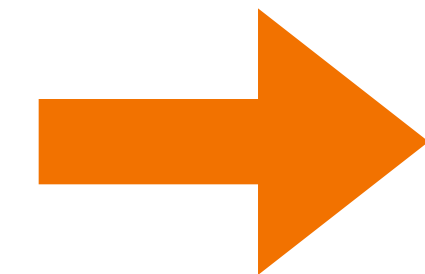


Strongly Anharmonic Regime



Ioffe-Regel Limit: $\tau > 1/\omega$

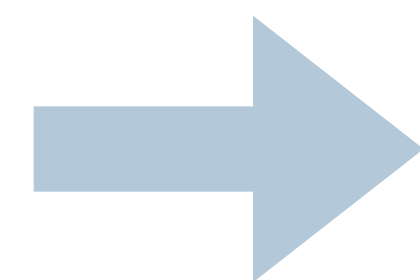
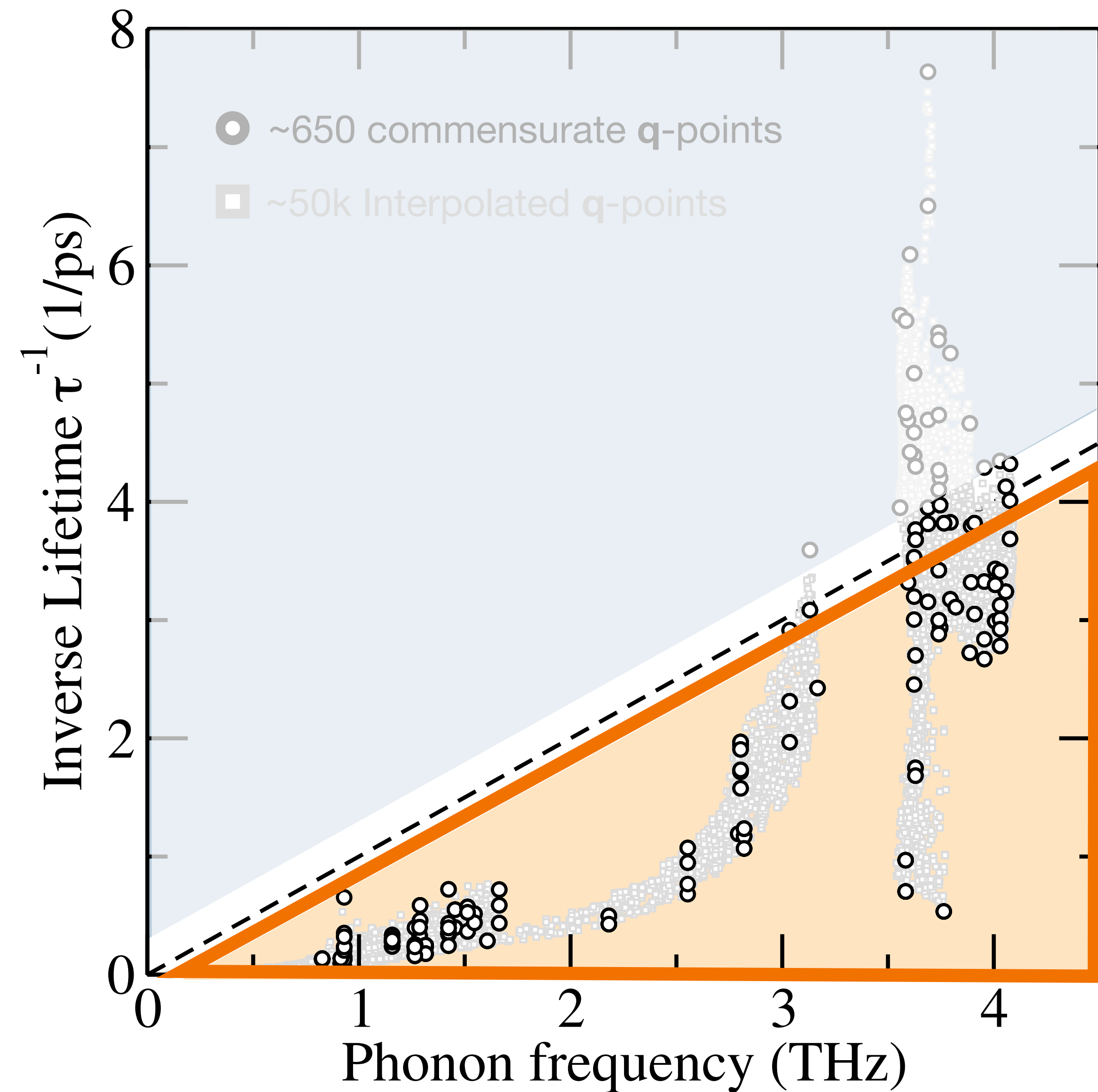
A. F. Ioffe and A. R. Regel, *Prog. Semicond.* **4**, 237 (1960).



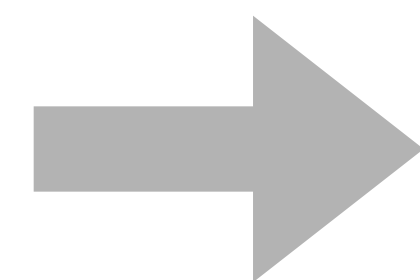
Phonon Picture

M. Simoncelli, N. Marzari, and F. Mauri,
Phys Rev X **12**, 041011 (2022).

Validity of Phonon-based Theories

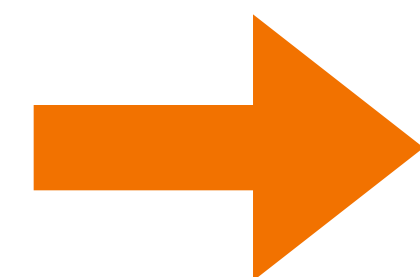


Strongly Anharmonic Regime



Ioffe-Regel Limit: $\tau > 1/\omega$

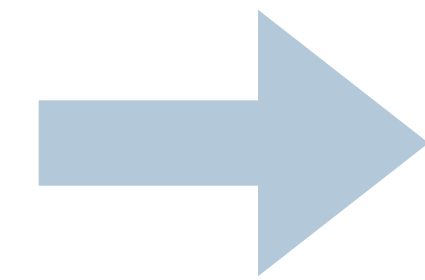
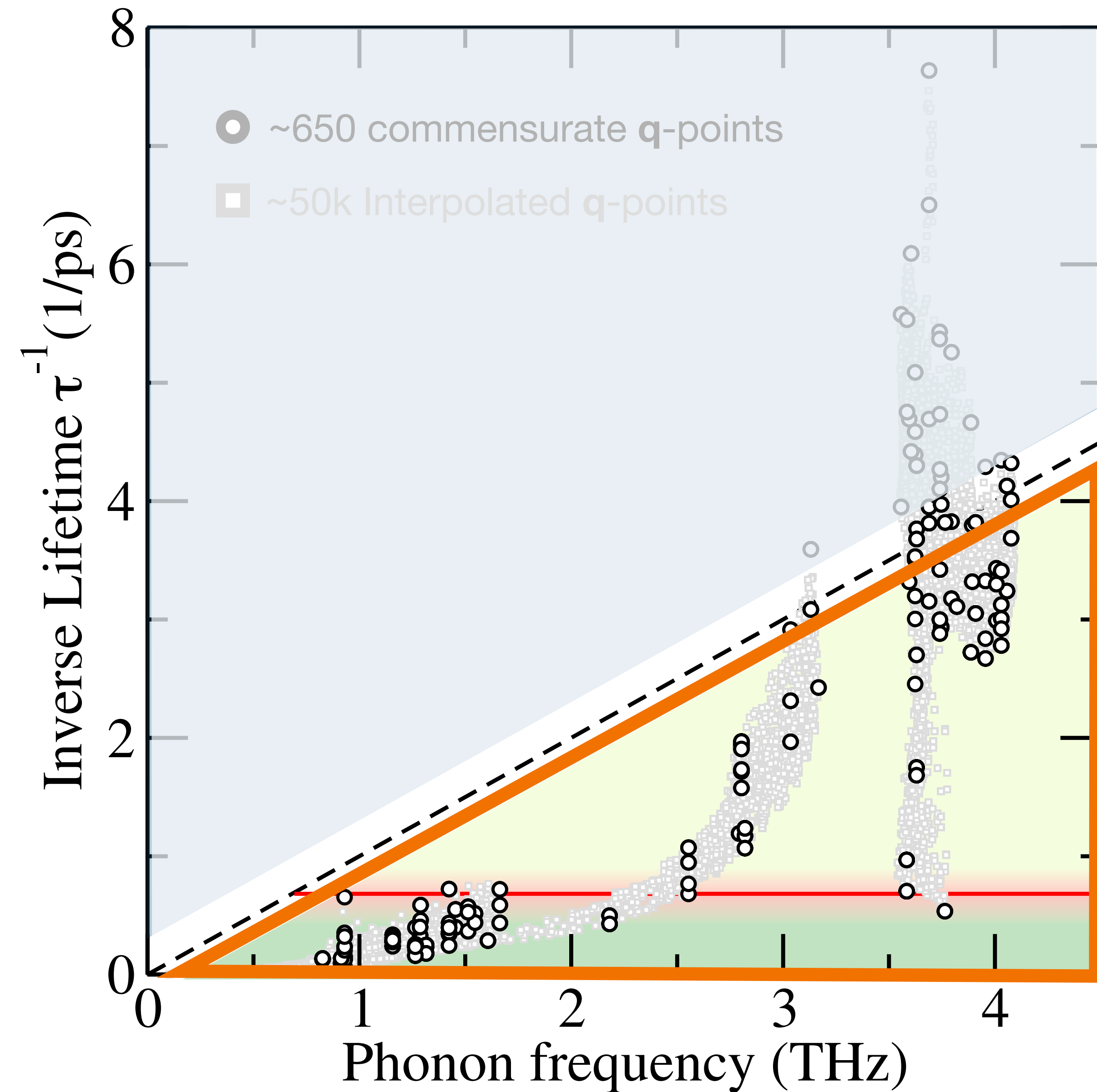
A. F. Ioffe and A. R. Regel, *Prog. Semicond.* **4**, 237 (1960).



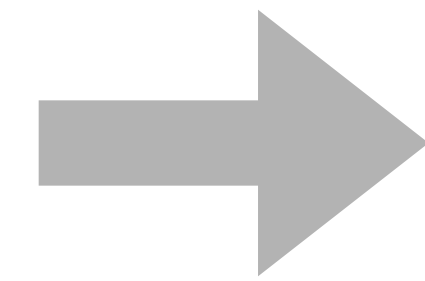
Phonon Picture

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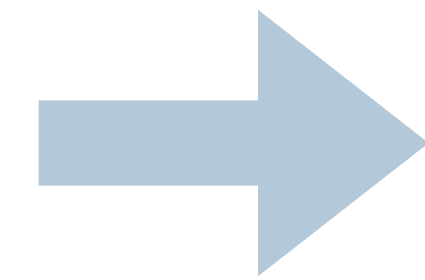
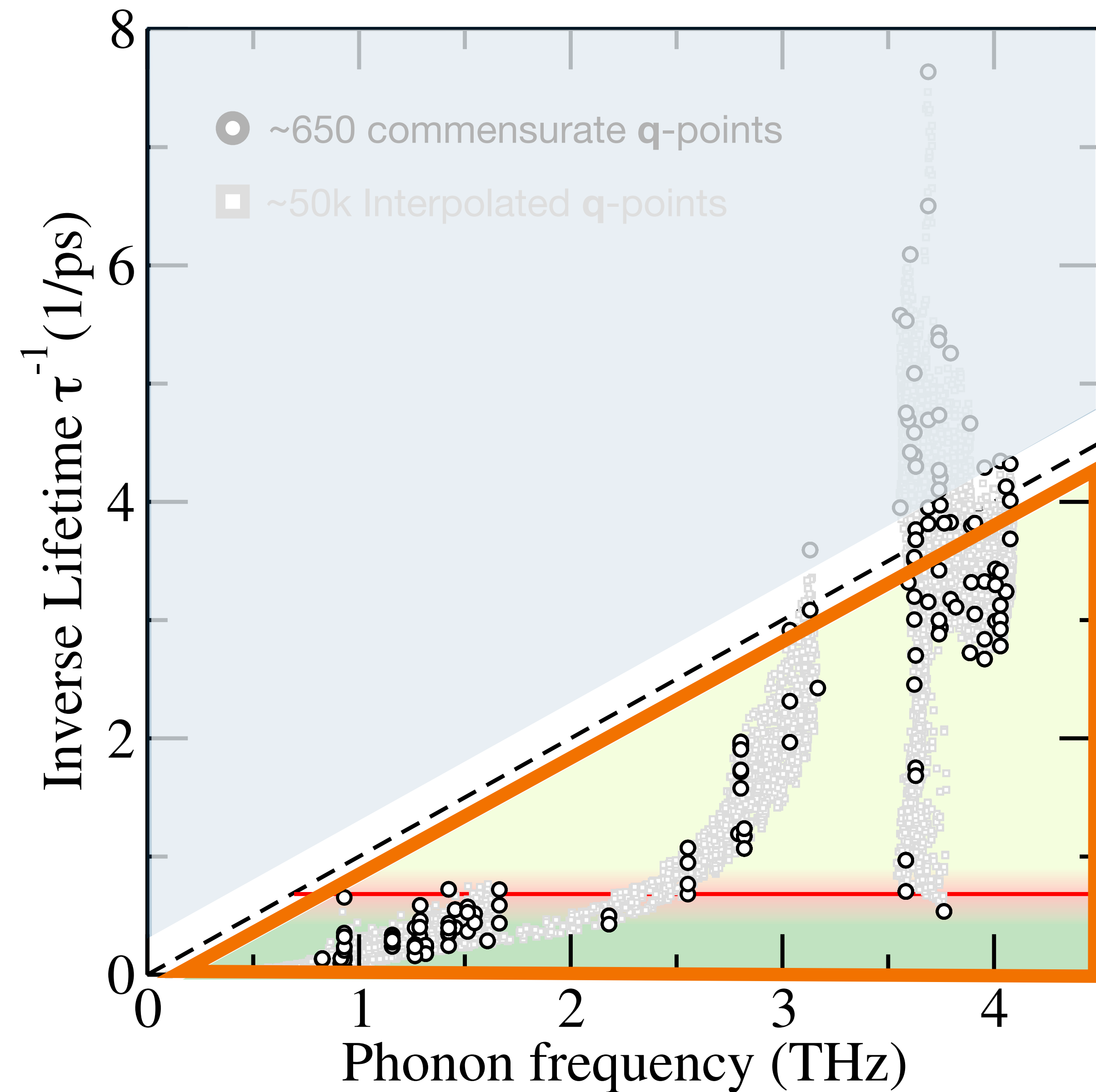
Strongly Anharmonic Regime



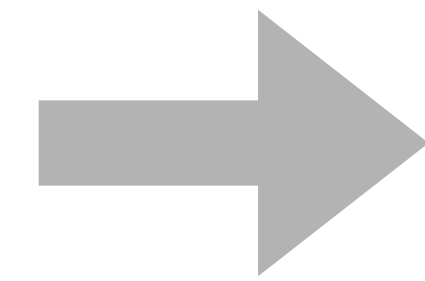
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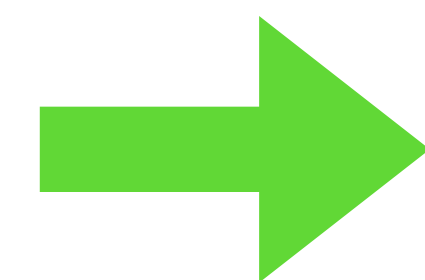


Strongly Anharmonic Regime



Ioffe-Regel Limit: $\tau > 1/\omega$

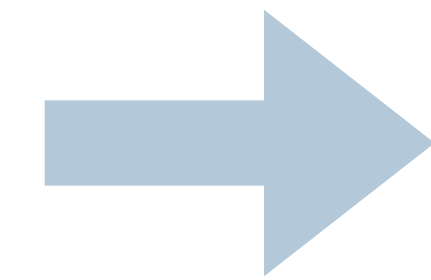
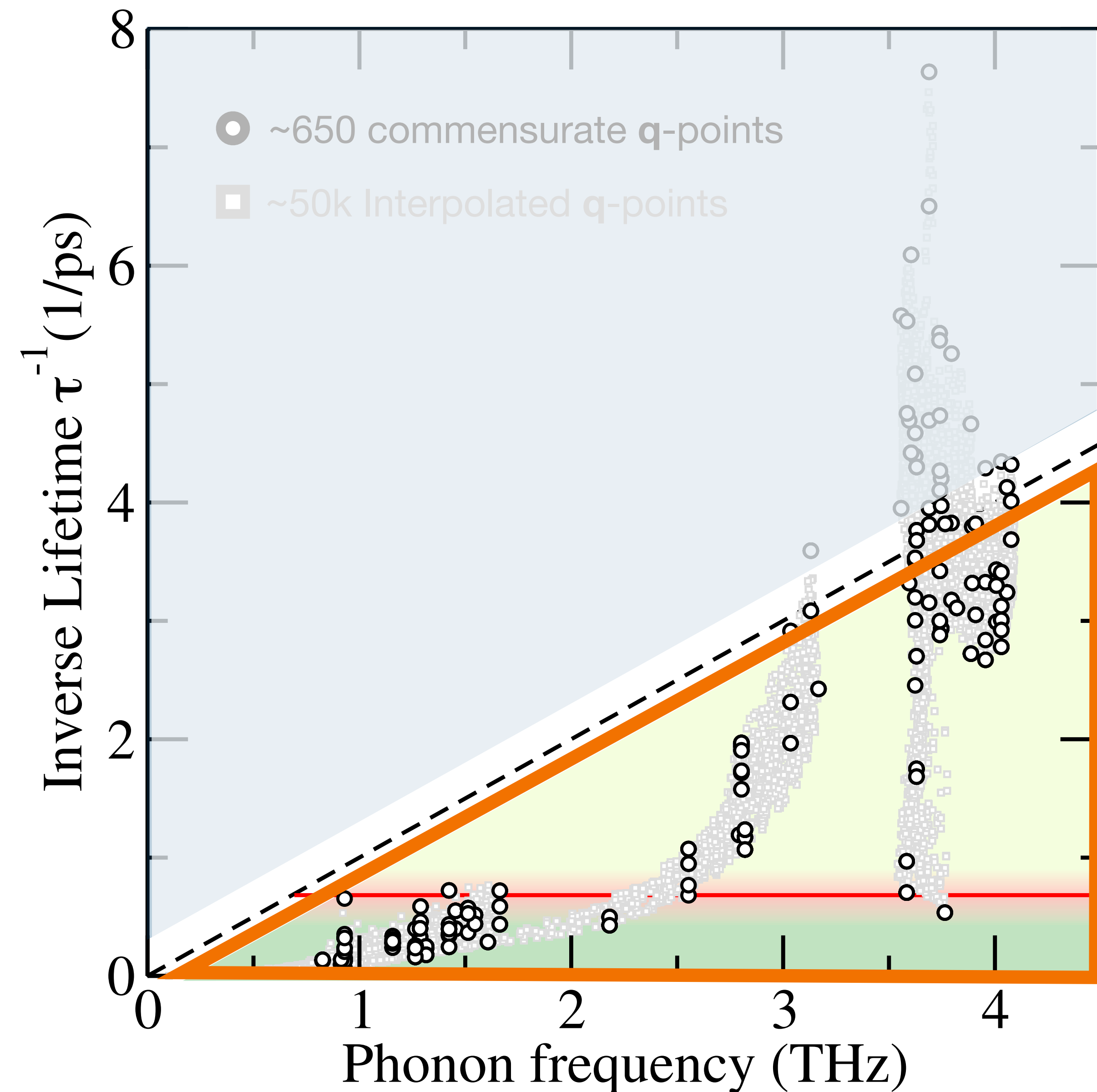
A. F. Ioffe and A. R. Regel, *Prog. Semicond.* **4**, 237 (1960).



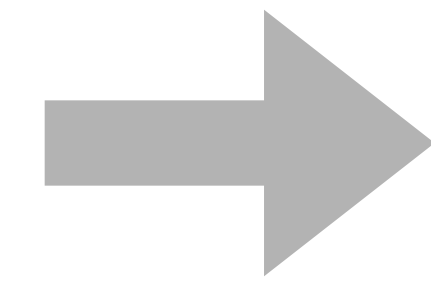
Wave-like Wigner Transport

M. Simoncelli, N. Marzari, and F. Mauri, *Phys Rev X* **12**, 041011 (2022).

Validity of Phonon-based Theories

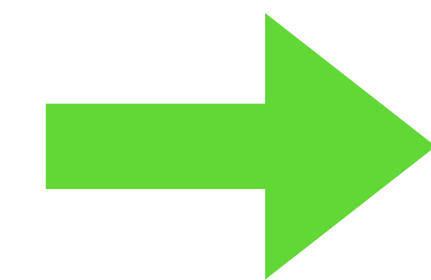


Strongly Anharmonic Regime



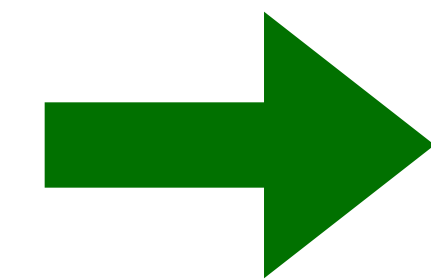
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Wave-like Wigner Transport

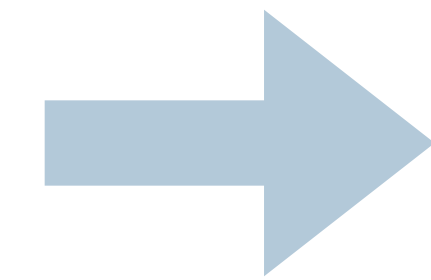
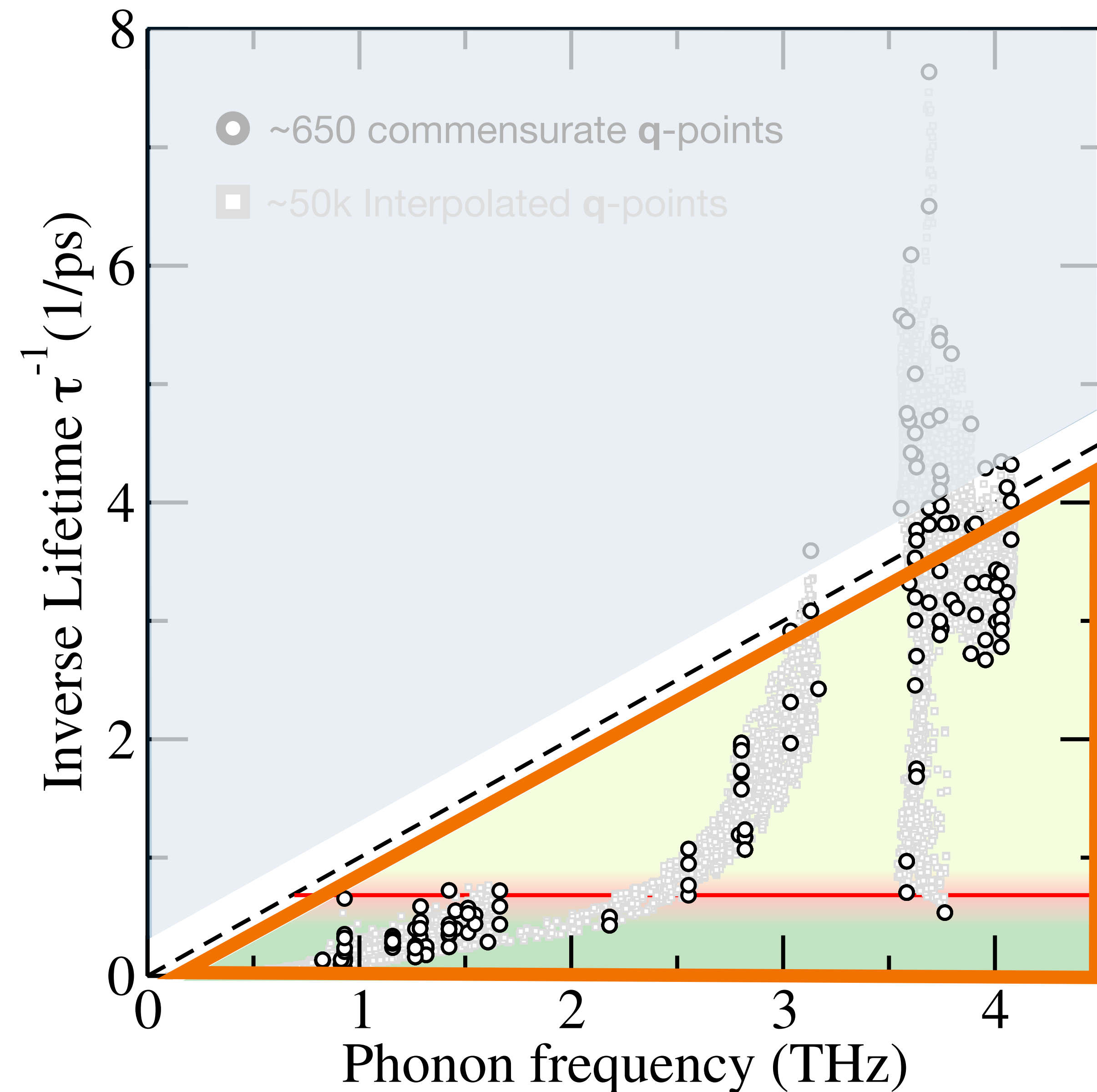
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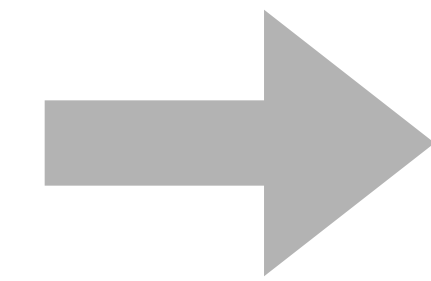
Particle-like Boltzmann Transport

D. A. Broido *et al.*, *Appl. Phys. Lett.* **91**, 231922 (2007).

Validity of Phonon-based Theories

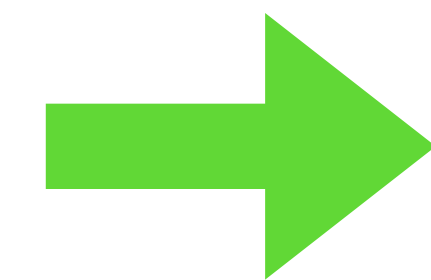


Strongly Anharmonic Regime



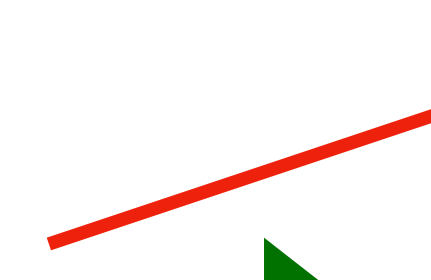
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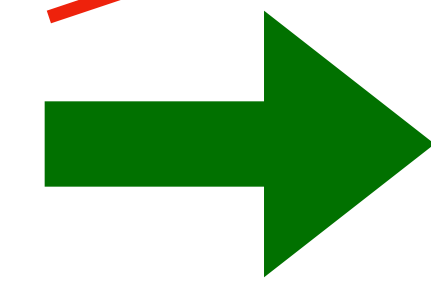


Wave-like Wigner Transport

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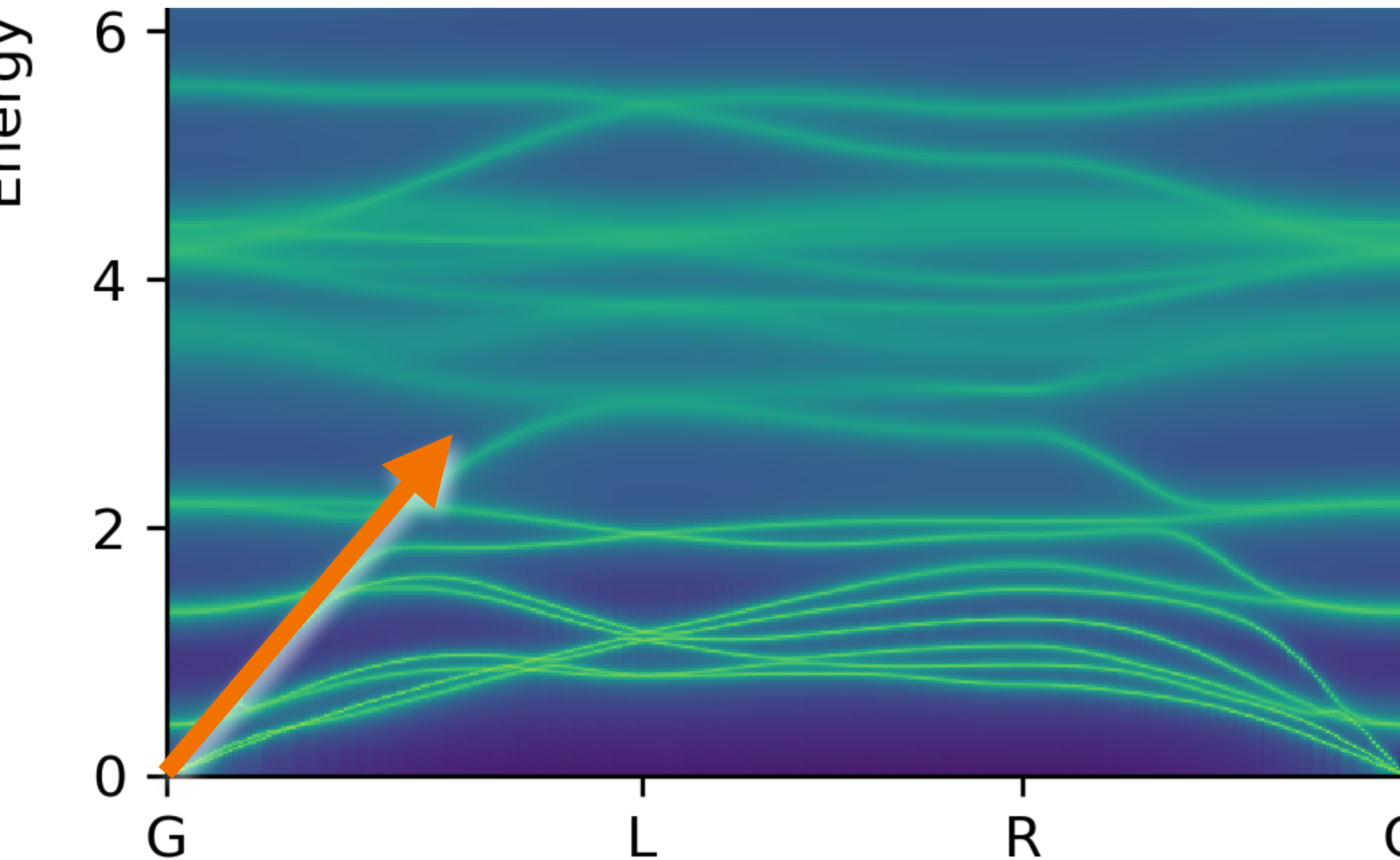
Average interband spacing



Particle-like Boltzmann Transport

D. A. Broido *et al.*, *Appl. Phys. Lett.* **91**, 231922 (2007).

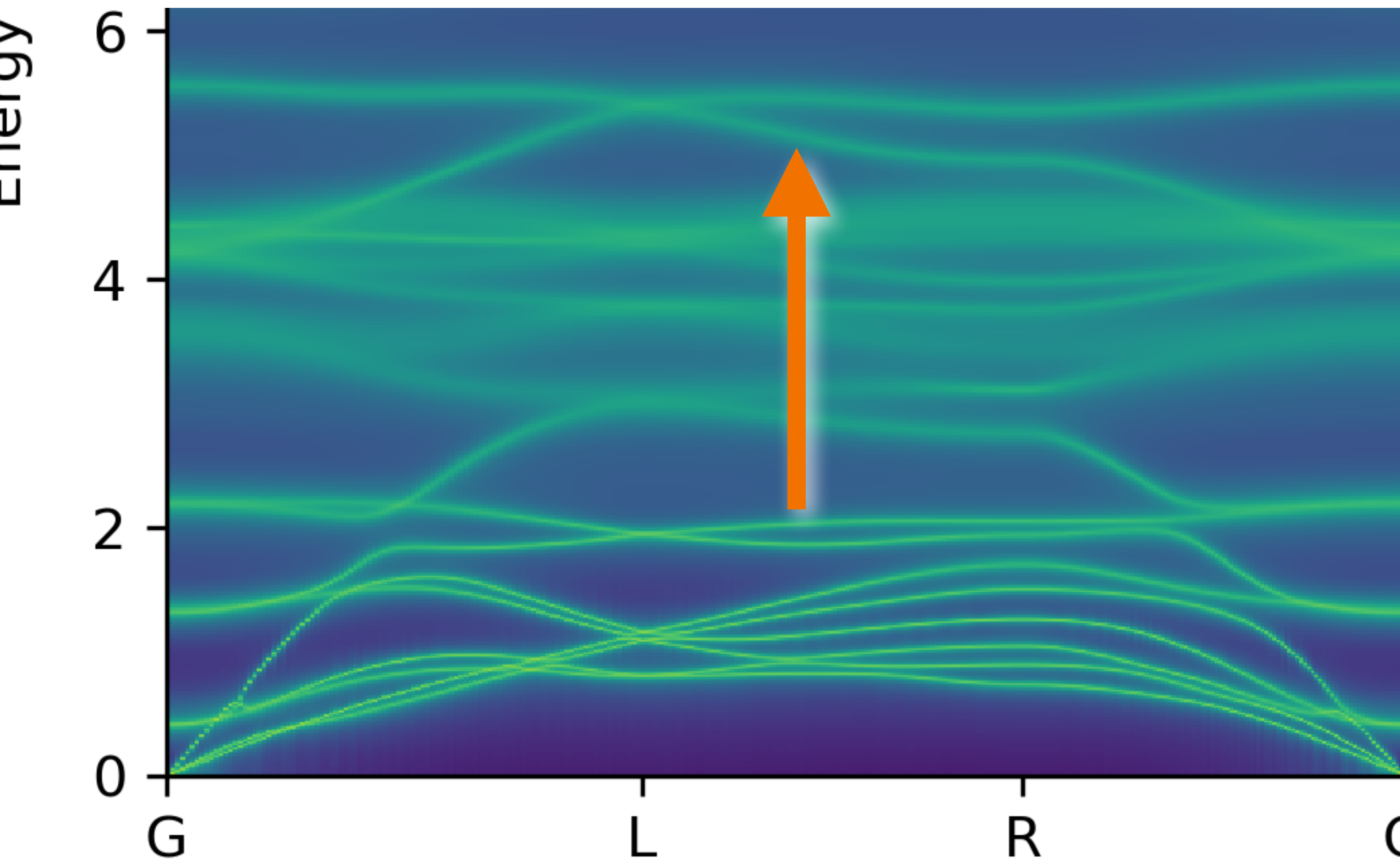
Wigner vs. Boltzmann Transport



Boltzmann Transport:

Transport determined by **group** velocity, i.e., the diagonal of the **momentum matrix**: $\mathbf{v}_g = \langle s | \mathbf{v} | s \rangle$

Wigner vs. Boltzmann Transport



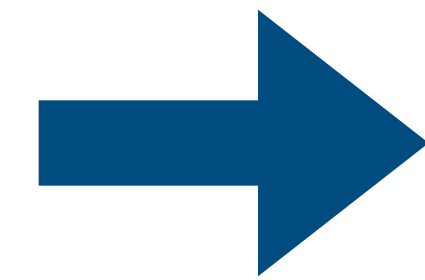
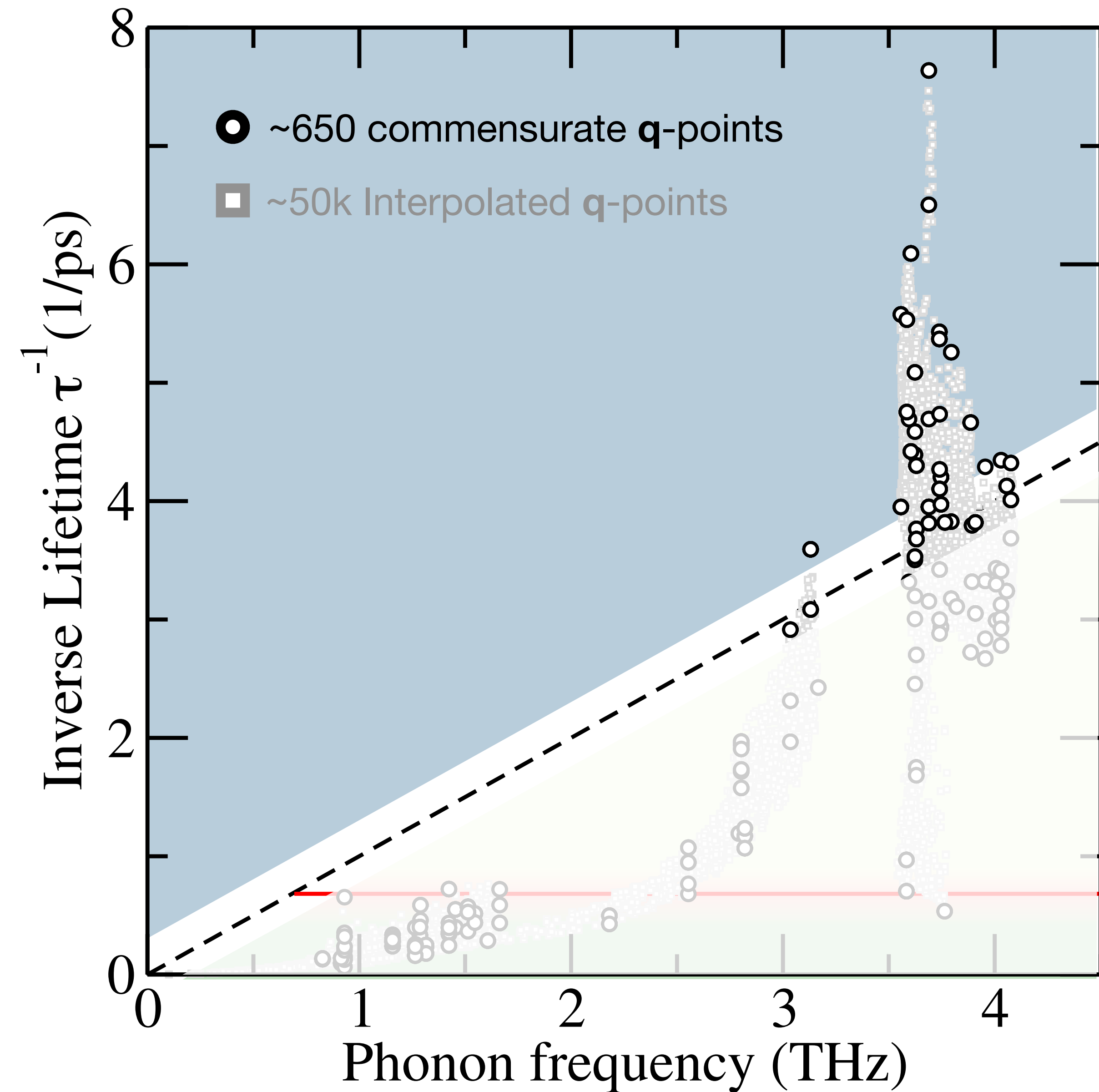
Boltzmann Transport:

Transport determined by **group** velocity, i.e., the diagonal of the **momentum matrix**: $\mathbf{v}_g = \langle s | \mathbf{v} | s \rangle$

Wigner Transport:

Transport determined by **direct transition**, i.e., the off-diagonal of the **momentum matrix**: $\langle s | \mathbf{v} | s' \rangle$

Validity of Phonon-based Theories



Strongly Anharmonic Regime

What happens with
electronic transport
in this regime?

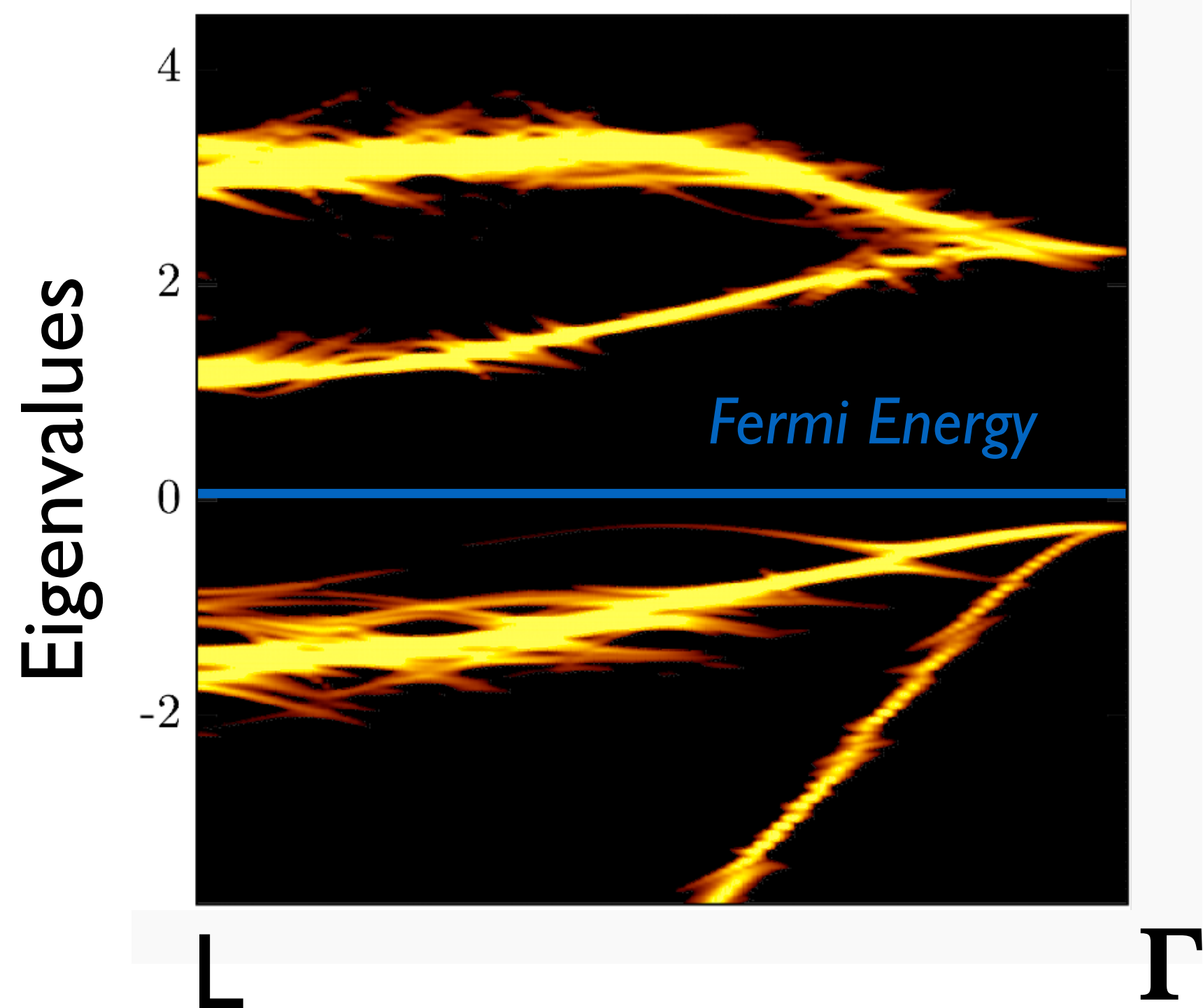
Anharmonic Effects in Electronic Transport

Ab initio MD

Band Structure
Unfolding

Fully Anharmonic, non-perturbative assessment of
vibronically renormalized electronic band structures.

M. Zacharias, M. Scheffler, and C. Carbogno,
Phys. Rev. B **102**, 045126 (2020).



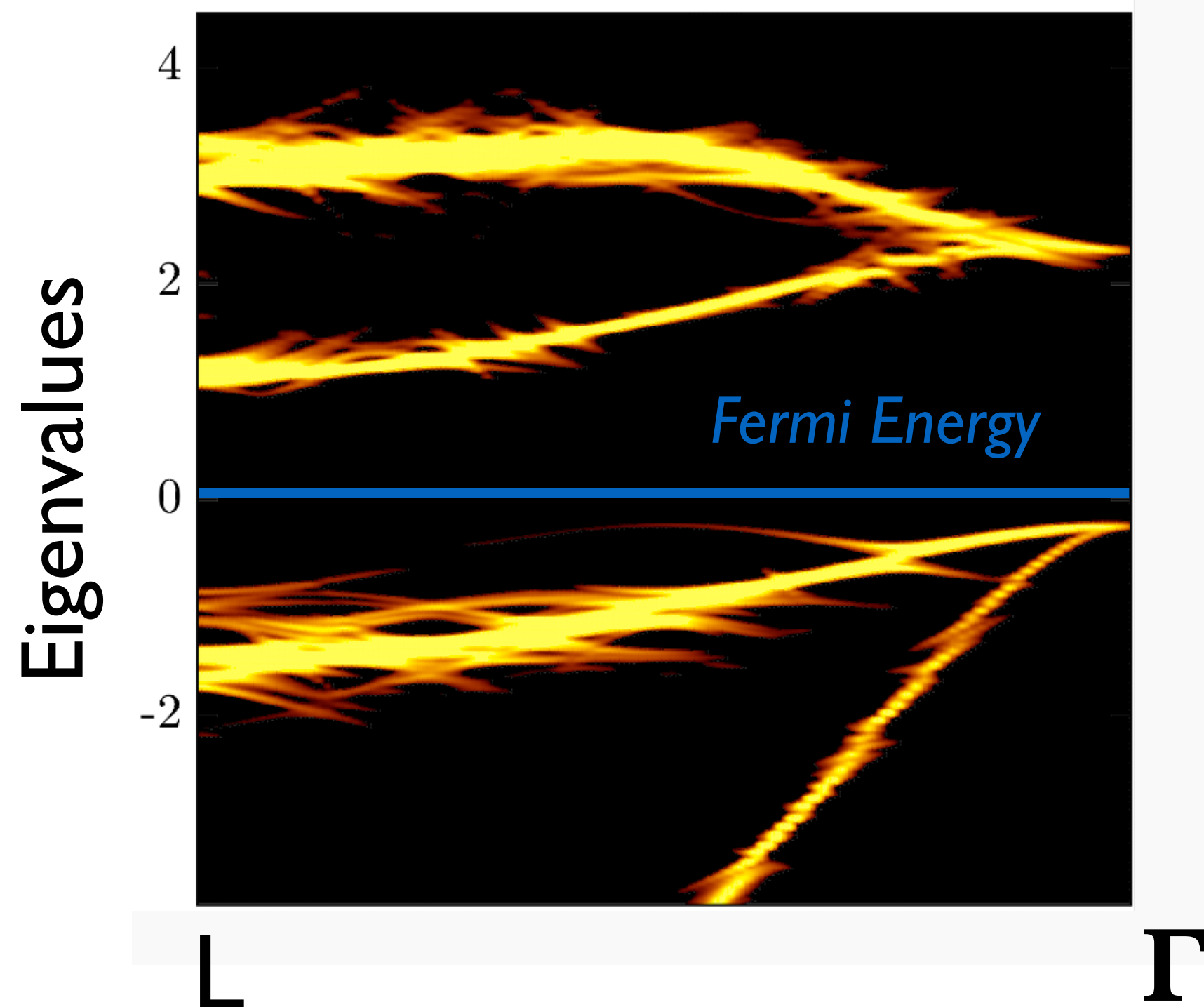
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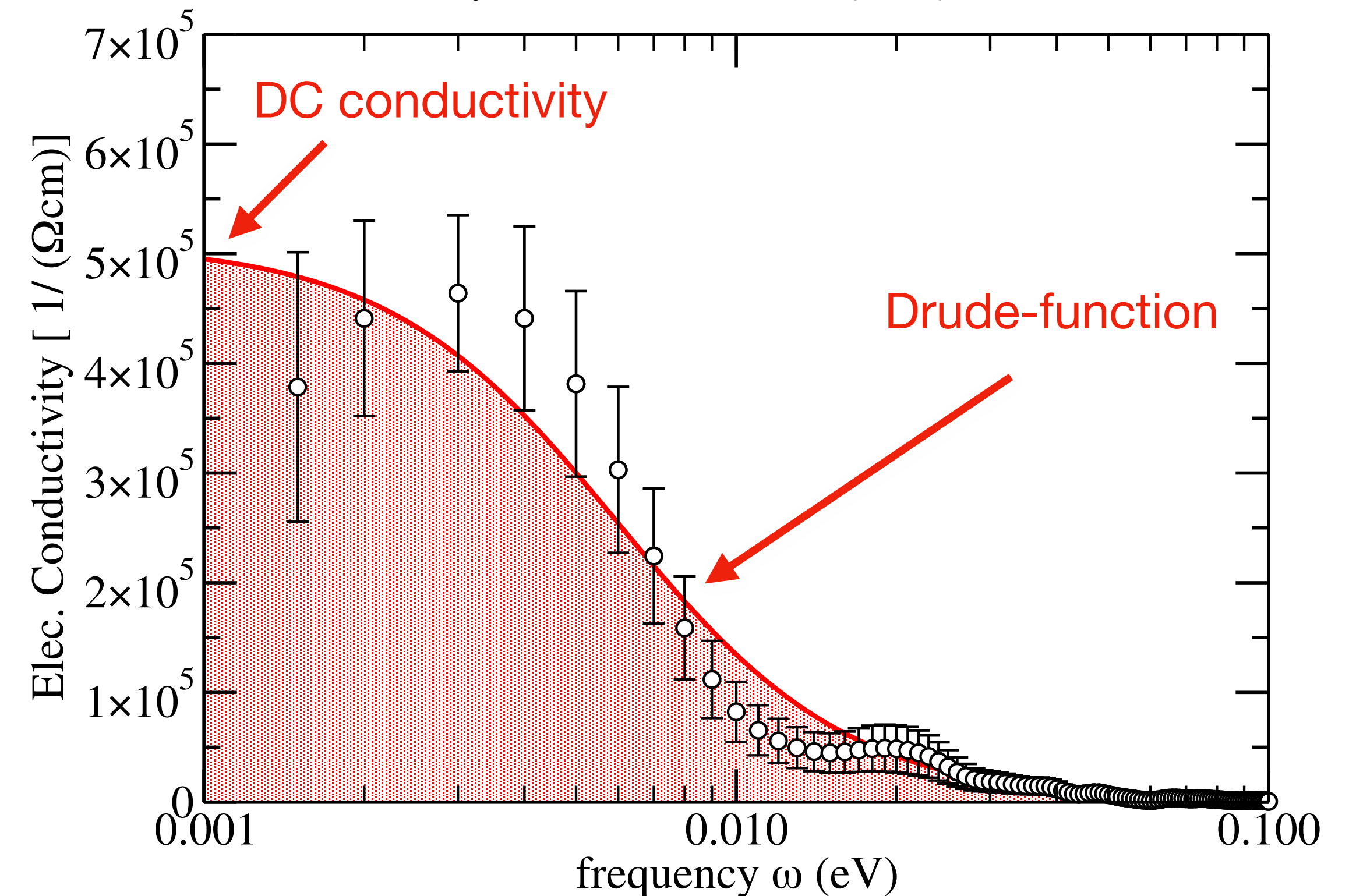


Ab initio MD

Kubo-Greenwood
Formalism

Fully Anharmonic, non-perturbative assessment of electronic transport coefficients.

B. Holst, M. French, and R. Redmer,
Phys. Rev. B **83**, 235120 (2011).

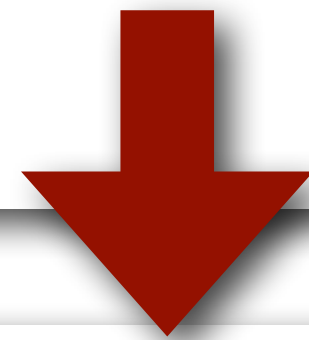


GREENWOOD-KUBO FORMALISM

D.A. Greenwood, *Proc. Phys. Soc.* **71**, 585 (1958).

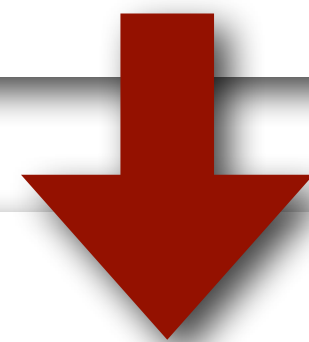
Kubo's Linear Response:

$$\sigma(\omega) = \frac{1}{V} \left\langle \lim_{\varepsilon \rightarrow 0} \int_0^{\infty} dt e^{i(\omega+i\varepsilon)t} \int_0^{(k_B T)^{-1}} d\tau \mathbf{Tr} [\hat{\rho}_0 \mathbf{j}_c(t - i\hbar\tau) \cdot \mathbf{j}_c(t)] \right\rangle_T$$



Independent Particle Picture:

$$\mathbf{j}_c = -\frac{e}{\hbar} \frac{\partial \varepsilon_n(\mathbf{k})}{\partial \mathbf{k}} \xrightarrow{\text{Heisenberg picture}} \mathbf{j}_c(t)$$

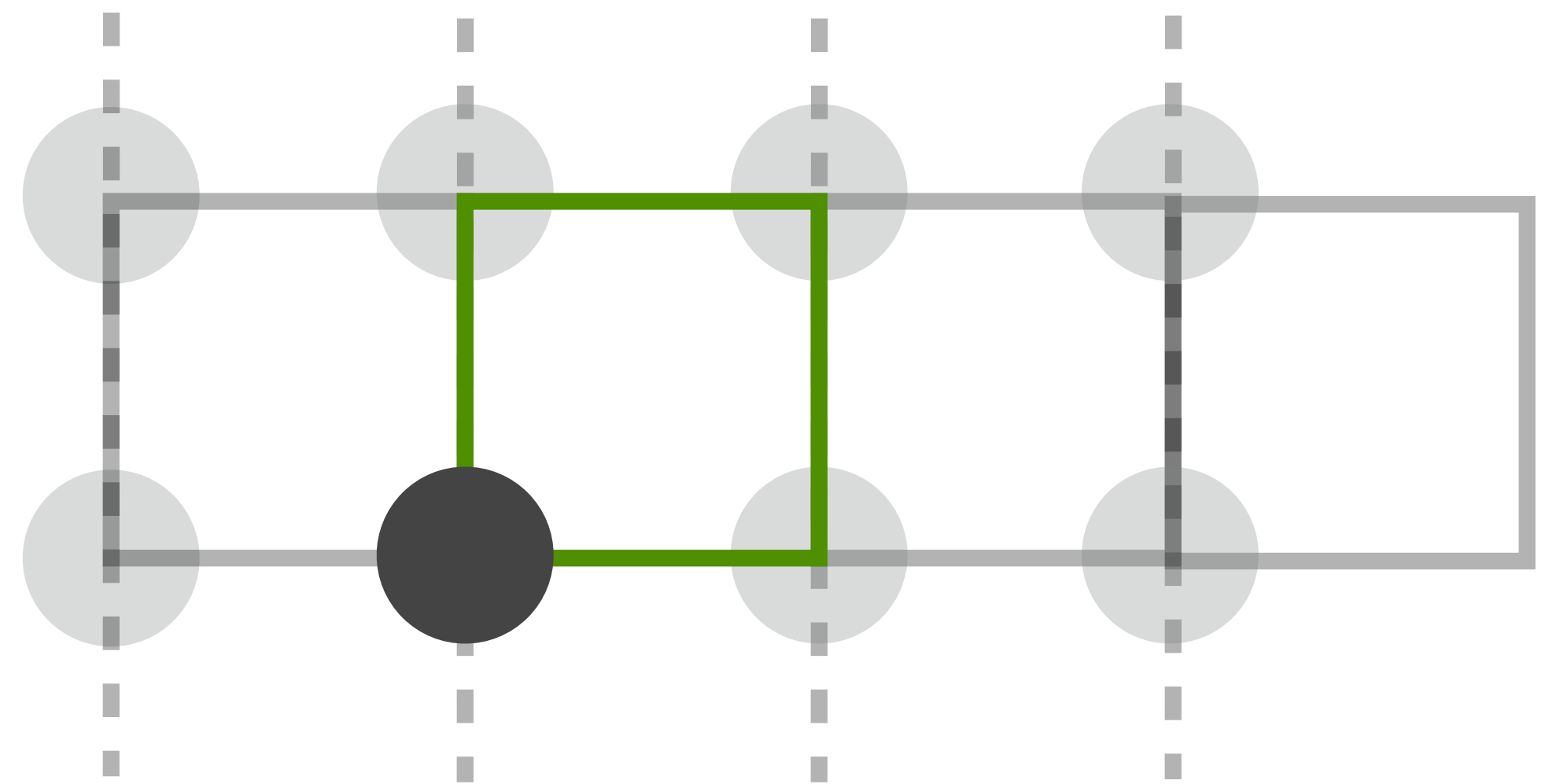
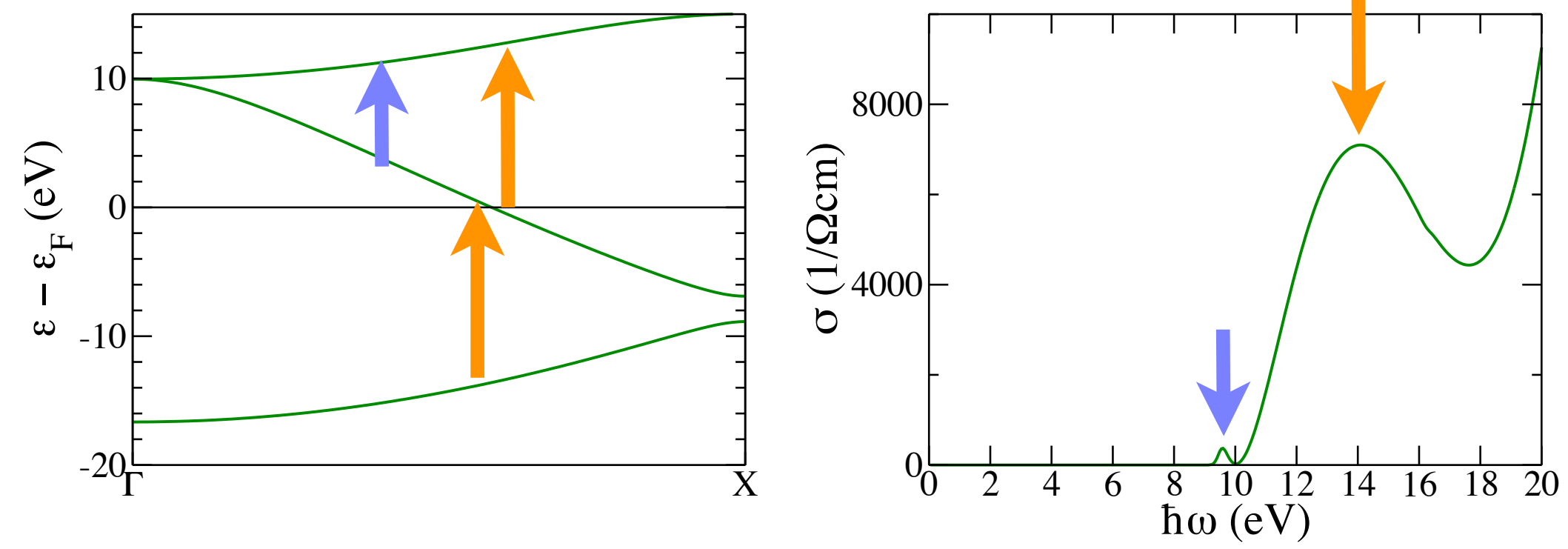


$$\sigma(\omega) = \frac{e^2 \hbar^2}{m_e^2 \omega} \frac{2\pi}{V} \left\langle \sum_{n, n \neq m} \sum_{\mathbf{k}} w_{\mathbf{k}} [f(\varepsilon_n) - f(\varepsilon_m)] |\langle n\mathbf{k} | \nabla | m\mathbf{k} \rangle|^2 \delta(\varepsilon_n - \varepsilon_m - \hbar\omega) \right\rangle_T$$

B. Holst, M. French, and R. Redmer, *Phys. Rev. B* **83**, 235120 (2011).

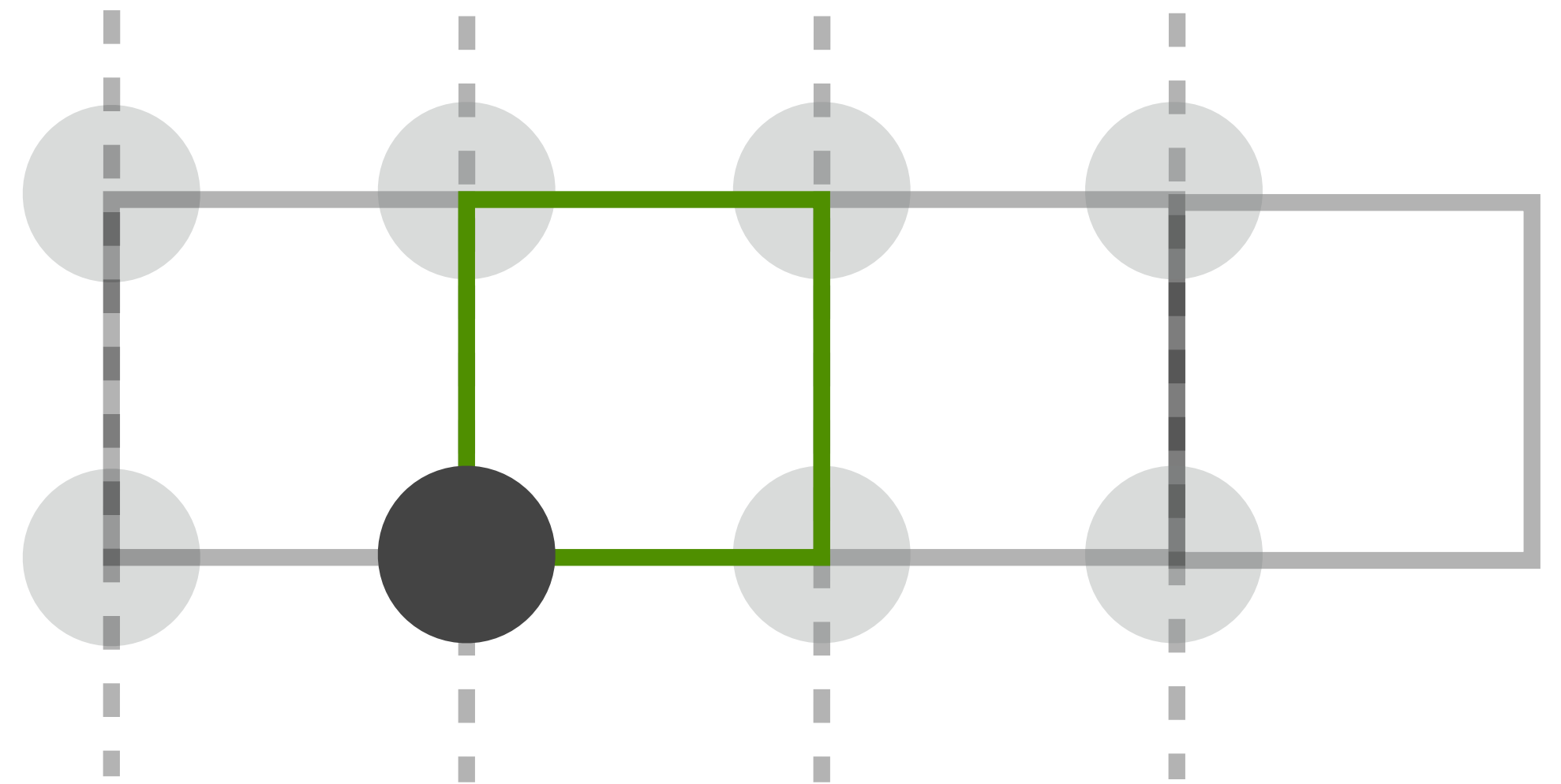
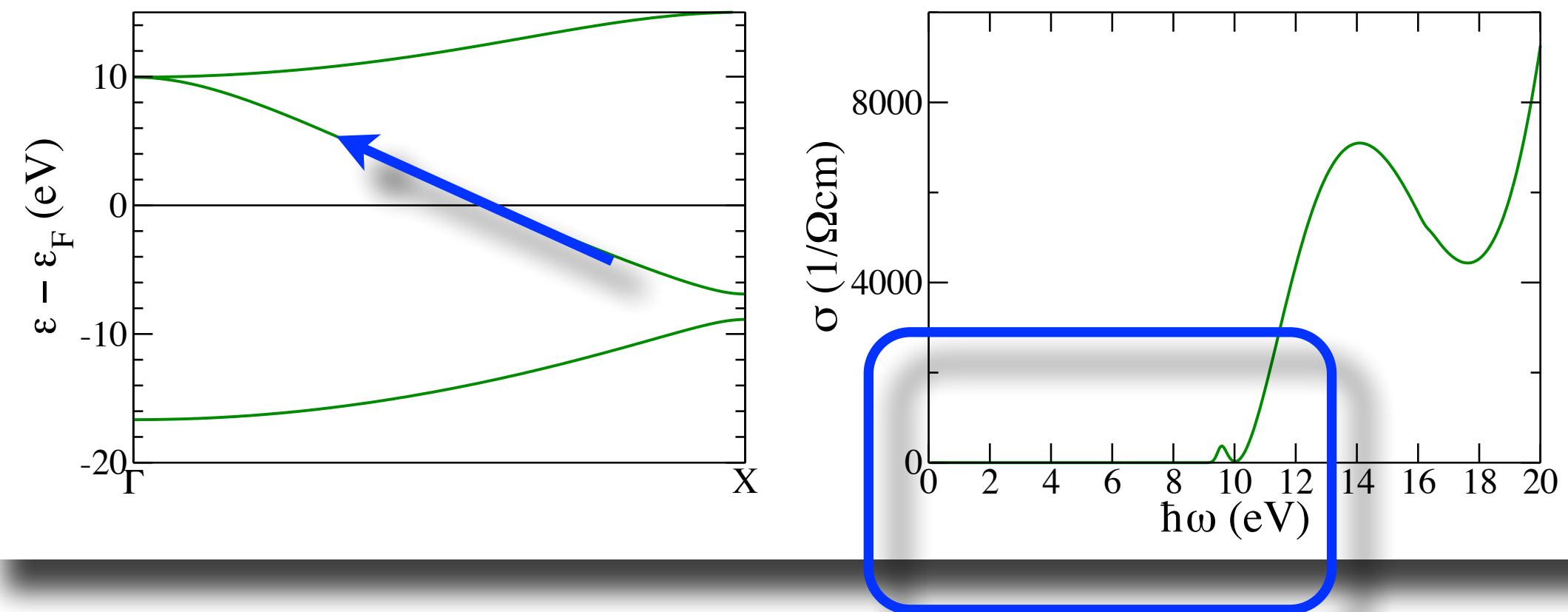
GREENWOOD-KUBO FORMALISM

D.A. Greenwood, *Proc. Phys. Soc.* **71**, 585 (1958).



GREENWOOD-KUBO FORMALISM

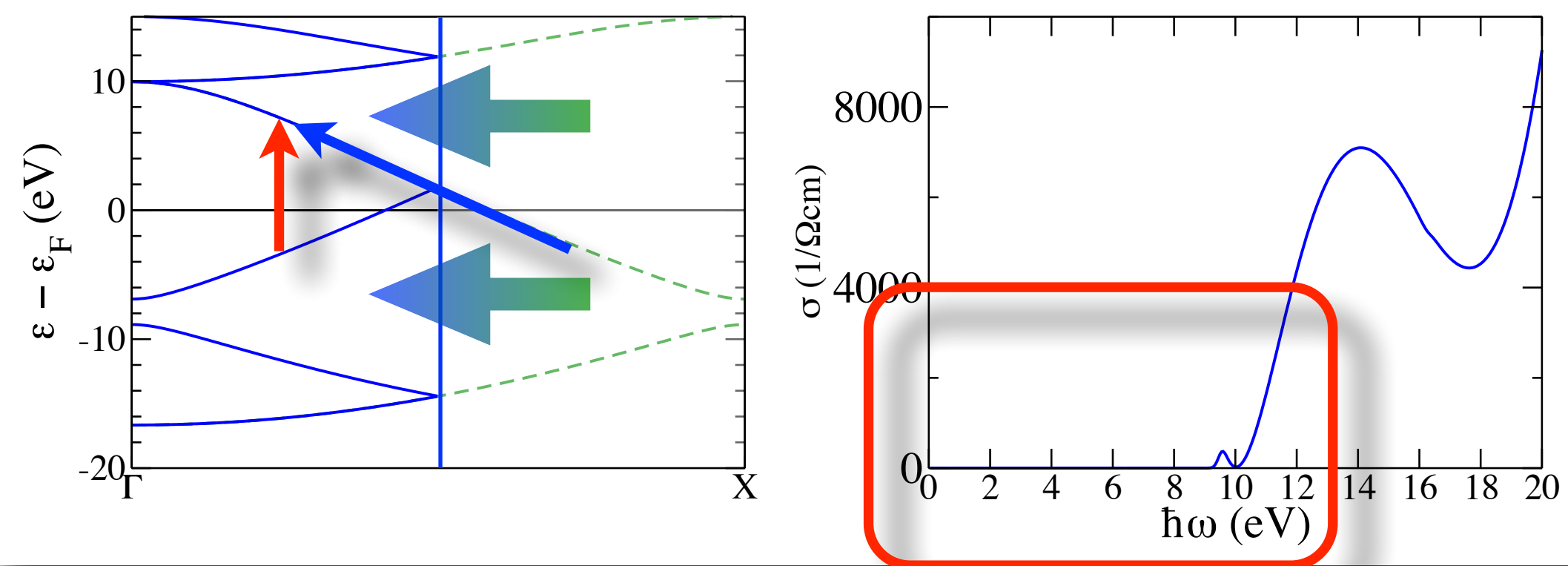
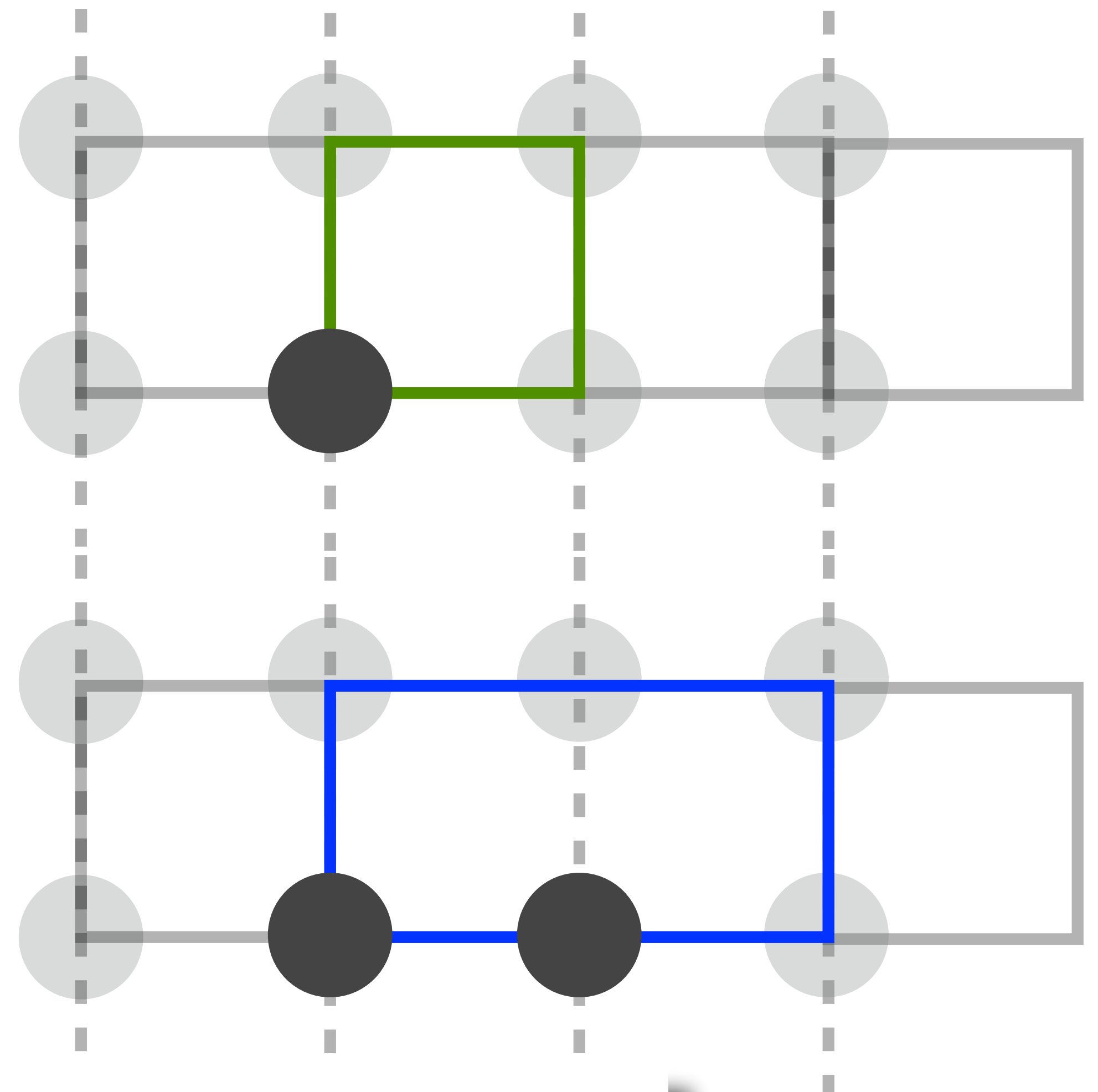
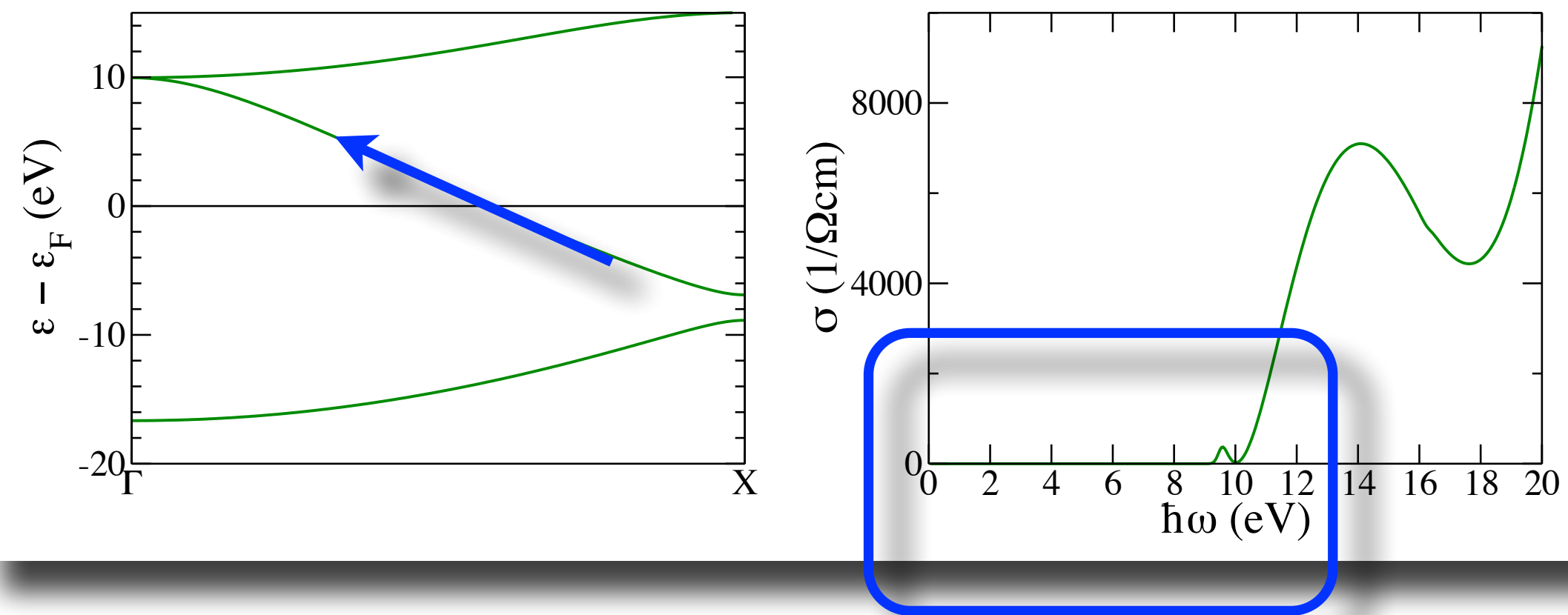
D.A. Greenwood, *Proc. Phys. Soc.* **71**, 585 (1958).



Crystal Momentum Conservation:
Non-vertical transitions require phonons

GREENWOOD-KUBO FORMALISM

D.A. Greenwood, *Proc. Phys. Soc.* **71**, 585 (1958).

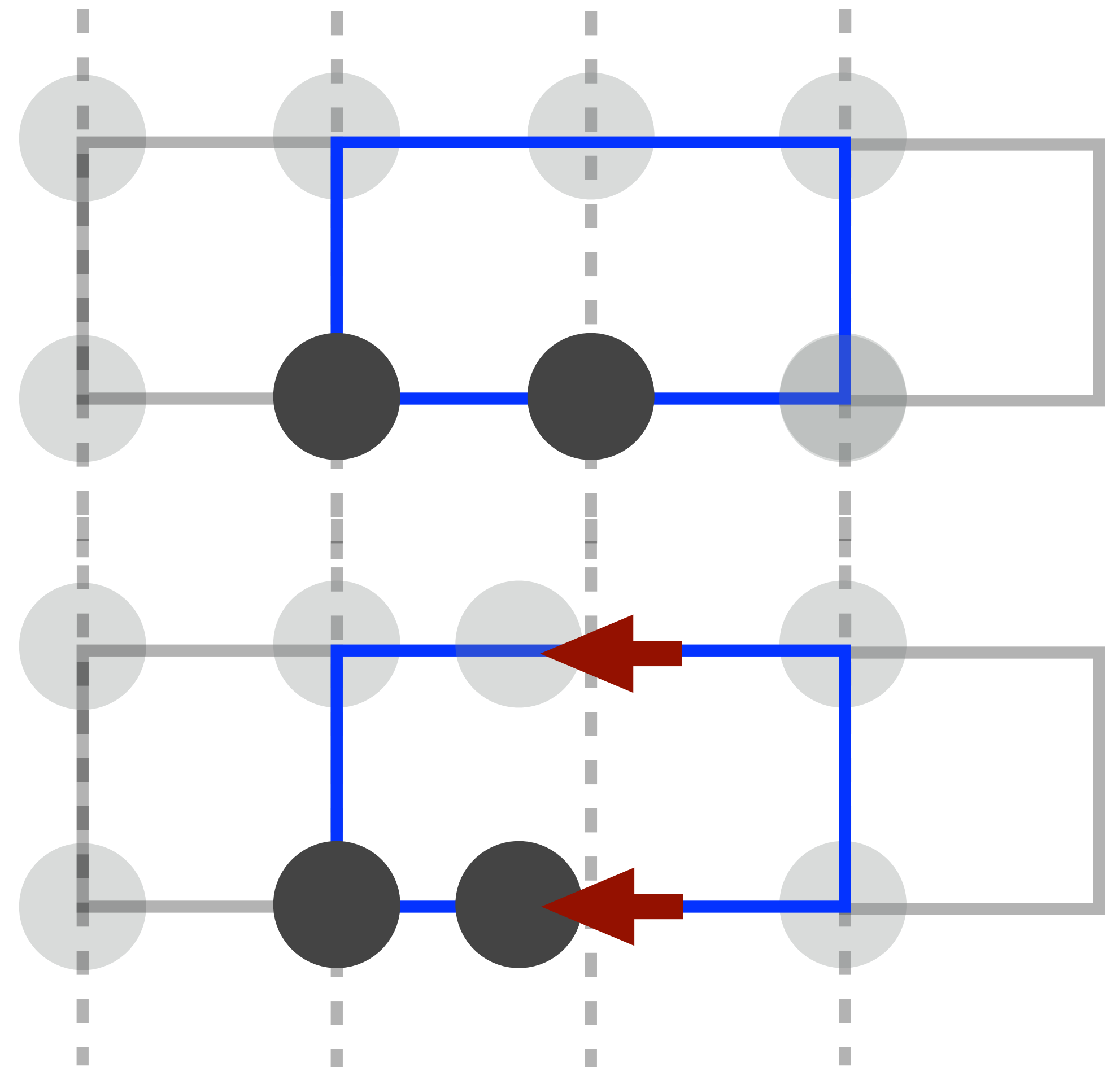
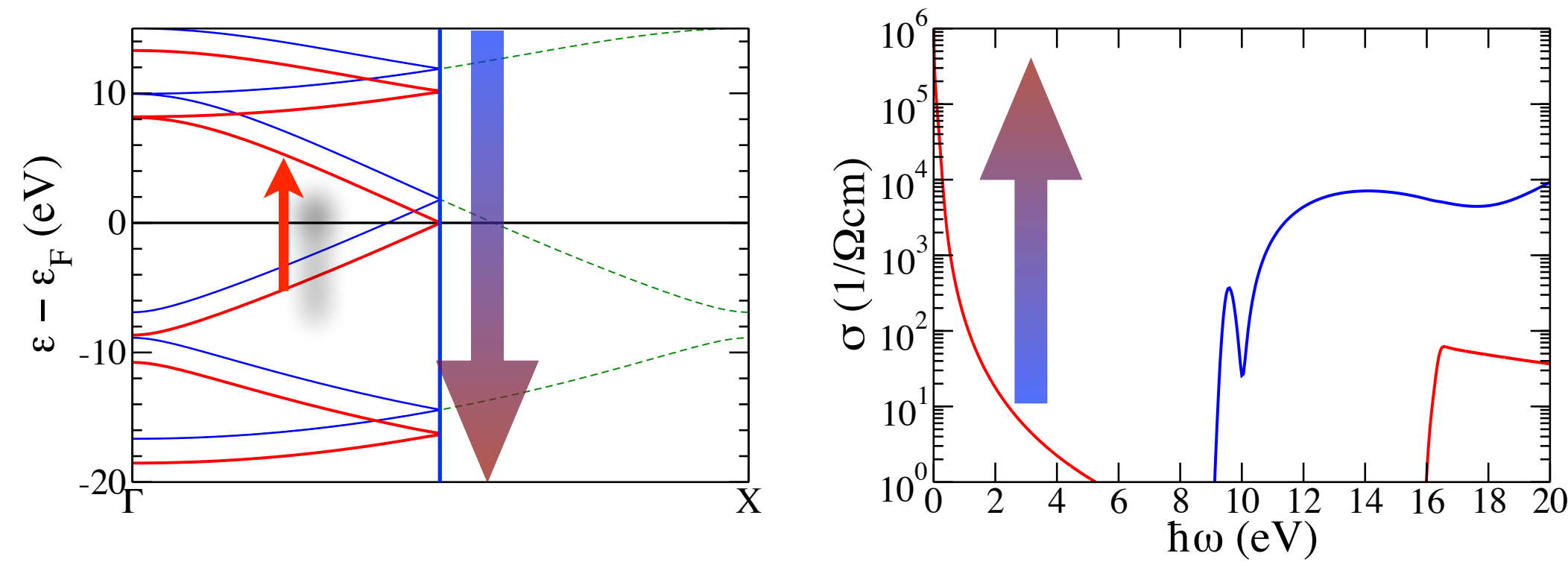
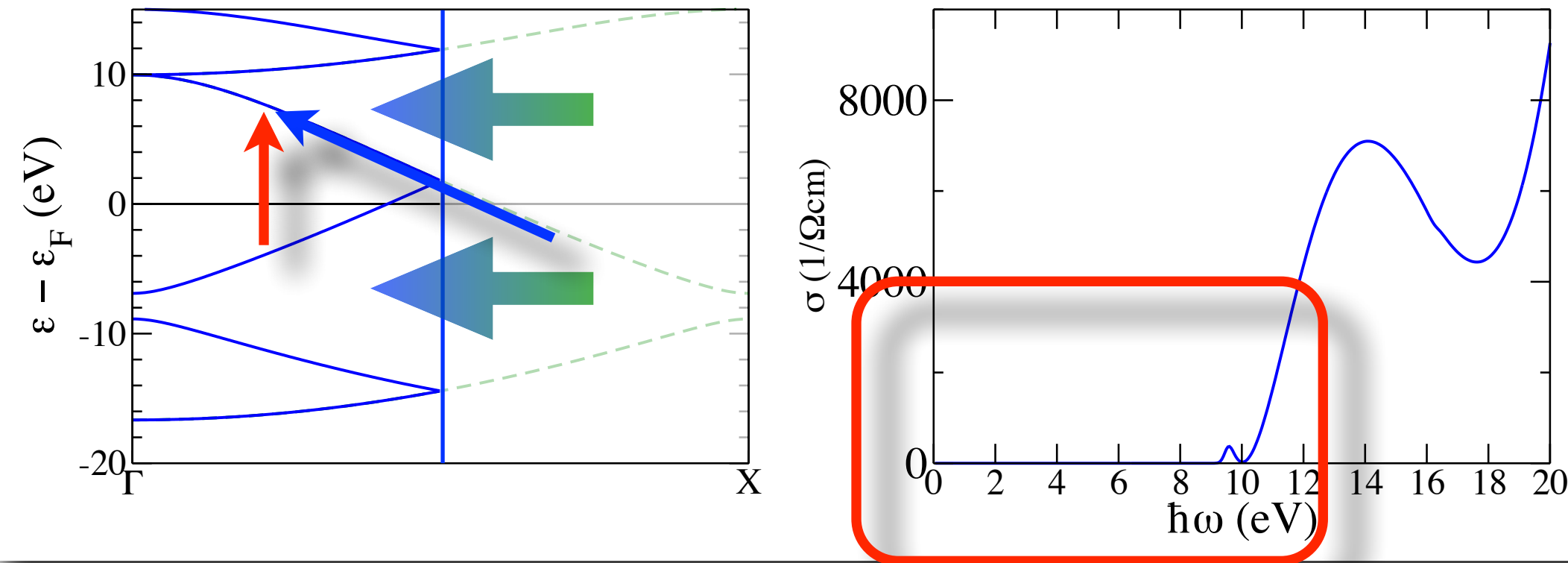


Brillouin zone folding:

Larger supercells allow for **direct transitions** that are however suppressed by **symmetry**.

GREENWOOD-KUBO FORMALISM

D.A. Greenwood, *Proc. Phys. Soc.* **71**, 585 (1958).

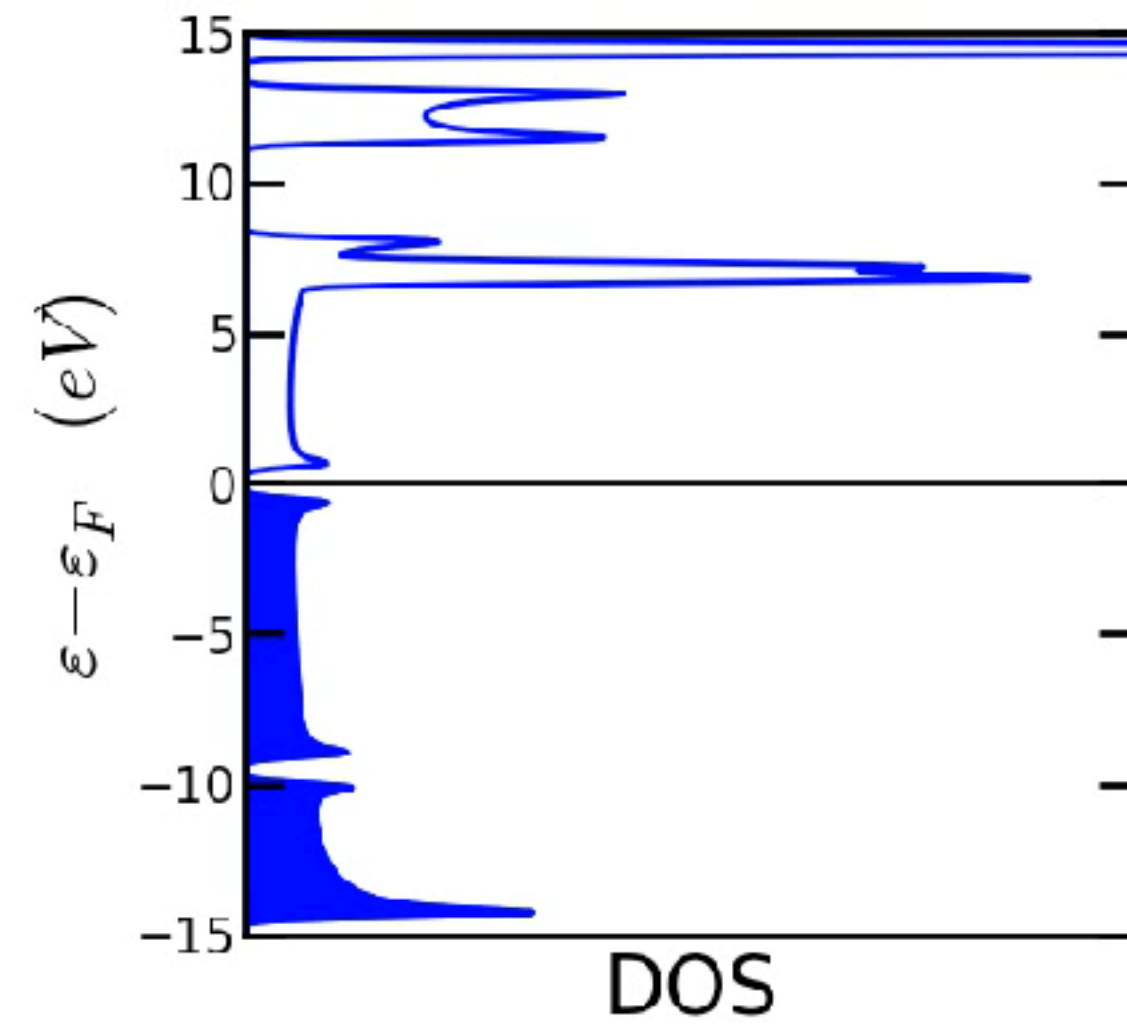
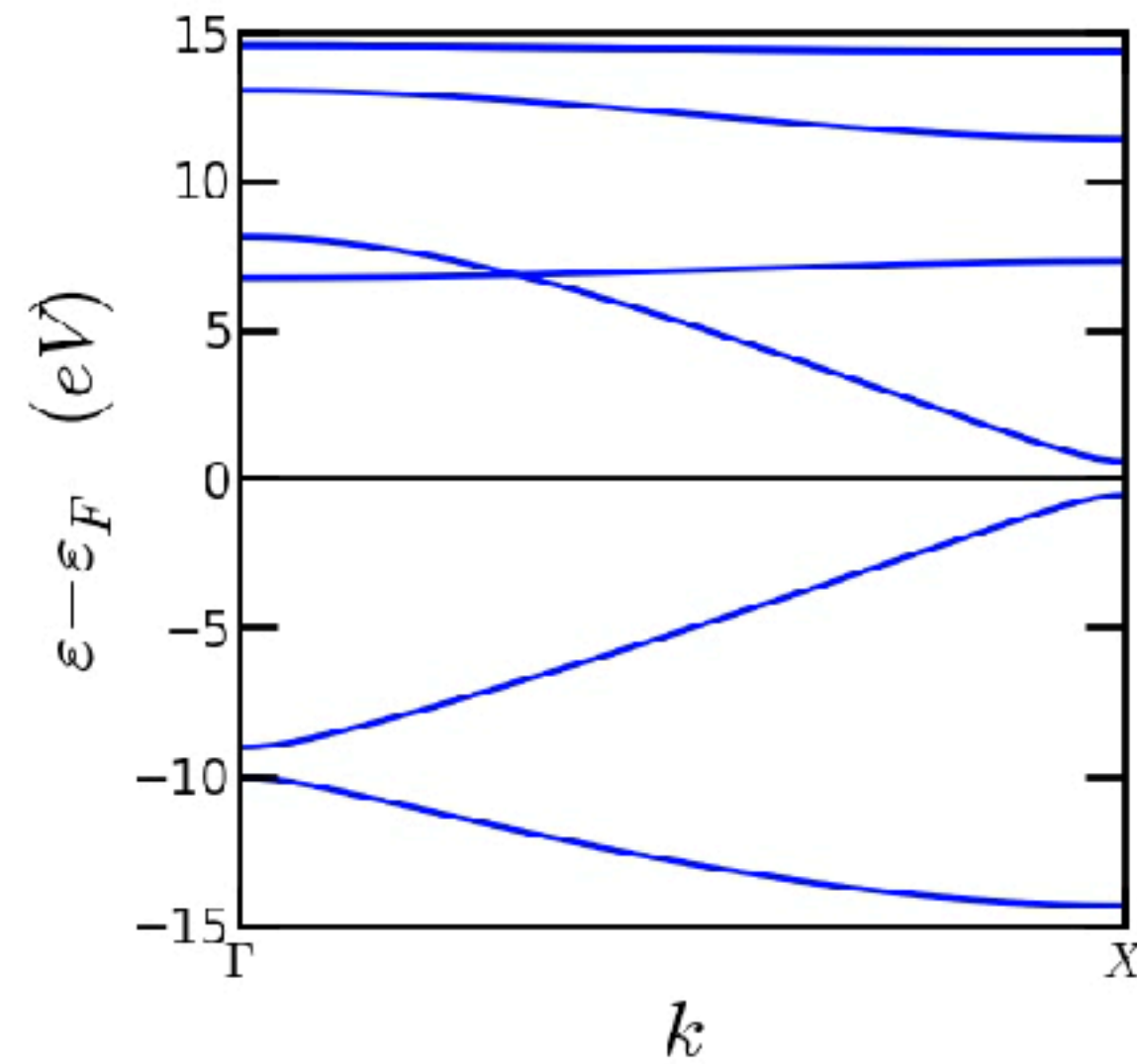
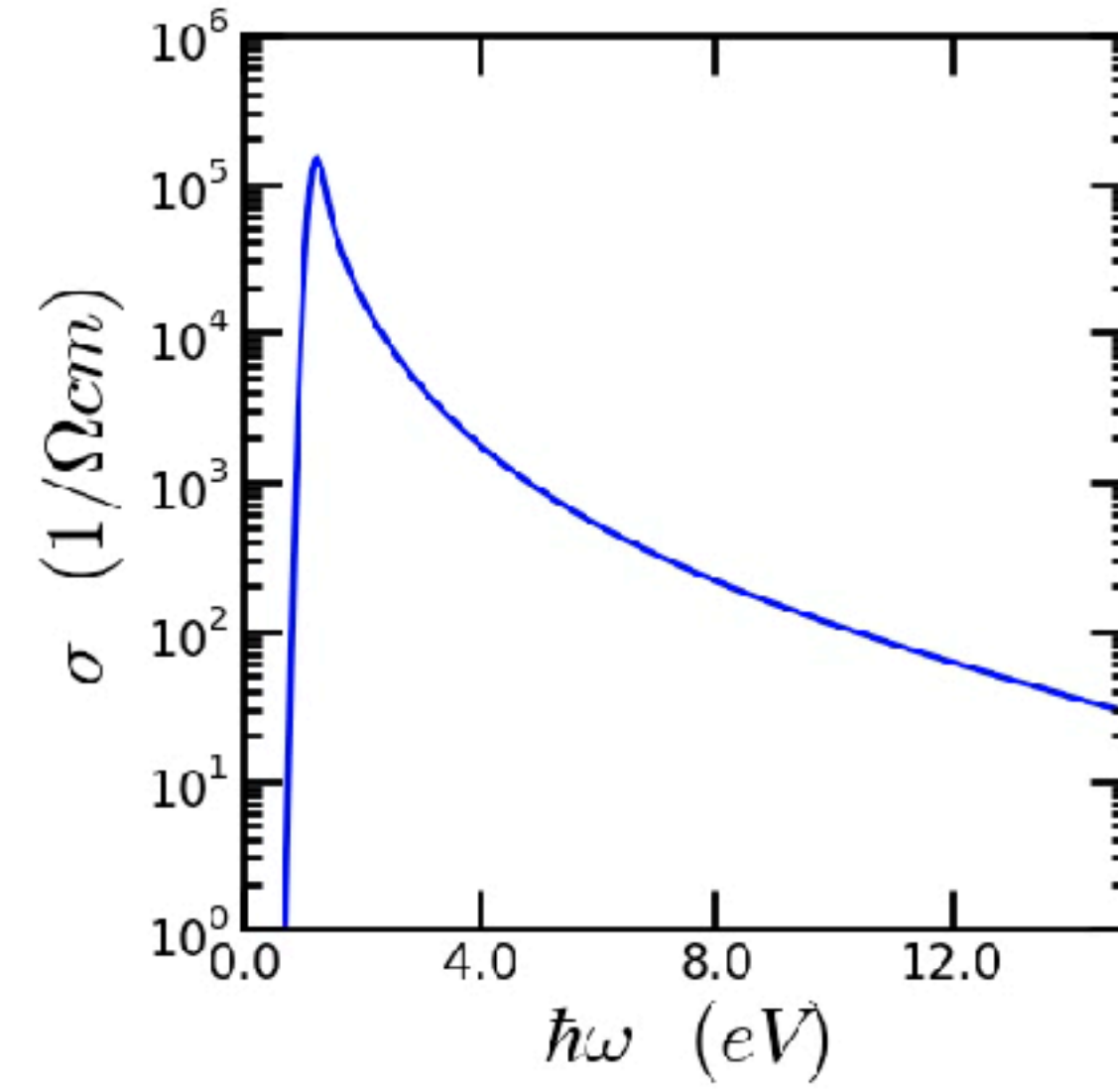
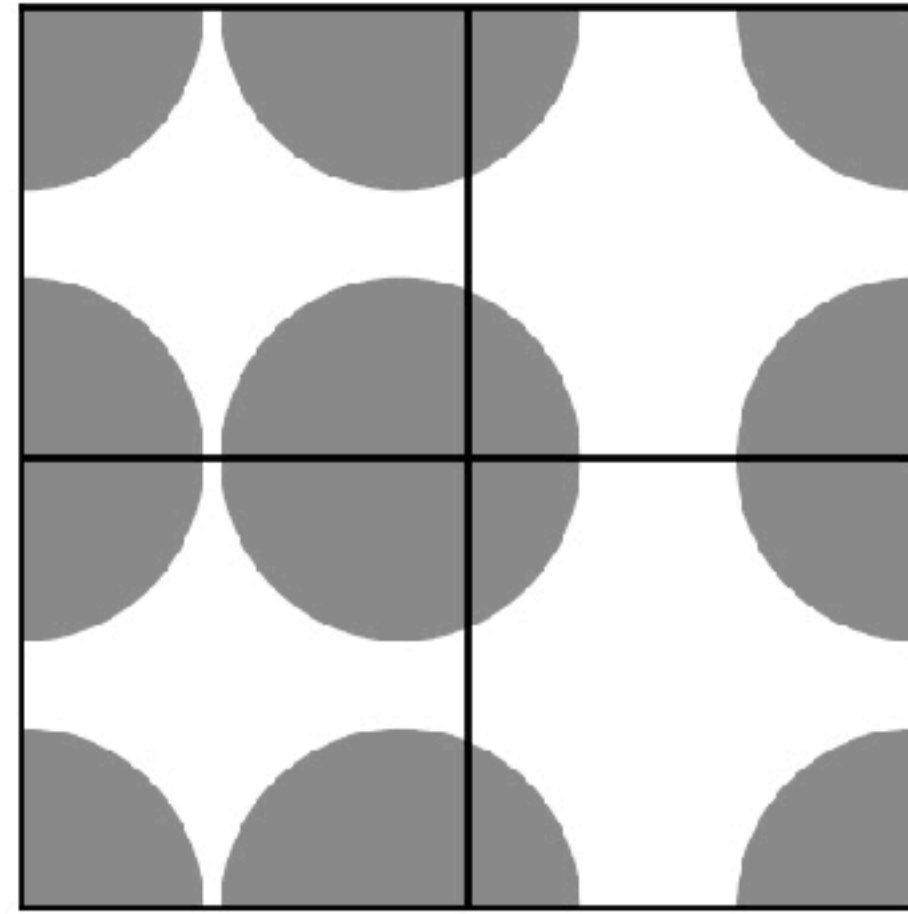


Thermal Motion of the nuclei:

Phonons momentarily break the **symmetry** and thus allow the **direct transitions** to become **active**.

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D.A. Greenwood, *Proc. Phys. Soc.* **71**, 585 (1958).



GREENWOOD-KUBO FORMALISM

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$$\sigma(\omega) = \frac{e^2 \hbar^2}{m_e^2 \omega} \frac{2\pi}{V} \left\langle \sum_{n, n \neq m} \sum_{\mathbf{k}} w_{\mathbf{k}} [f(\varepsilon_n) - f(\varepsilon_m)] |\langle n\mathbf{k} | \nabla | m\mathbf{k} \rangle|^2 \delta(\varepsilon_n - \varepsilon_m - \hbar\omega) \right\rangle_T$$

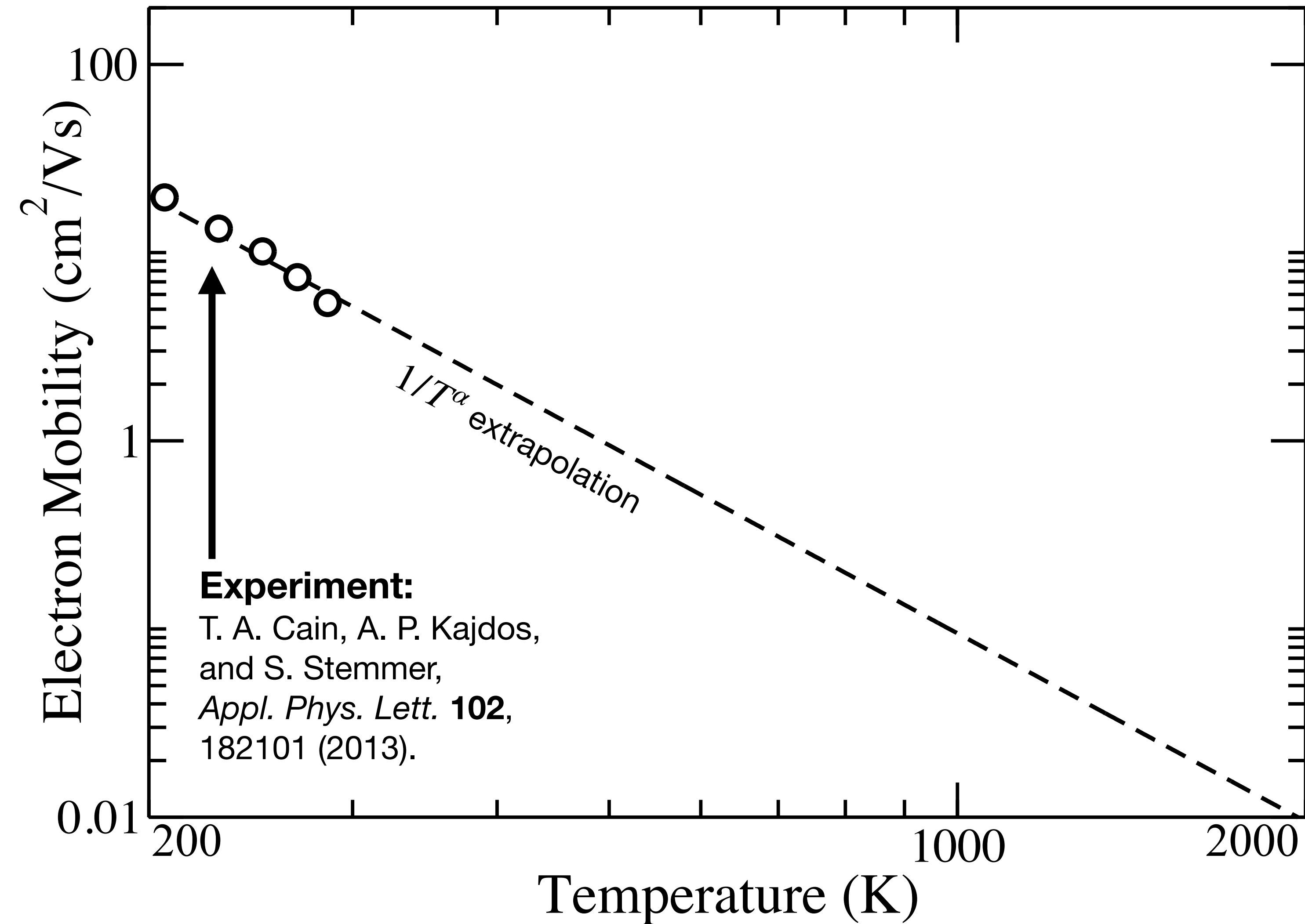
B. Holst, M. French, and R. Redmer, *Phys. Rev. B* **83**, 235120 (2011).

Challenges:

- **Extremely dense k-grids needed to resolve $\omega \rightarrow 0$ limit**
- **Large supercell needed to resolve phonon dispersion**

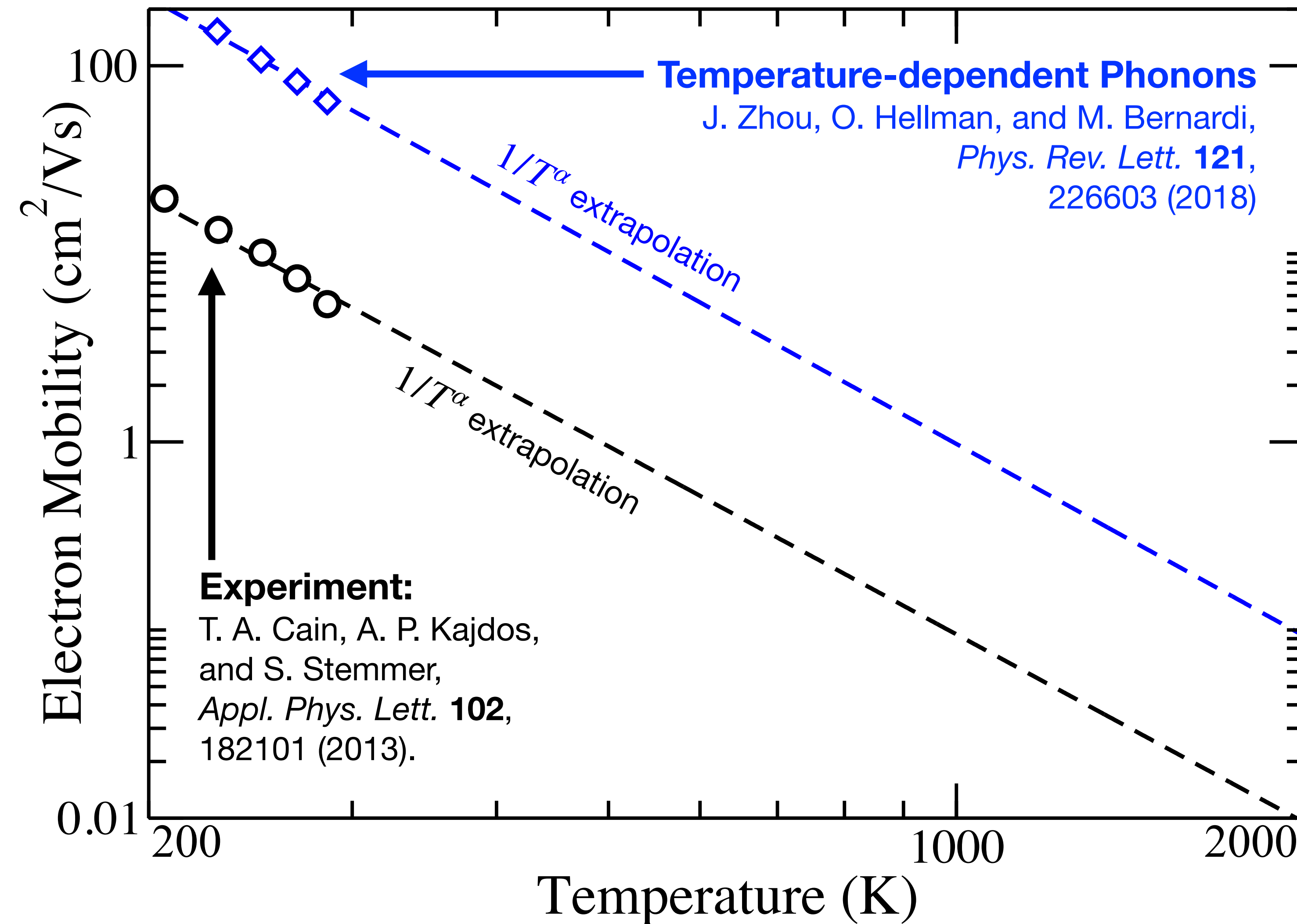
Mobility in Highly Anharmonic SrTiO₃

Doping: $8 \times 10^{17} \text{ cm}^{-3}$



Mobility in Highly Anharmonic SrTiO₃

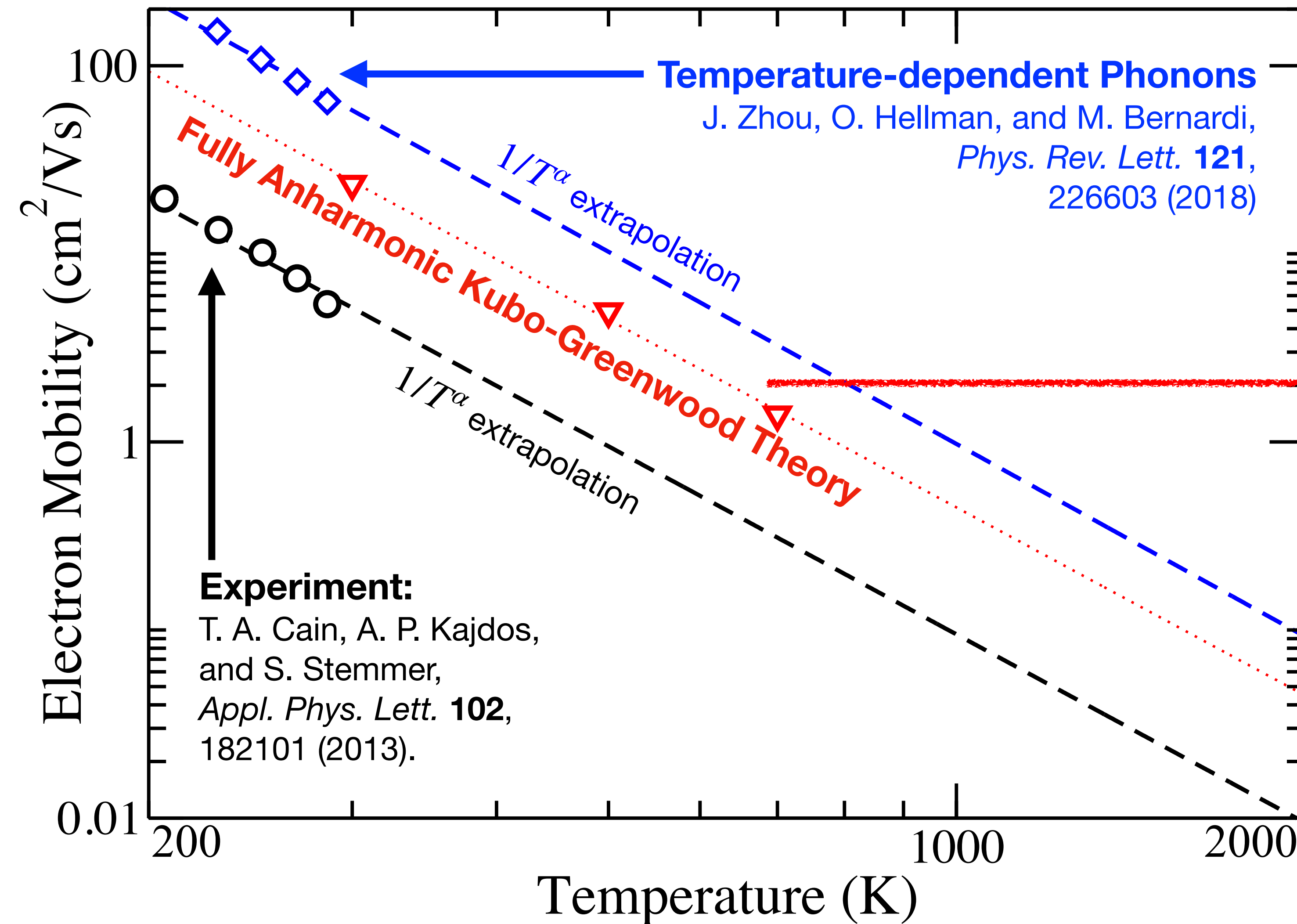
Doping: $8 \times 10^{17} \text{ cm}^{-3}$



*T-dependent Phonons +
Perturbative El.-Pho. Coupling
reproduces T-dependence of
the mobility at low temperatures.*

Mobility in Highly Anharmonic SrTiO₃

Doping: $8 \times 10^{17} \text{ cm}^{-3}$



T-dependent Phonons + Perturbative El.-Pho. Coupling reproduces T-dependence of the mobility at low temperatures.

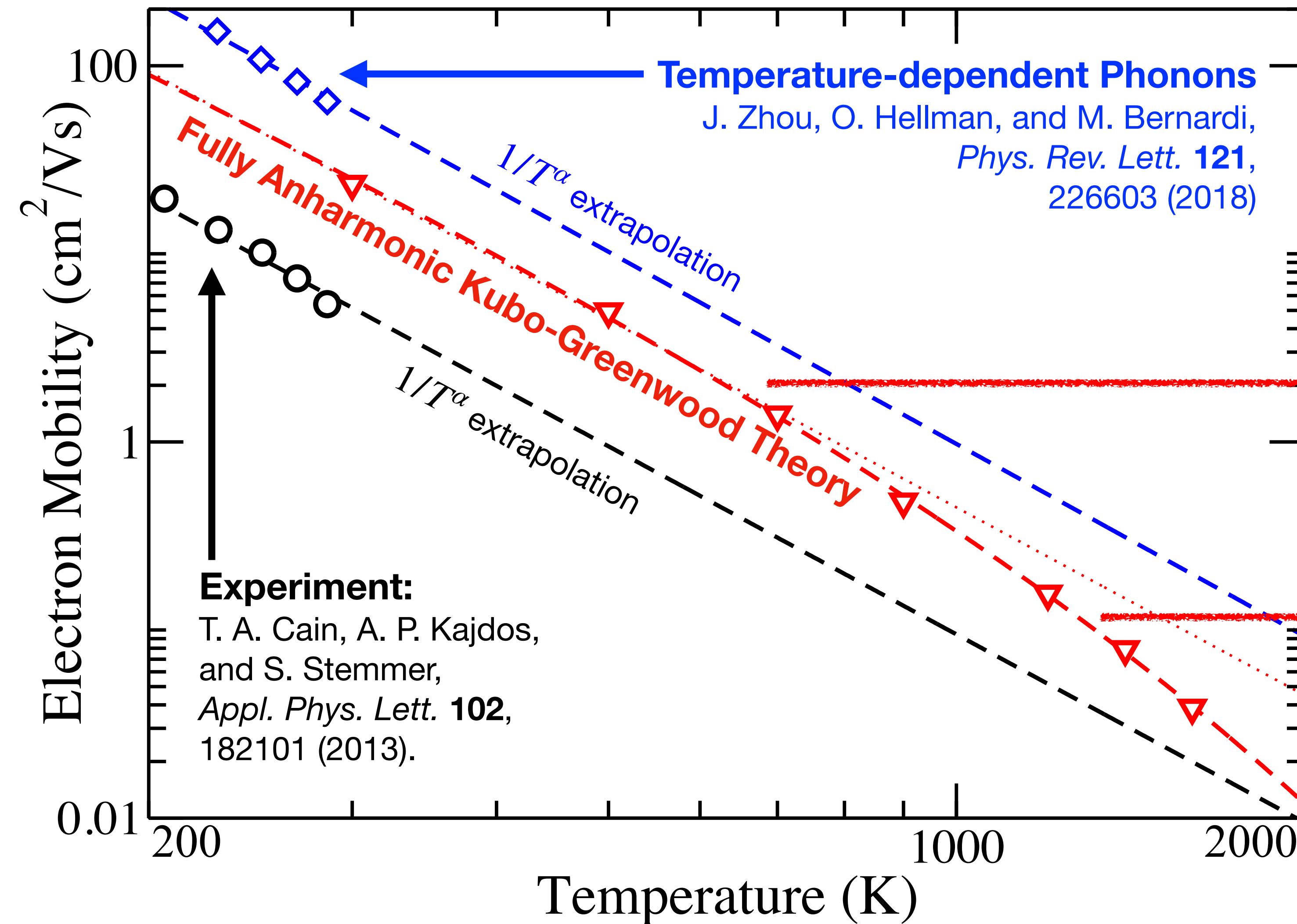
Fully-anharmonic KG Theory in line with literature results at low temperatures.

Experiment:
T. A. Cain, A. P. Kajdos,
and S. Stemmer,
Appl. Phys. Lett. **102**,
182101 (2013).

Temperature-dependent Phonons
J. Zhou, O. Hellman, and M. Bernardi,
Phys. Rev. Lett. **121**,
226603 (2018)

Mobility in Highly Anharmonic SrTiO₃

Doping: $8 \times 10^{17} \text{ cm}^{-3}$



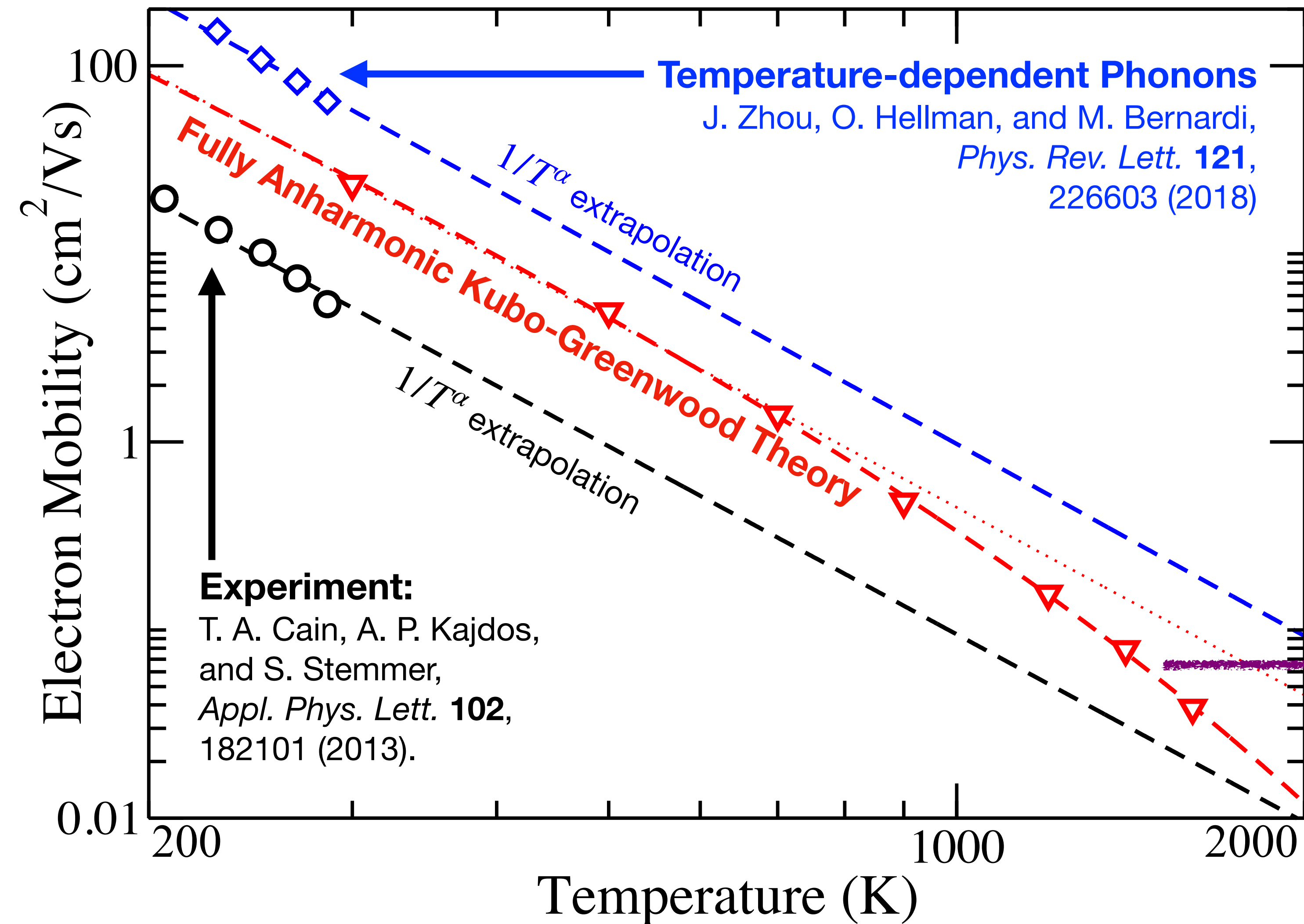
T-dependent Phonons +
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Fully-anharmonic KG Theory
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at low temperatures.

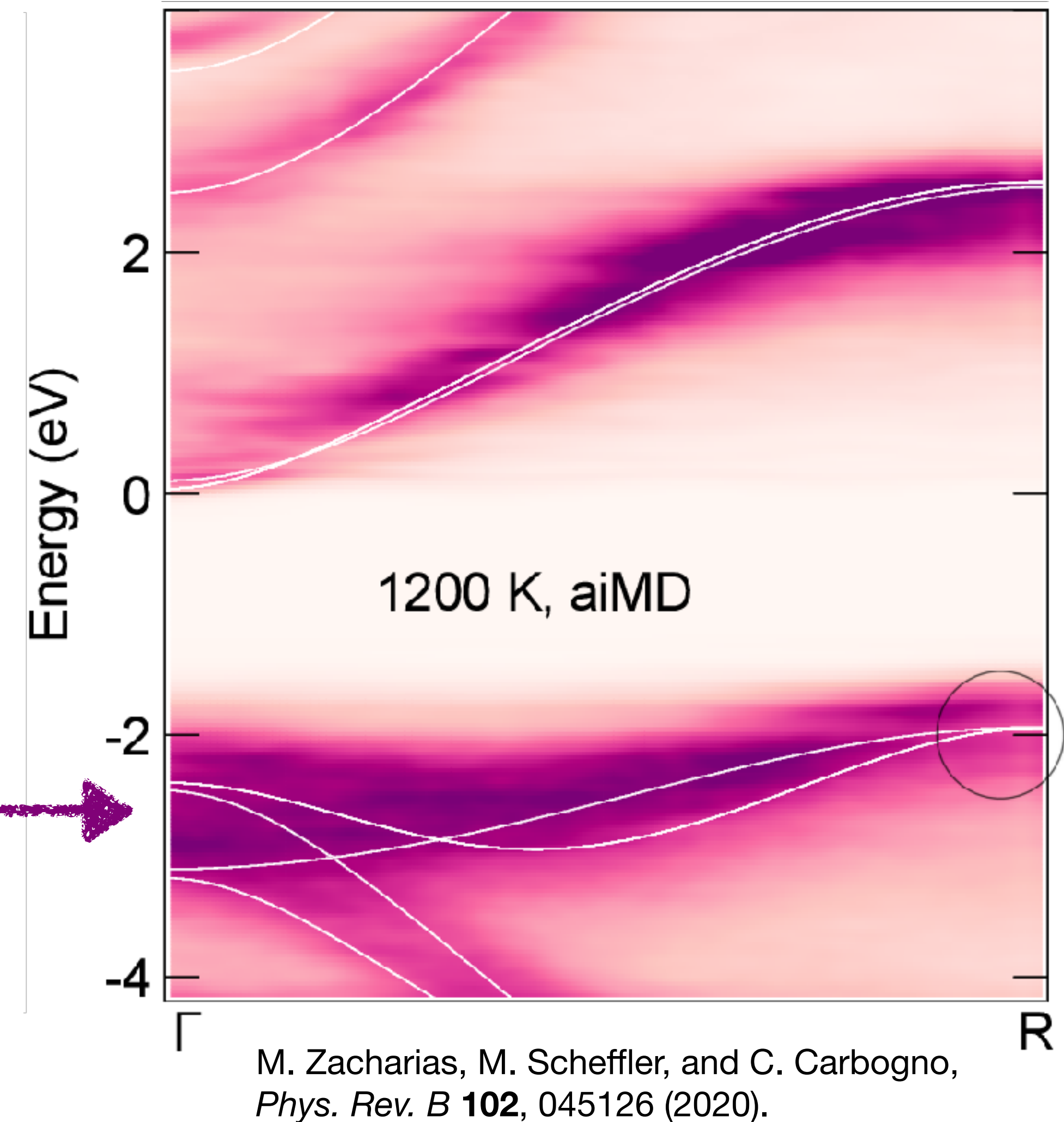
Fully-anharmonic KG Theory
predicts deviations from $1/T^\alpha$
behaviour at *high temperatures*.

Mobility in Highly Anharmonic SrTiO₃

Doping: $8 \times 10^{17} \text{ cm}^{-3}$

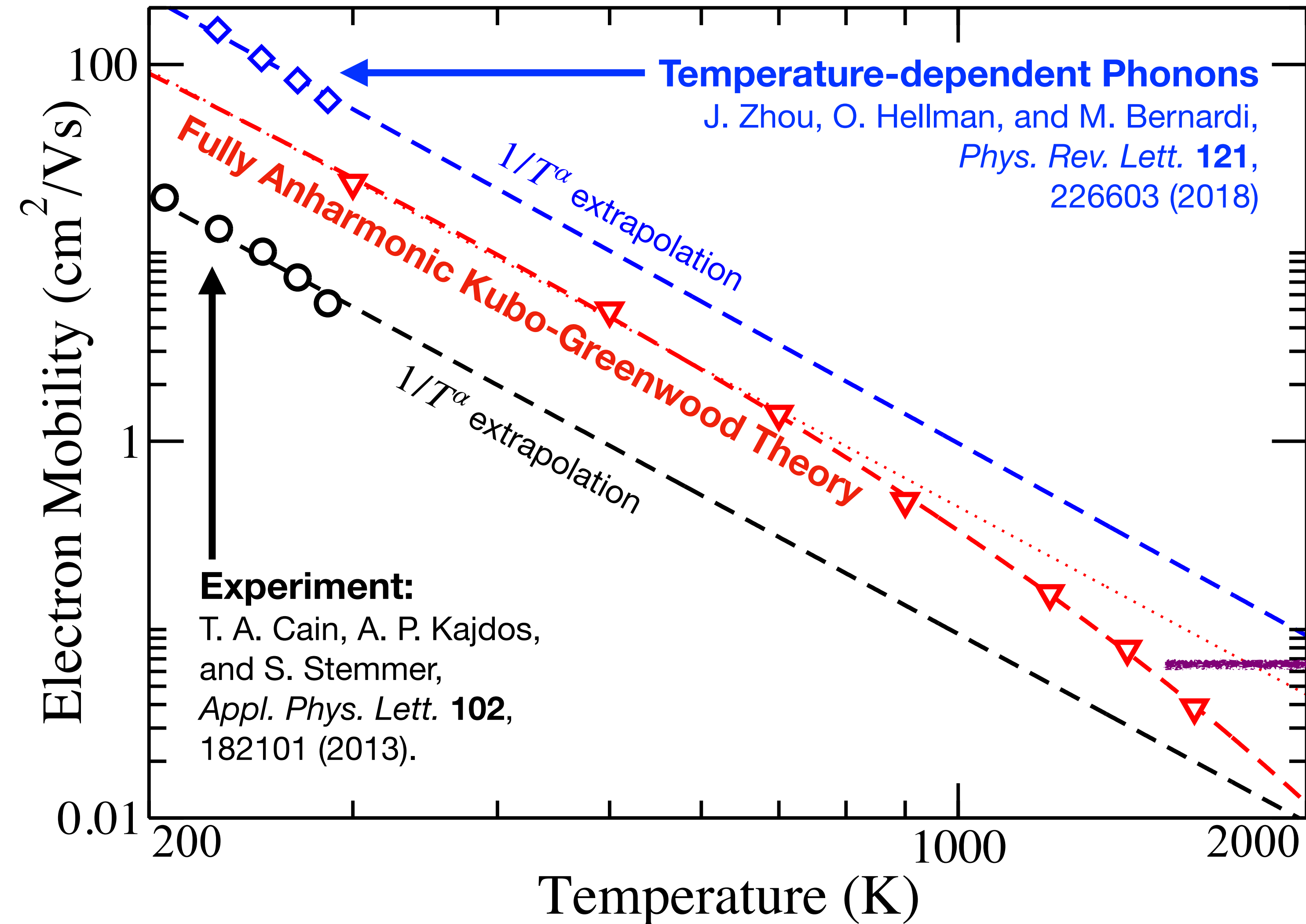


Fully Anharmonic Spectral Function

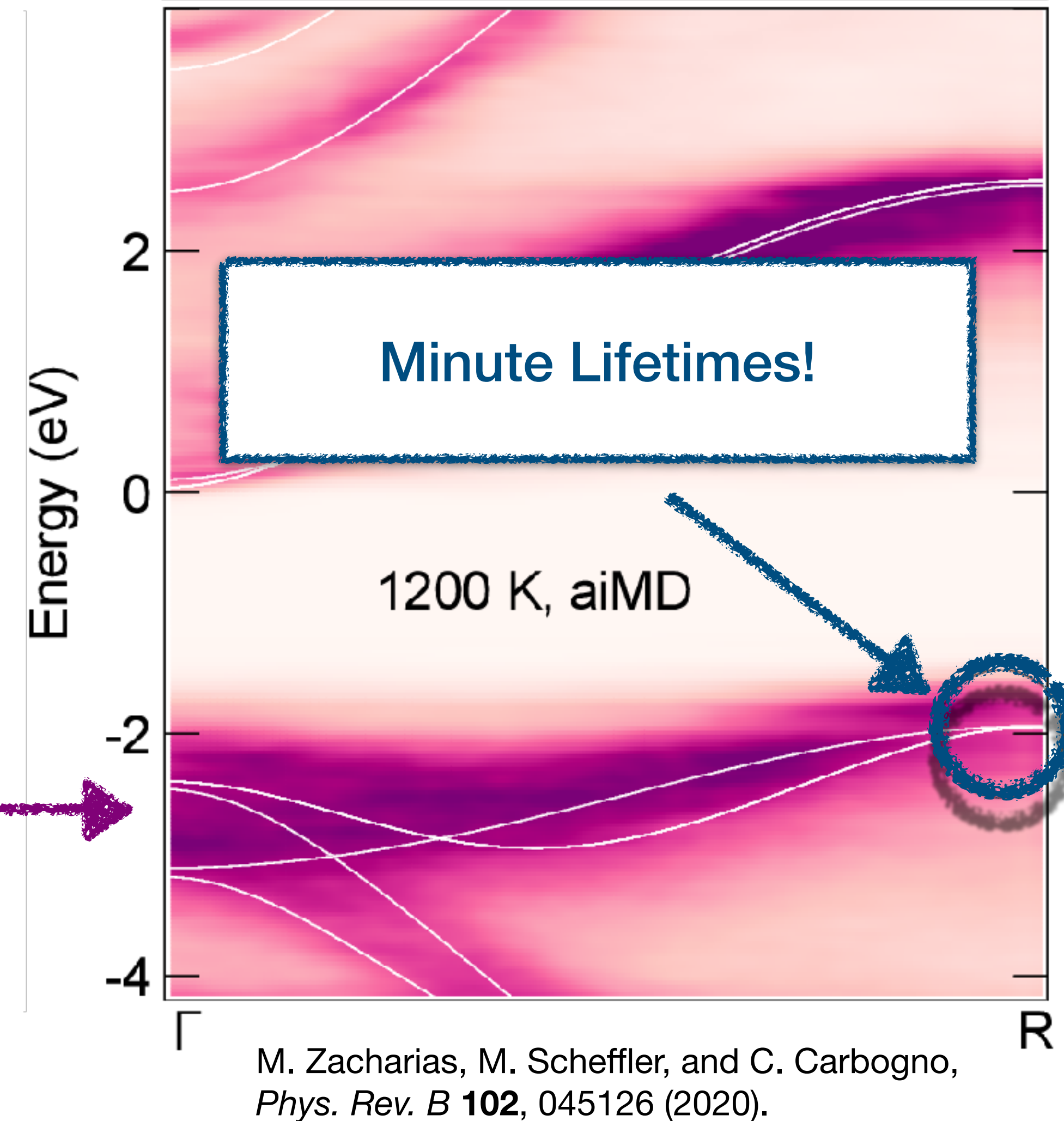


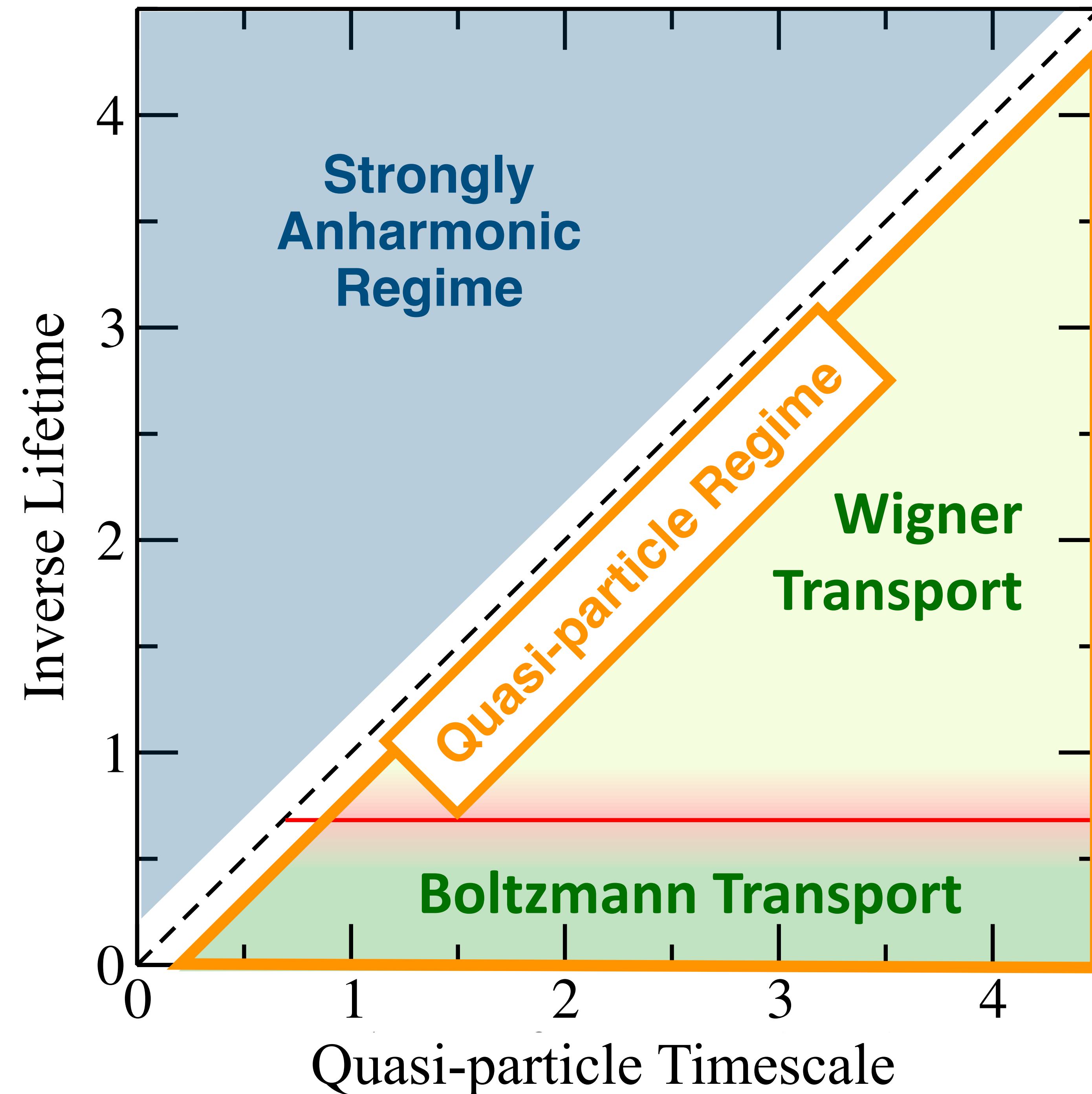
Mobility in Highly Anharmonic SrTiO₃

Doping: $8 \times 10^{17} \text{ cm}^{-3}$



Fully Anharmonic Spectral Function





Summary

- Going beyond the **quasi-particle picture** is crucial for computing **transport coefficients** in **complex** materials at **realistic conditions**.
- For **vibrational** transport, the implemented **Green-Kubo** formalism now allows the routine assessment of **fully anharmonic** thermal conductivities.
- For **electronic** transport, the improved **Kubo-Greenwood** formalism now allows to obtain **converged** mobilities in **anharmonic** materials.

Acknowledgements



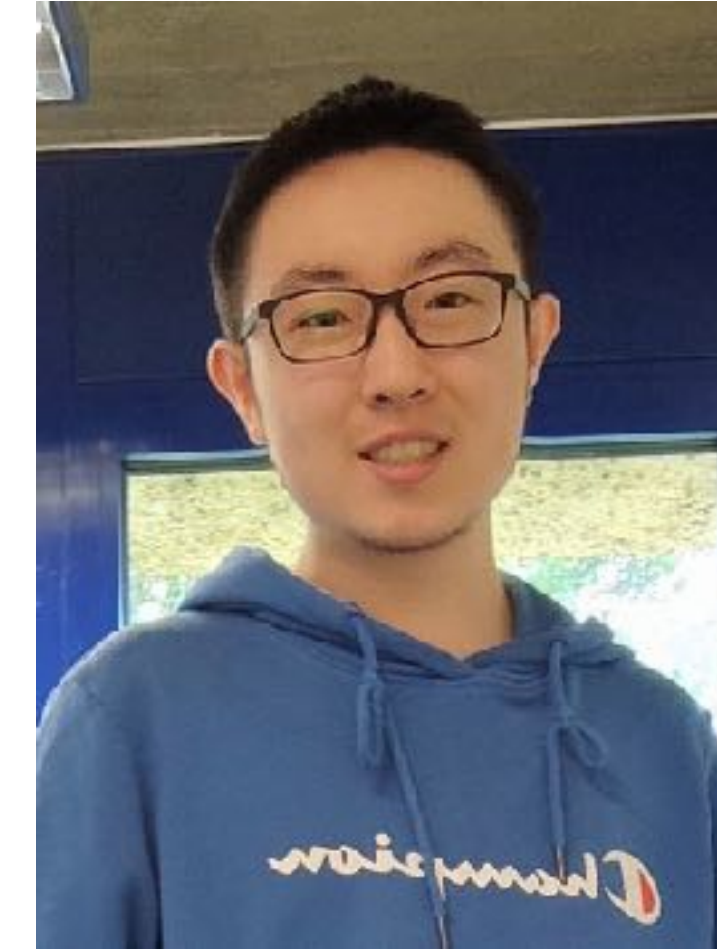
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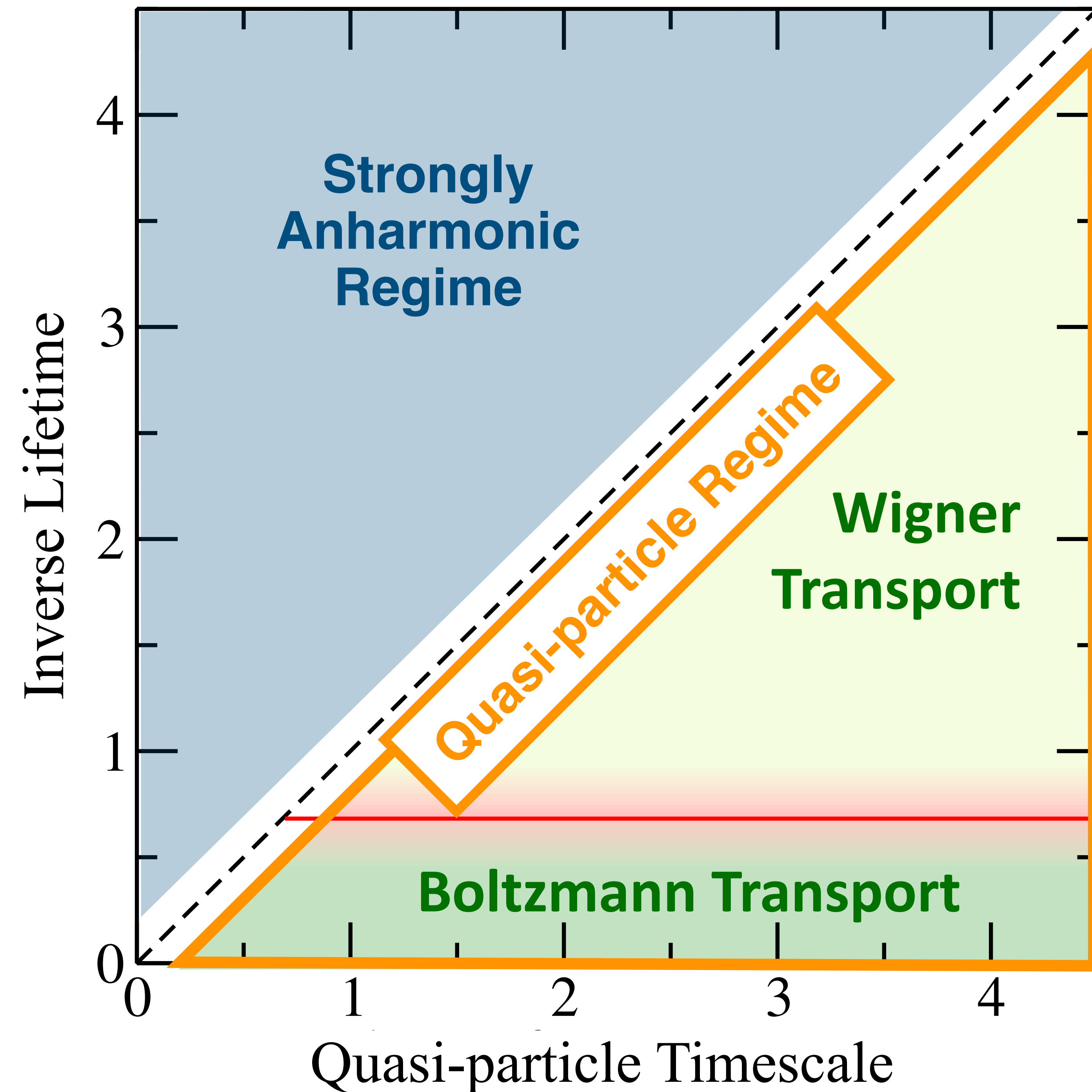
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Summary

- Going beyond the **quasi-particle picture** is crucial for computing **transport coefficients** in **complex** materials at **realistic conditions**.
- For **vibrational** transport, the implemented **Green-Kubo** formalism now allows the routine assessment of **fully anharmonic** thermal conductivities.
- For **electronic** transport, the improved **Kubo-Greenwood** formalism now allows to obtain **converged** mobilities in **anharmonic** materials.